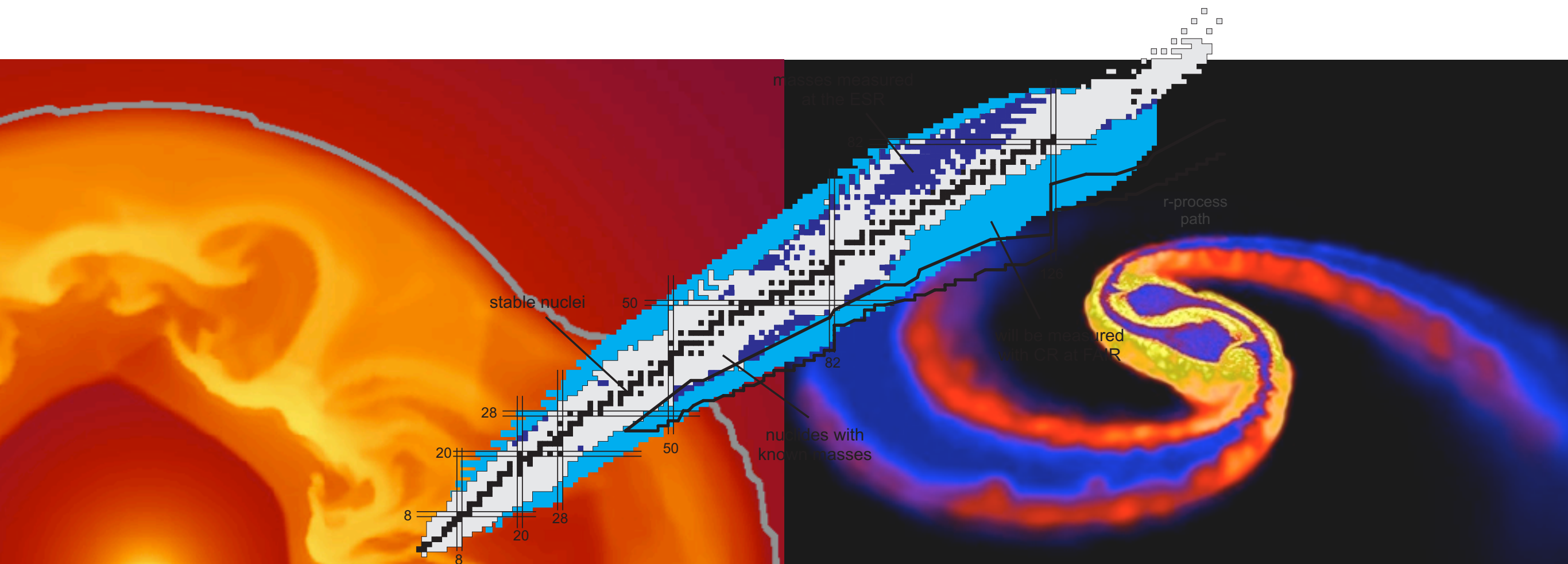


Impact of nuclear physics input on the r-process



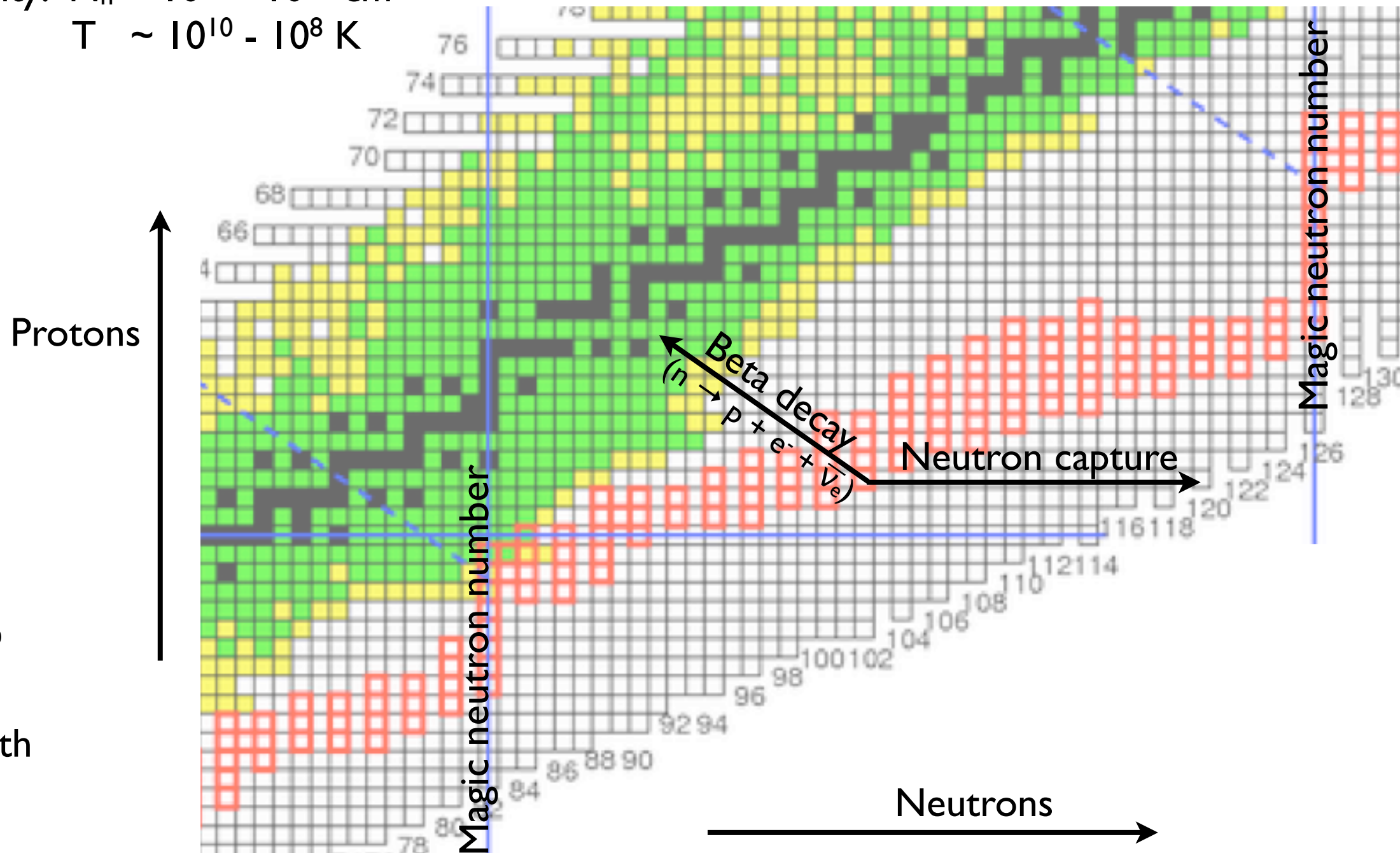
Almudena Arcones
Feodor Lynen Fellow, Basel University

r-process

Rapid neutron capture compared to beta decay

Neutron density: $N_n \sim 10^{27} - 10^{20} \text{ cm}^{-3}$

Temperature: $T \sim 10^{10} - 10^8 \text{ K}$



Ultra metal-poor stars = very old stars

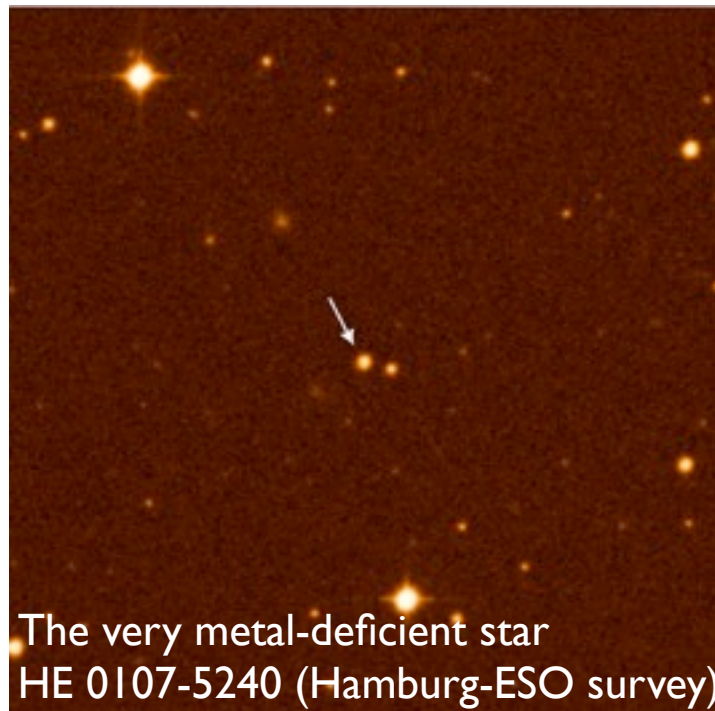
Their atmospheres show fingerprints of only few nucleosynthesis events that enriched the interstellar medium.

Abundances of r-process elements in:

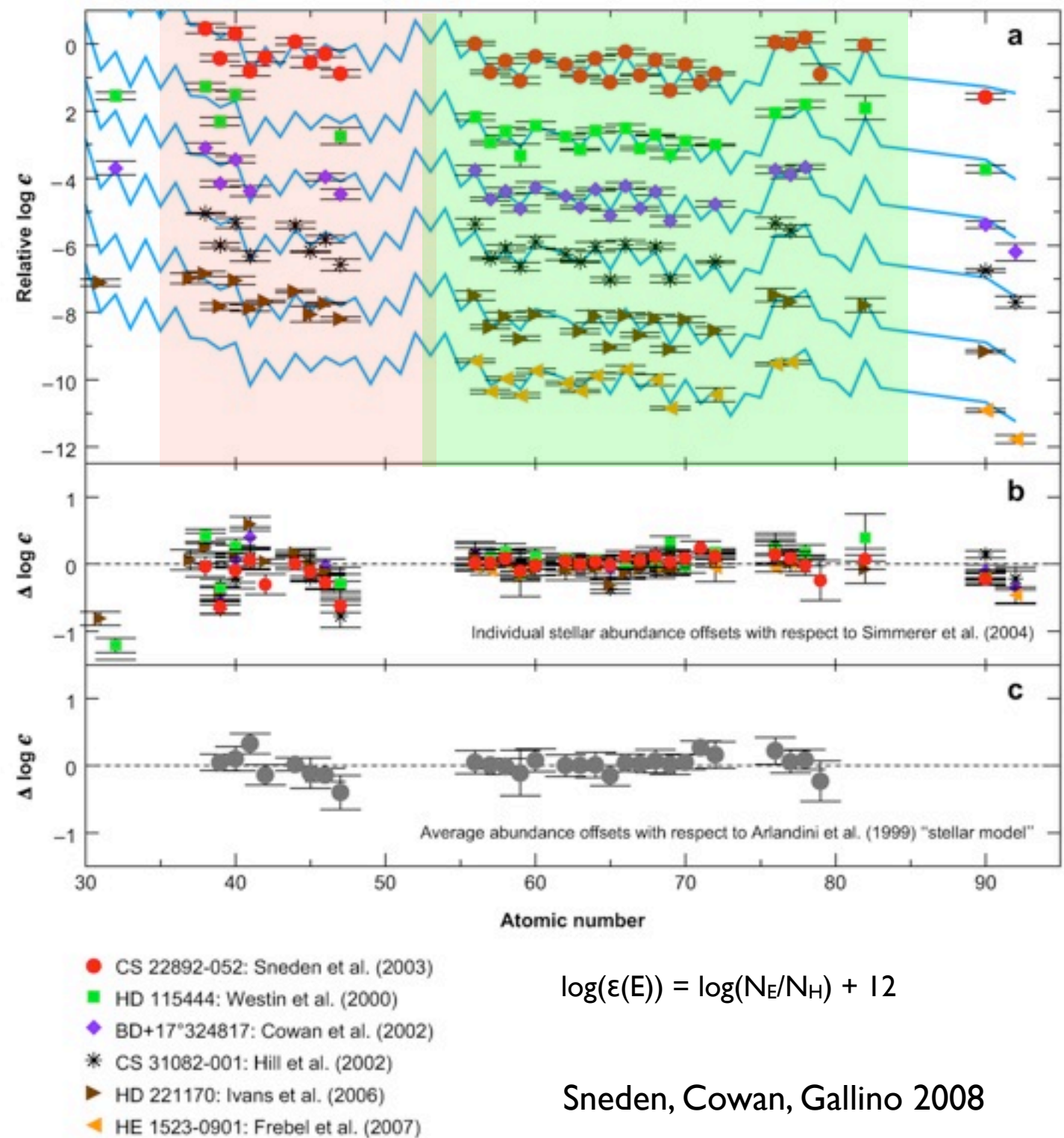
- ultra metal-poor stars and
- solar system

Two components or sites:

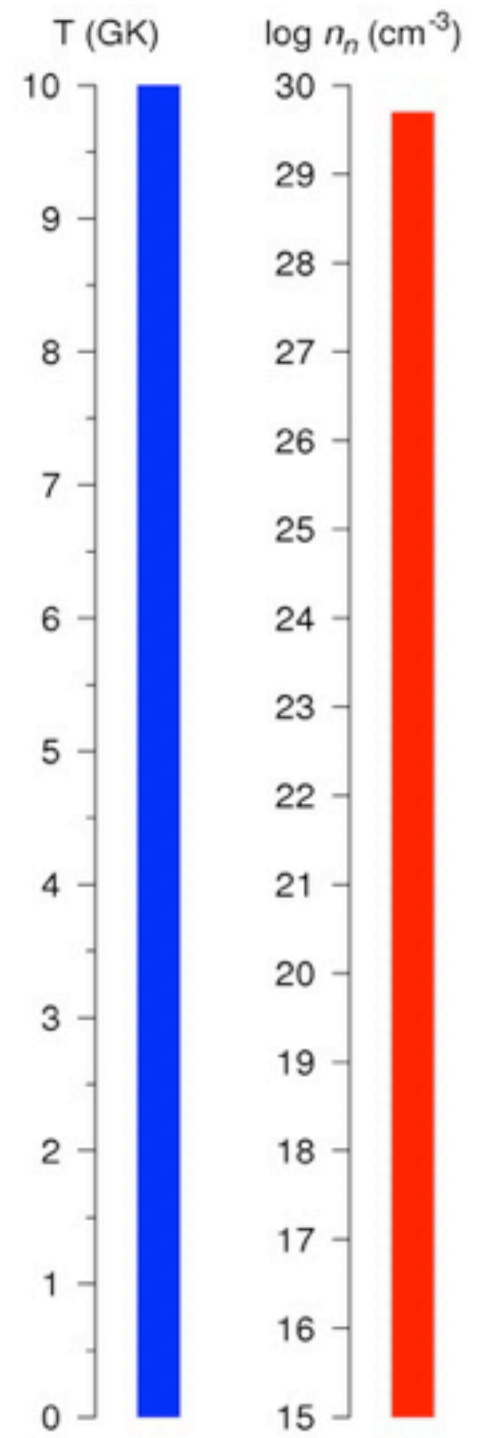
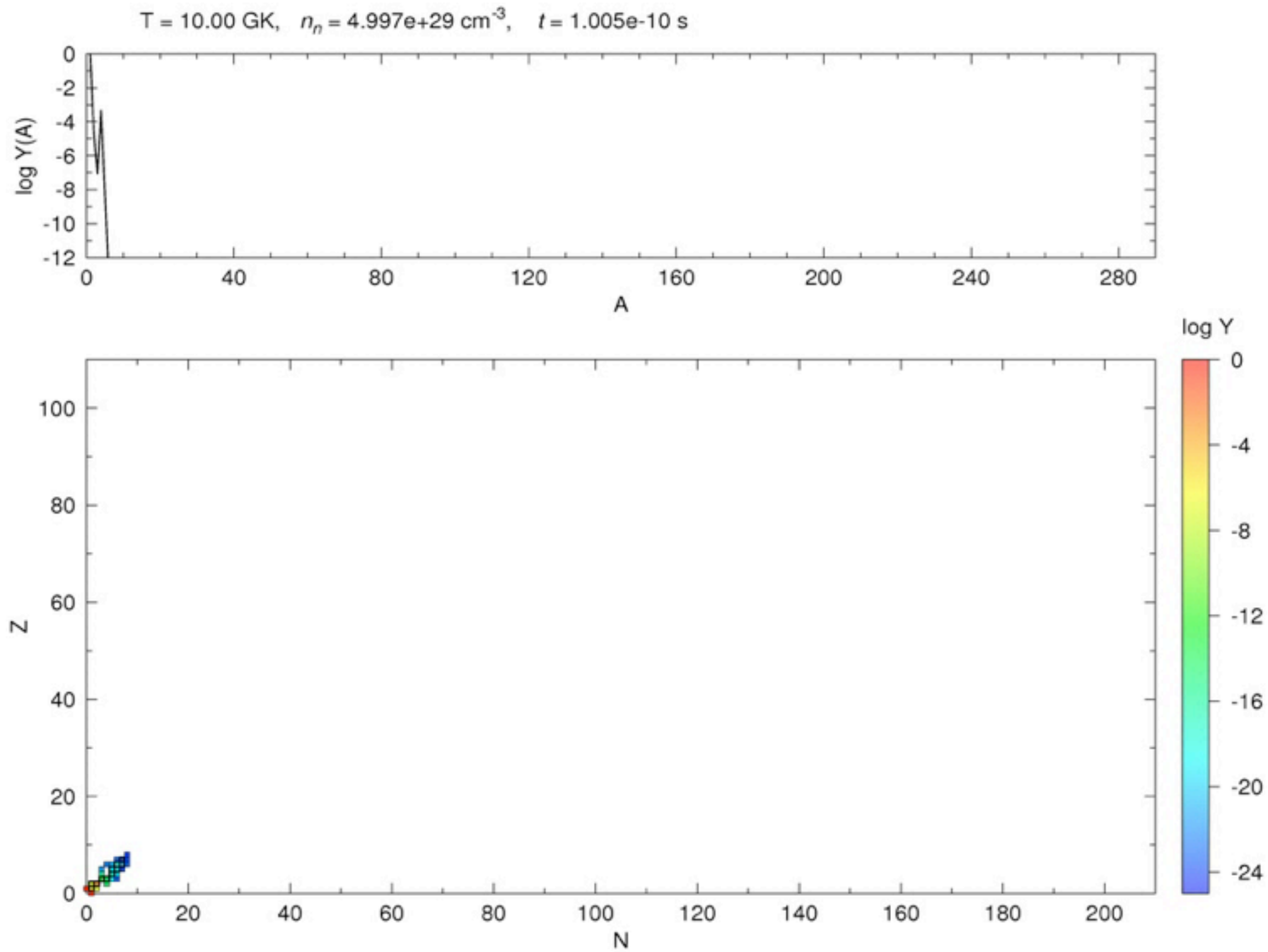
- robust r-process for $56 < Z < 83$
- scatter for lighter heavy elements $Z \sim 40$



The very metal-deficient star
HE 0107-5240 (Hamburg-ESO survey)

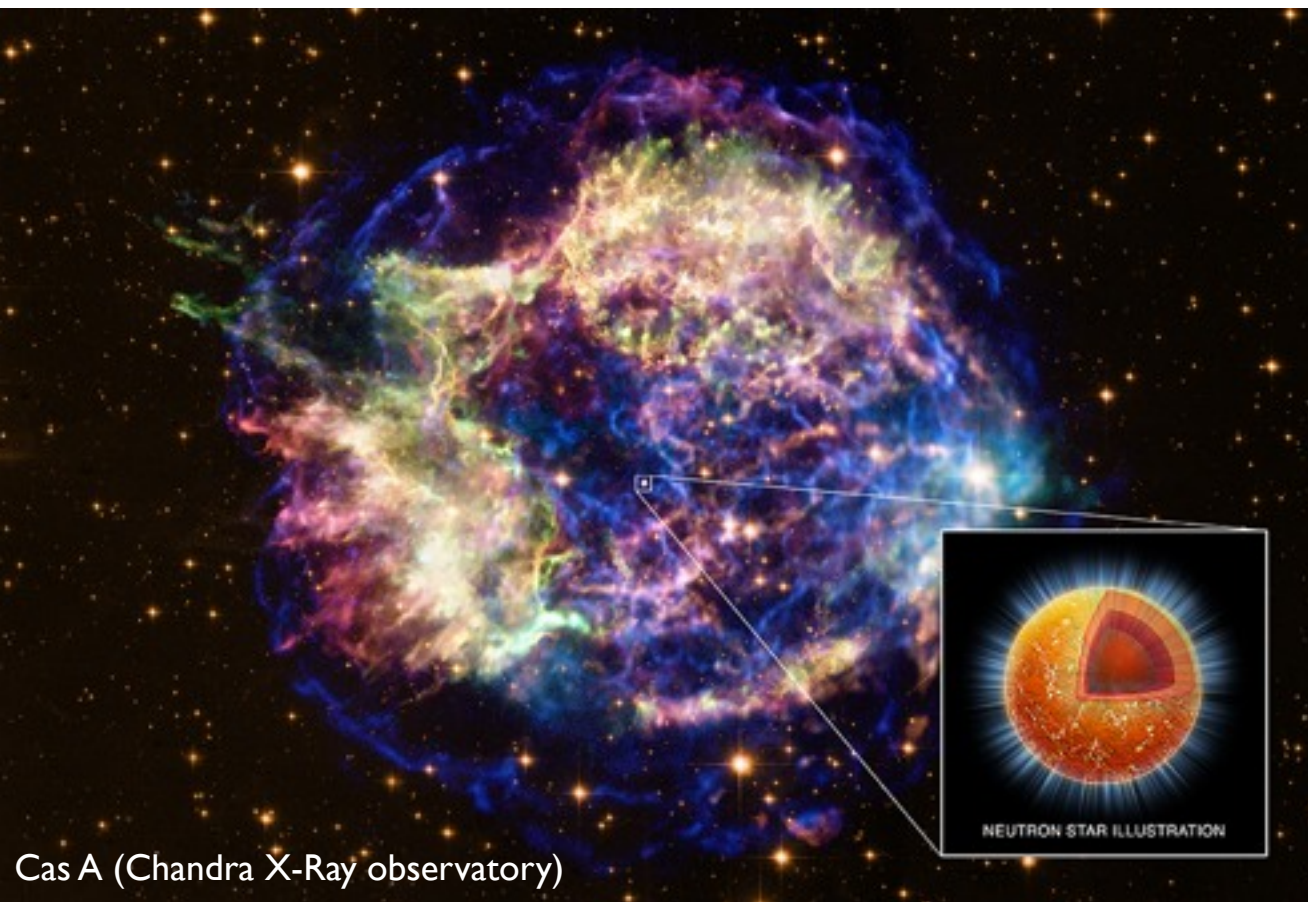


r-process



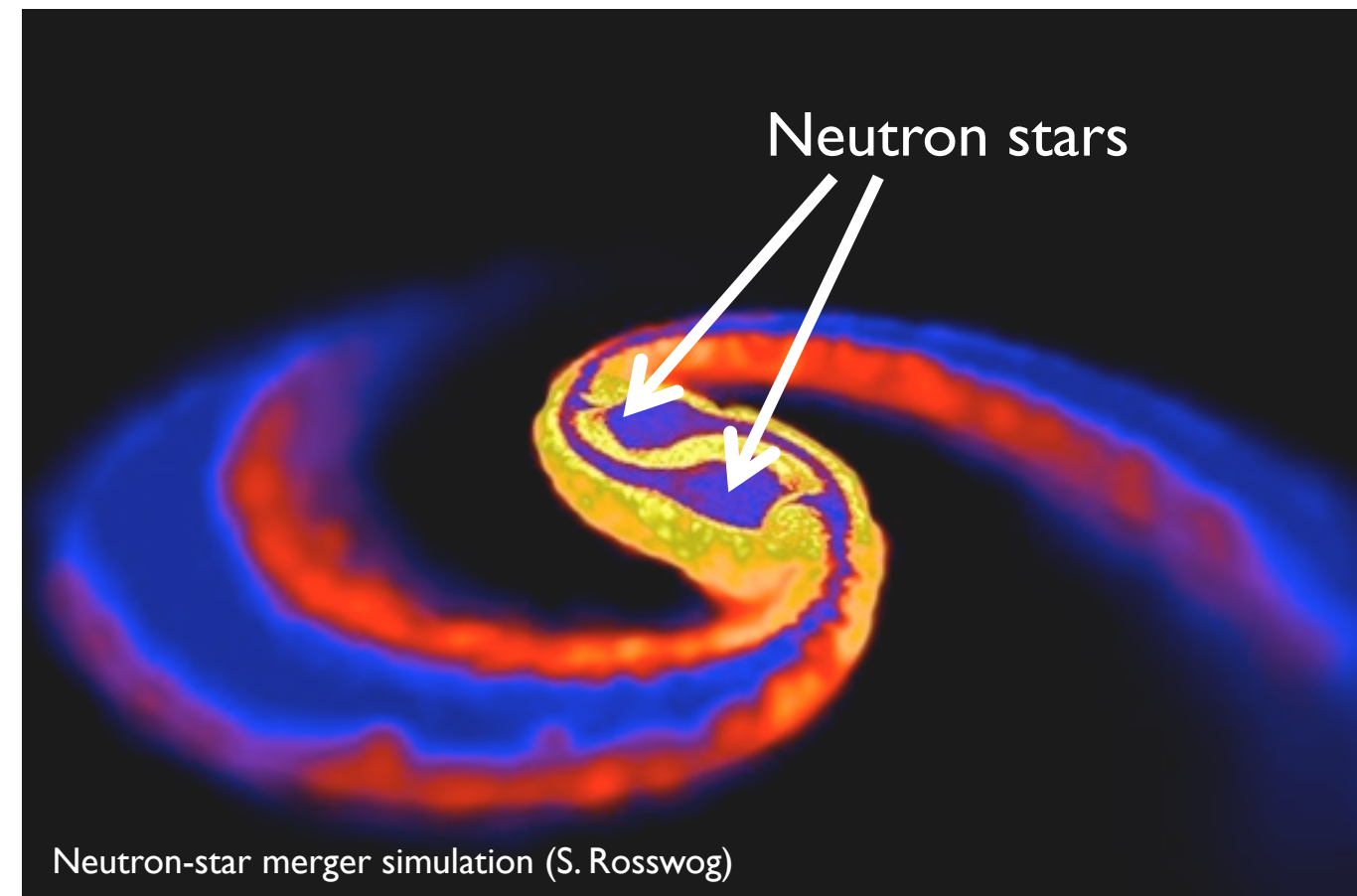
Where does the r-process occur?

Core-collapse supernovae



neutrino-driven wind (Woosley et al. 1994):
proton rich (Fischer et al. 2010, Hudepohl et al. 2010)
entropy too low (Woosley et al. 1994 → Roberts et al. 2010)
→ multidimensional effects,
neutrino collective oscillations, ...?

Neutron star mergers

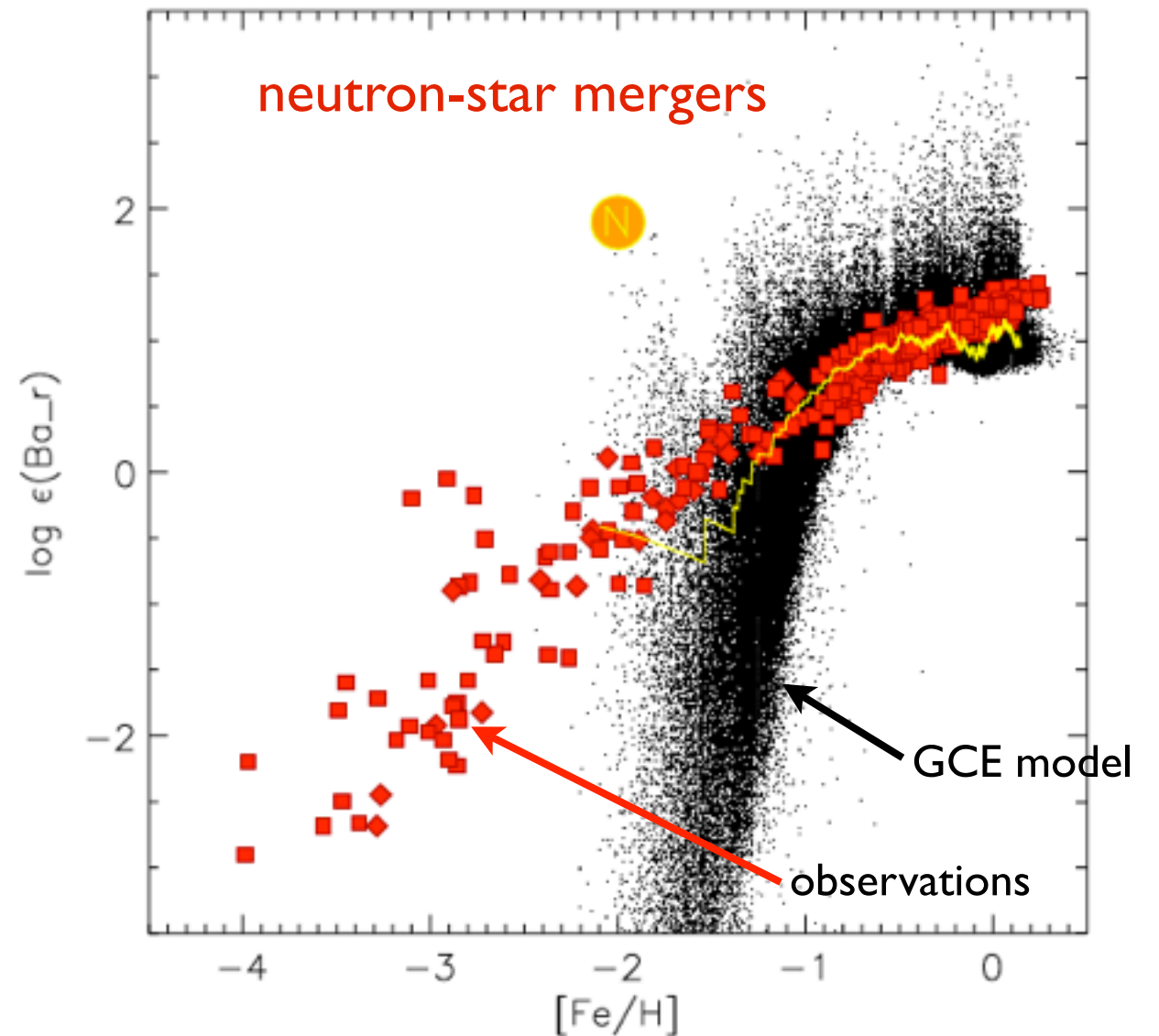
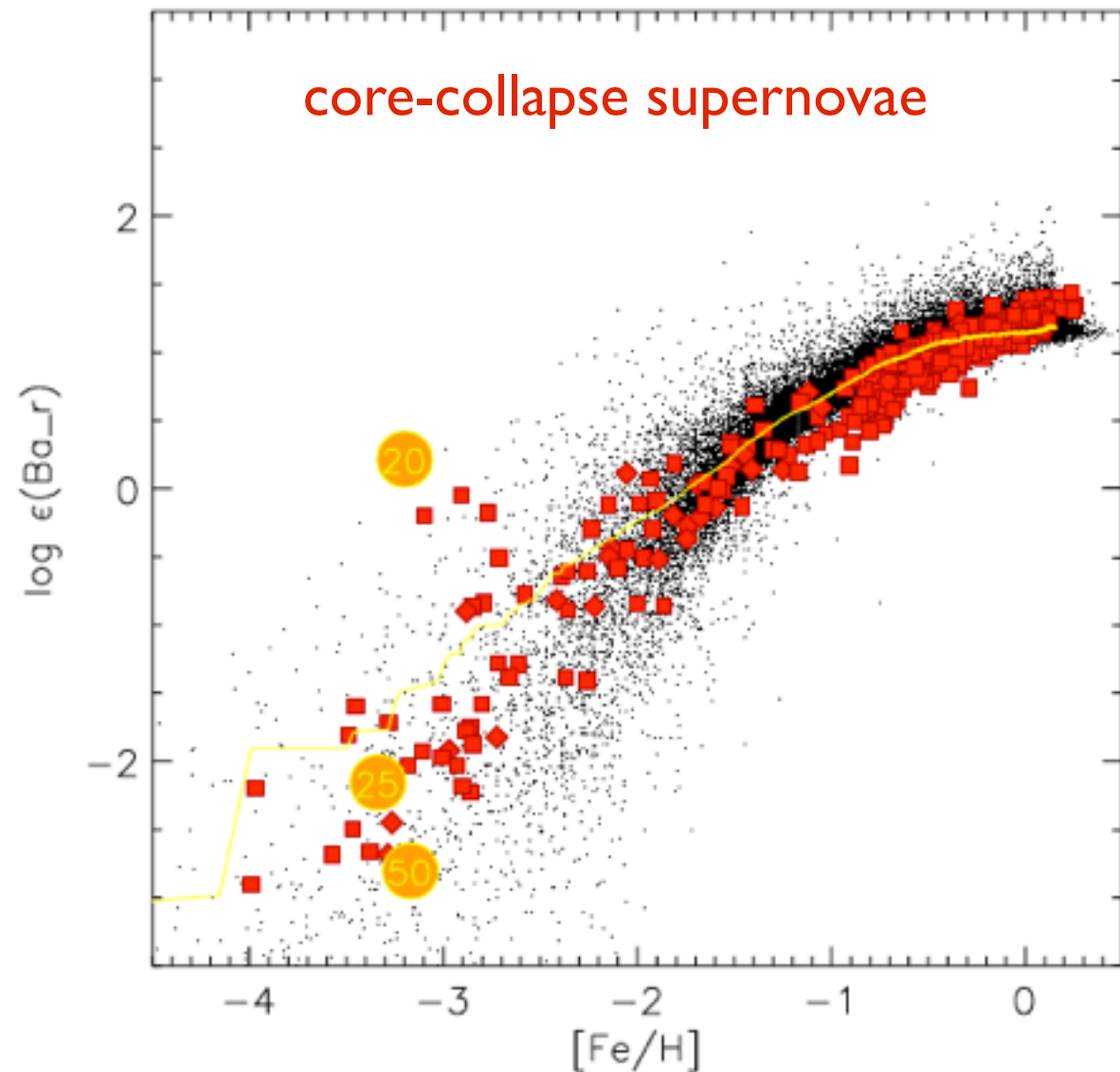


Right conditions for a successful r-process
(Freiburghaus et al. 1999, ..., Goriely et al. 2011)

They do not occur early enough to explain UMP
star abundances (Qian 2000, Argast et al. 2004)

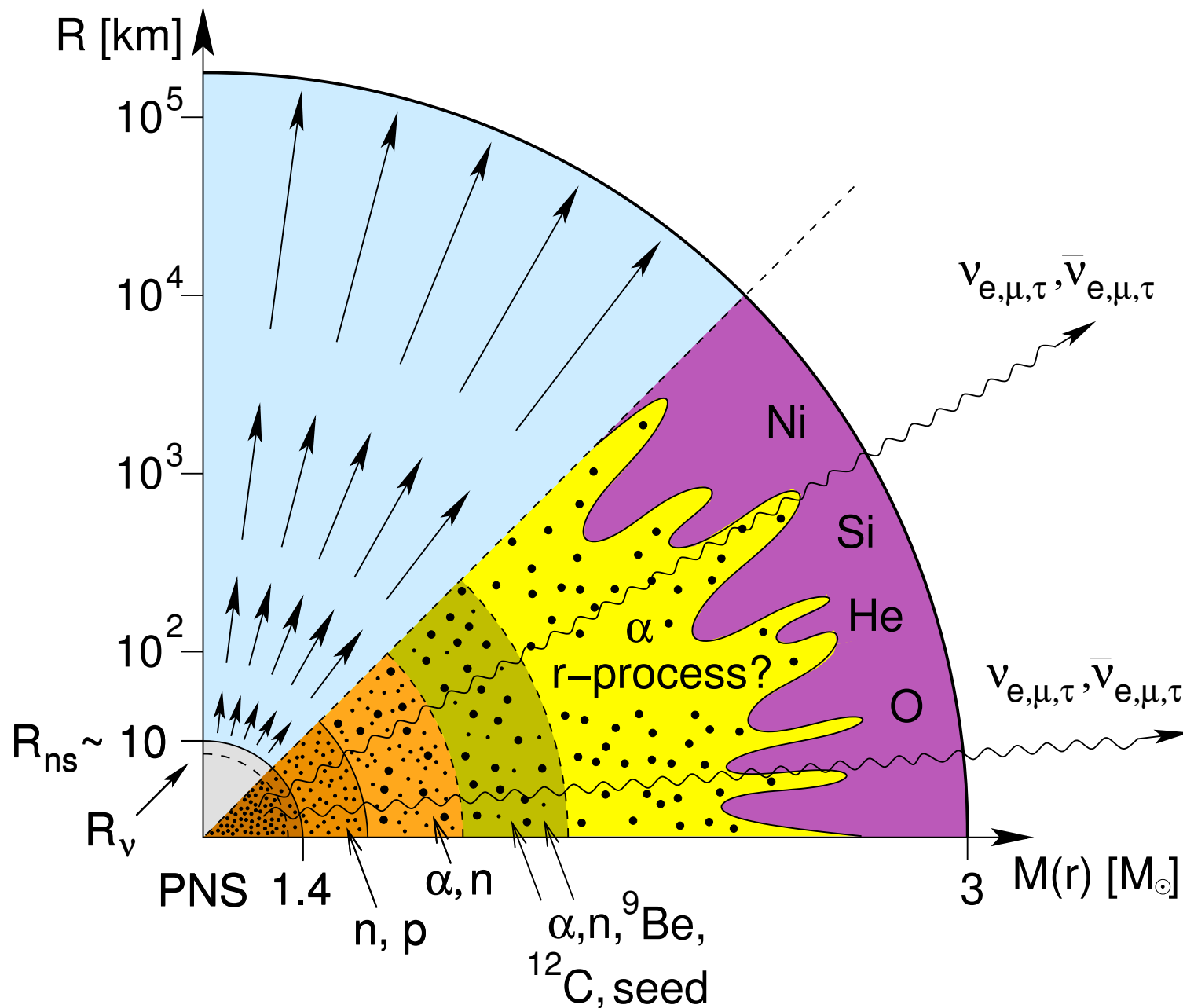
Chemical evolution: supernovae vs. mergers

Argast et al. 2004: galactic chemical evolution models r-process from:



Open questions: amount of mass ejected
event rate

Nucleosynthesis in neutrino-driven winds



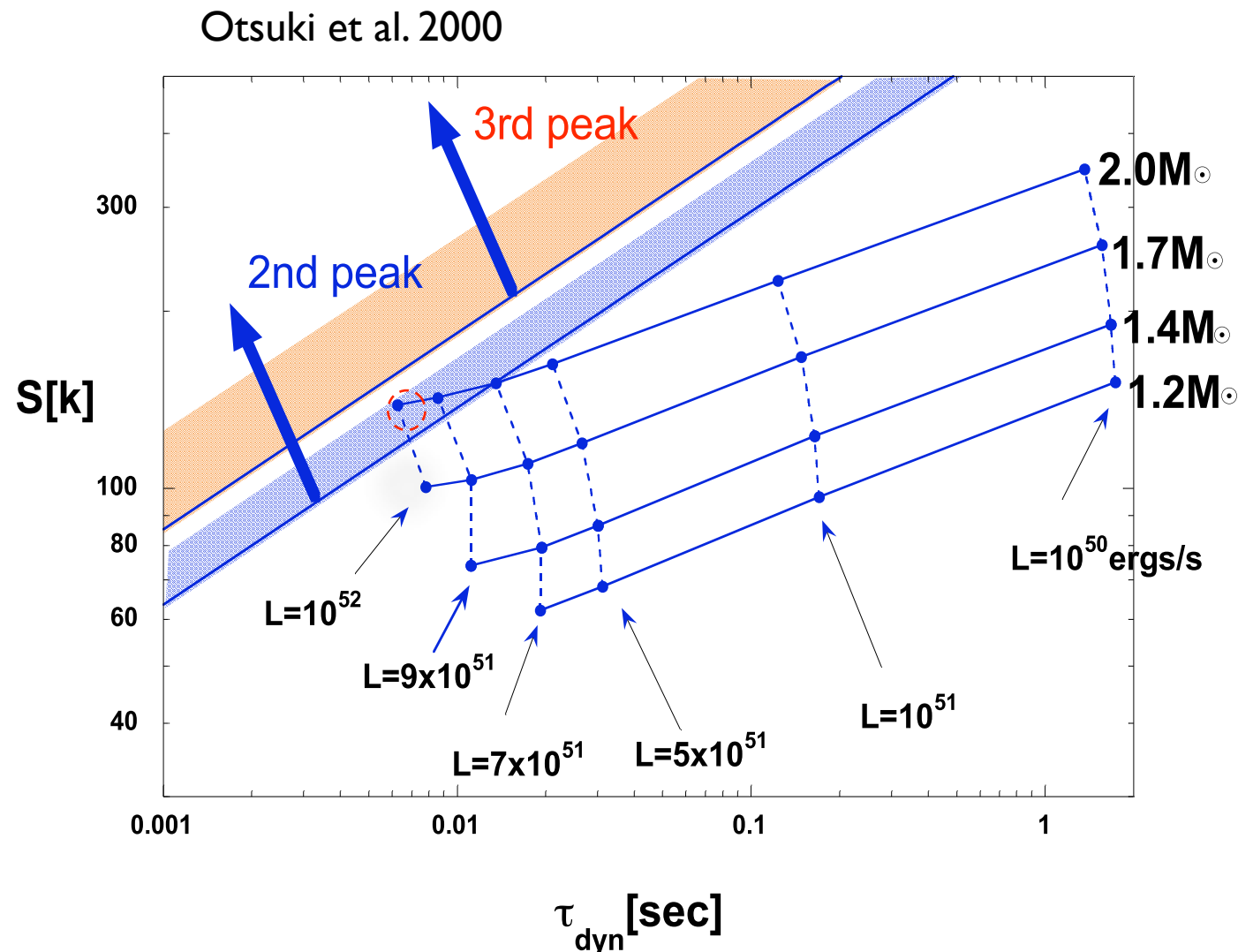
Production of heavy elements ($A > 130$) requires high neutron-to-seed ratio ($Y_n/Y_{\text{seed}} \sim 100$).

Necessary conditions for the r-process:

- **fast expansion**: inhibits the alpha-process and thus the formation of seed nuclei
- neutron rich ejecta: $Y_e < 0.5$
- **high entropy** is equivalent to high photon-to-baryon ratio. Photons dissociate seed nuclei into nucleons.

(Meyer et al. 1992, Hoffman et al. 1997, Otsuki et al. 2000, Thompson et al. 2001...)

Nucleosynthesis in neutrino-driven winds



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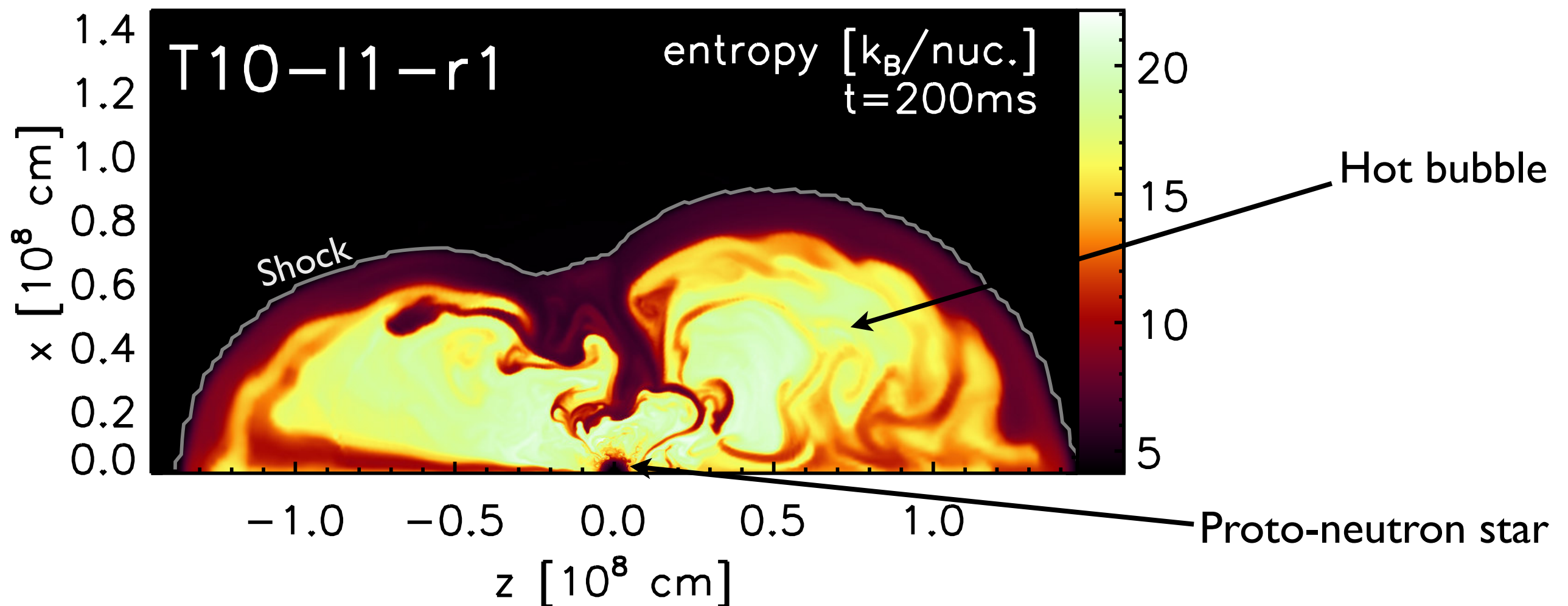
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Necessary conditions identified by steady-state models (e.g. Otsuki et al. 2000, Thompson et al. 2001) are not realized in recent simulations (Arcones et al. 2007, Fischer et al. 2010, Hudepohl et al. 2010, Roberts et al. 2010)

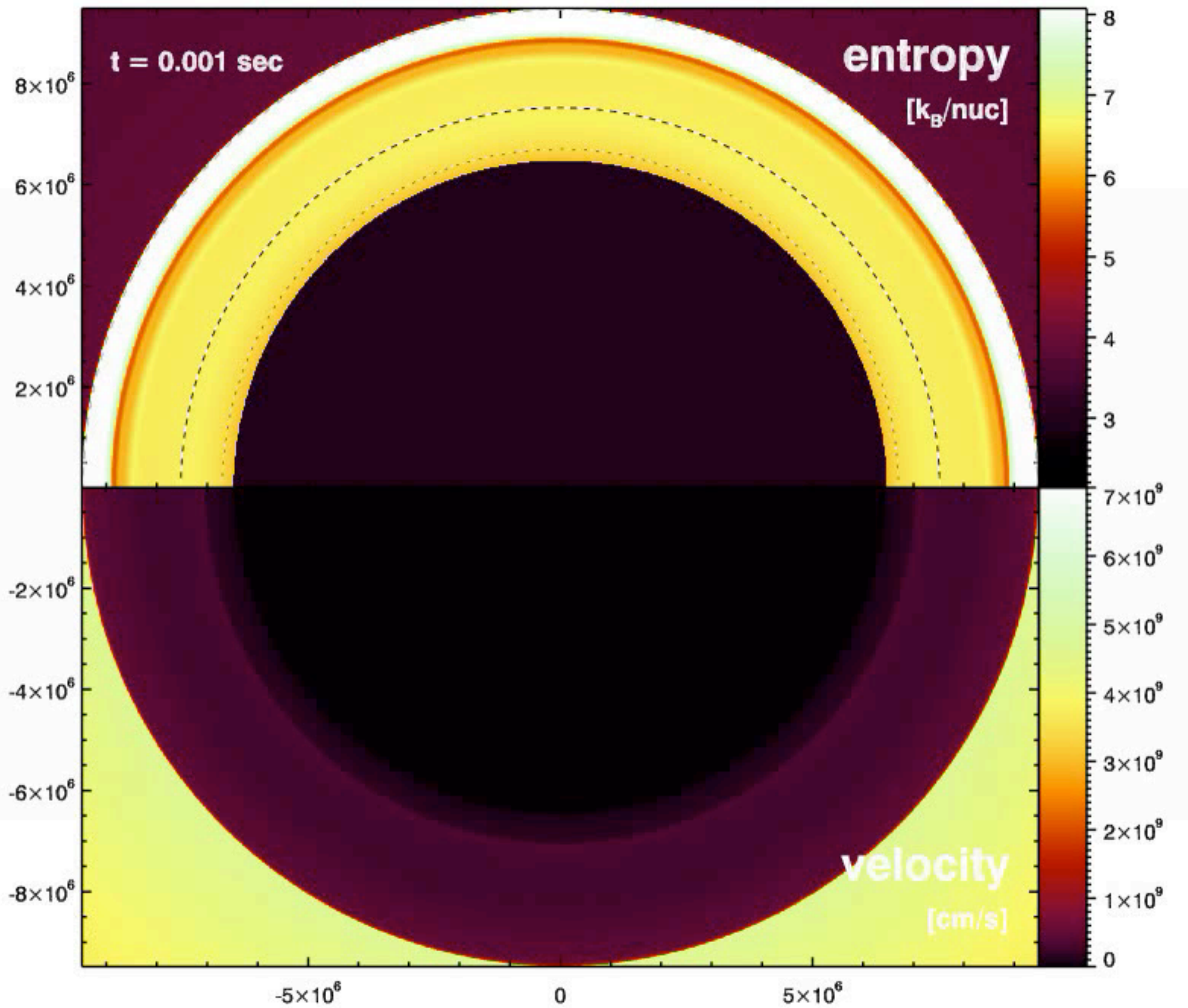
(Meyer et al. 1992, Hoffman et al. 1997, Otsuki et al. 2000, Thompson et al. 2001...)

Core-collapse supernova simulations



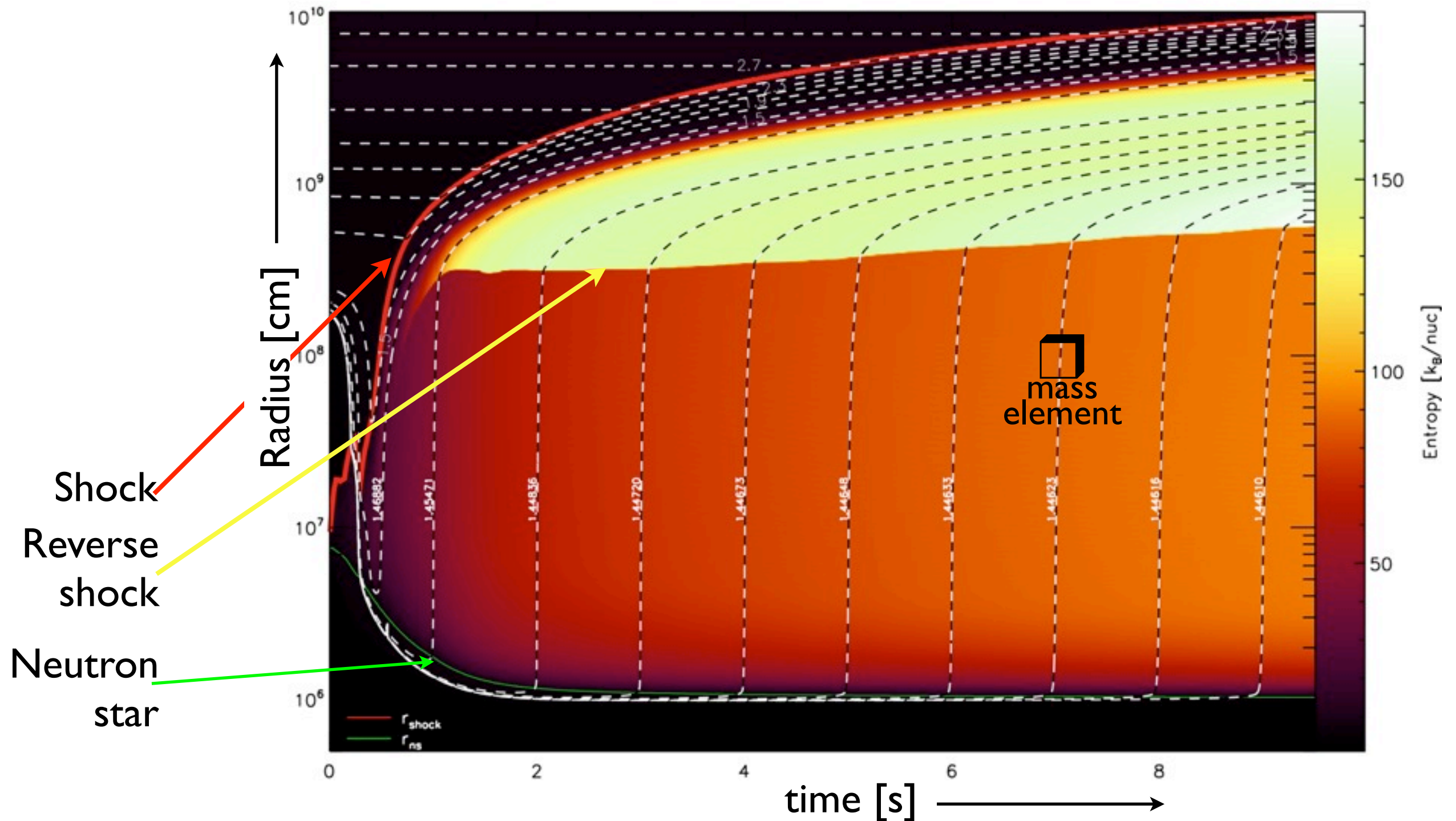
Long-time hydrodynamical simulations:

- ejecta evolution from ~ 5 ms after bounce to ~ 3 s in 2D (Arcones & Janka 2011) and ~ 10 s in 1D (Arcones et al. 2007)
- explosion triggered by neutrinos
- detailed study of nucleosynthesis-relevant conditions

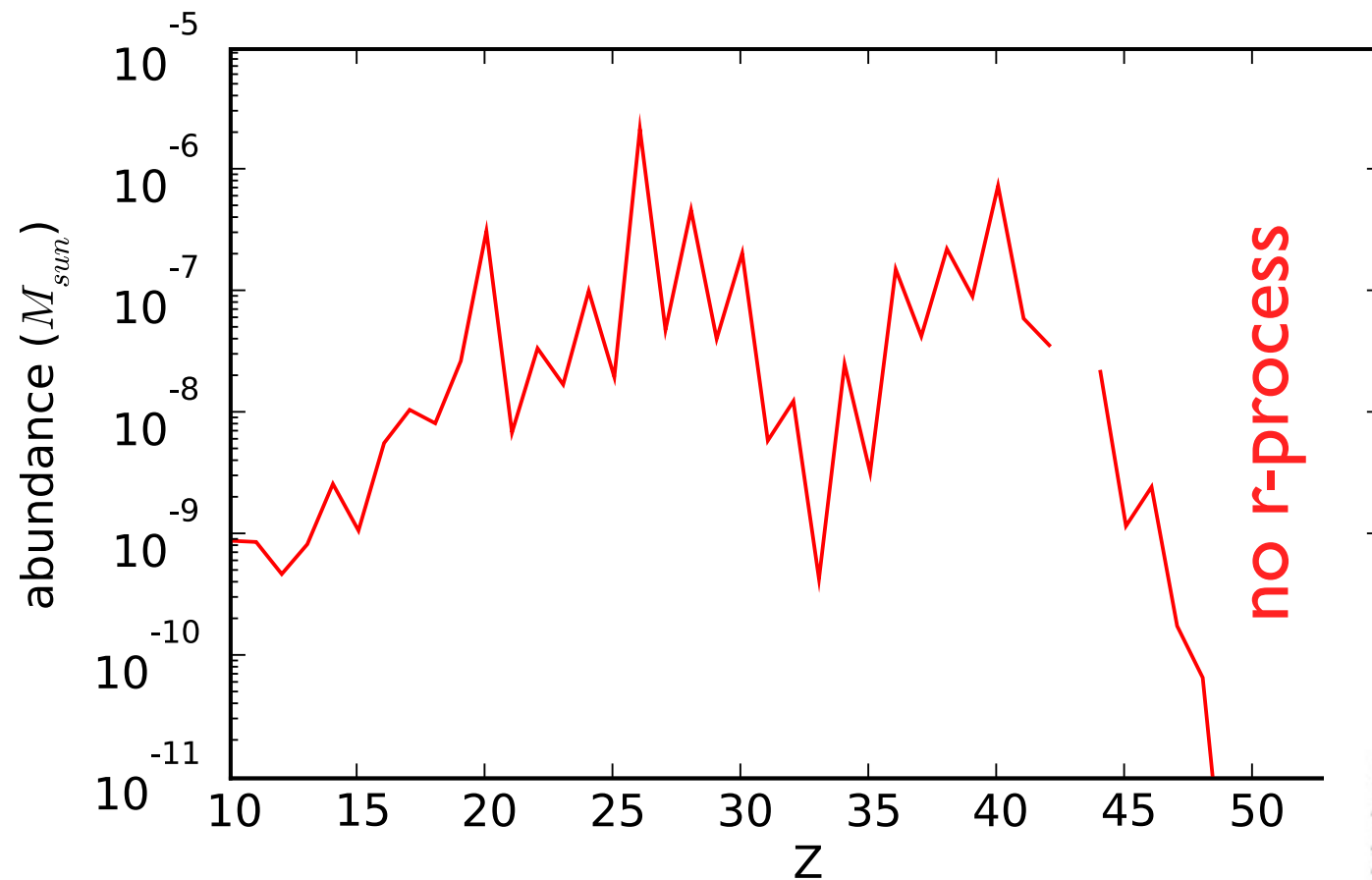


1D simulations for nucleosynthesis studies

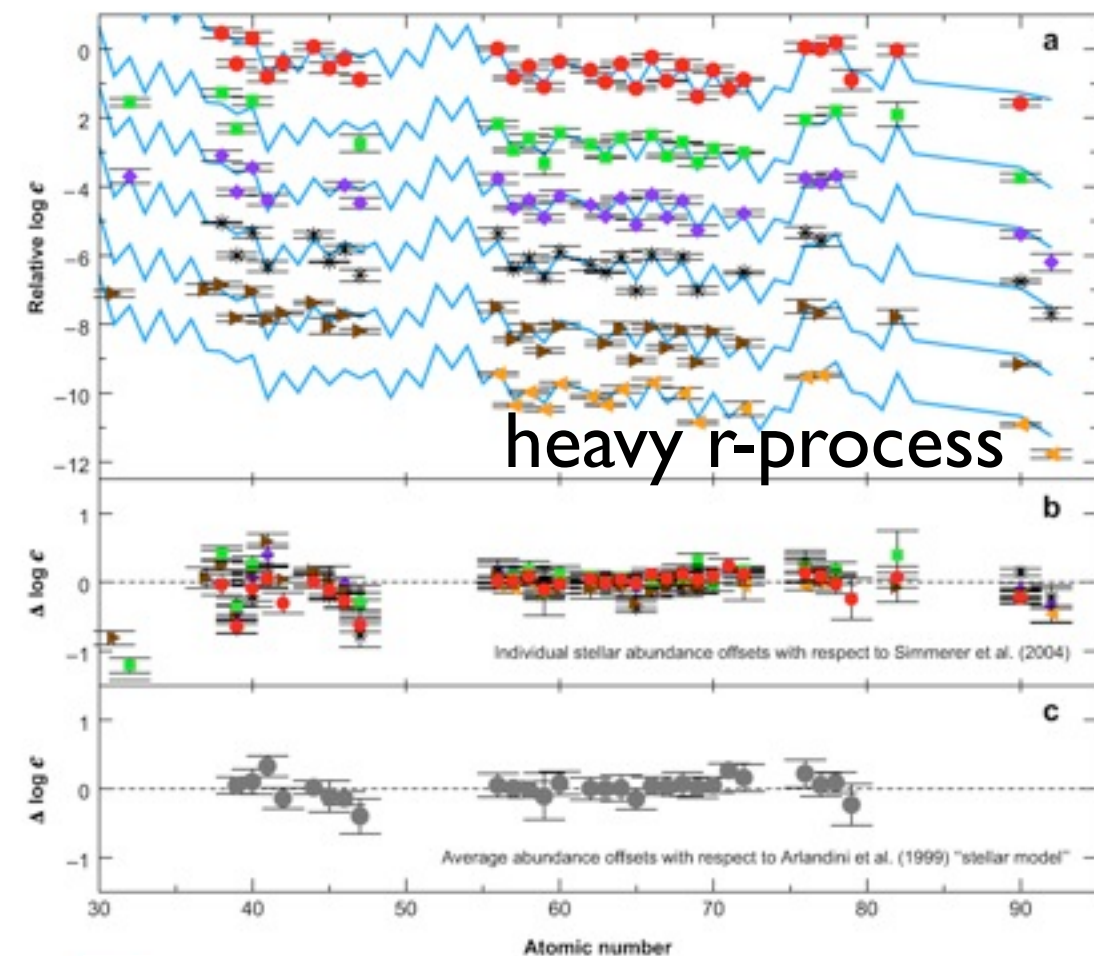
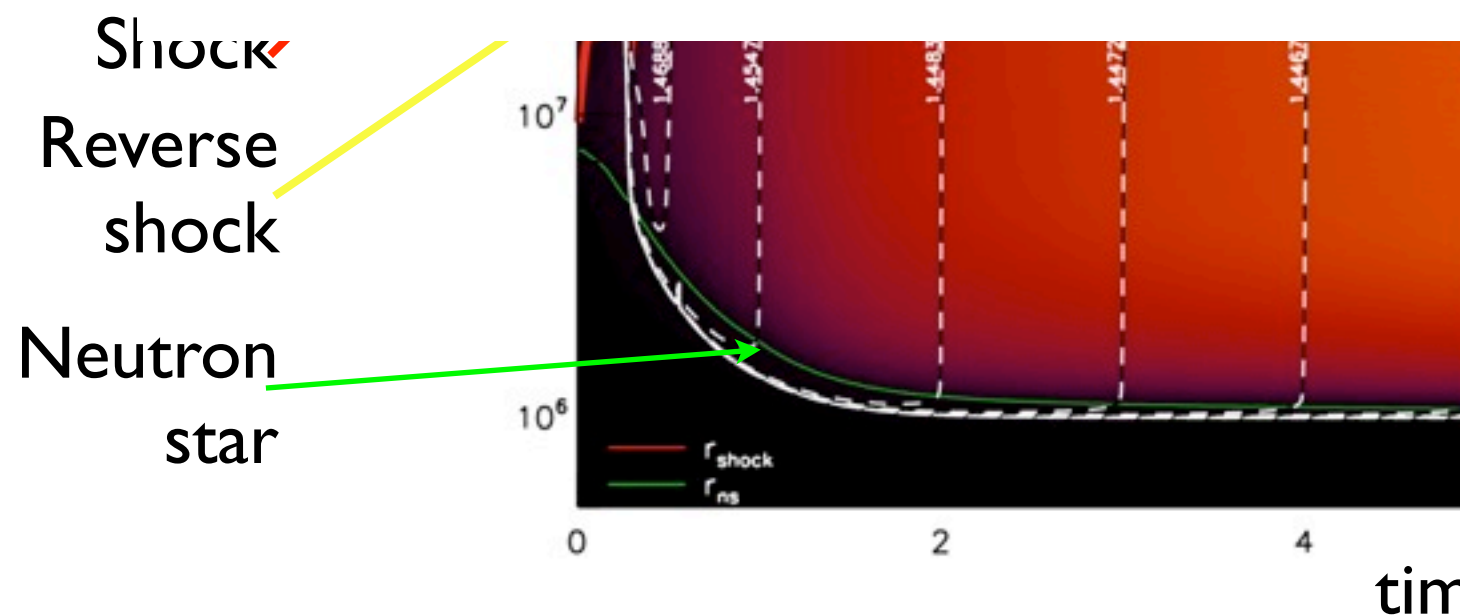
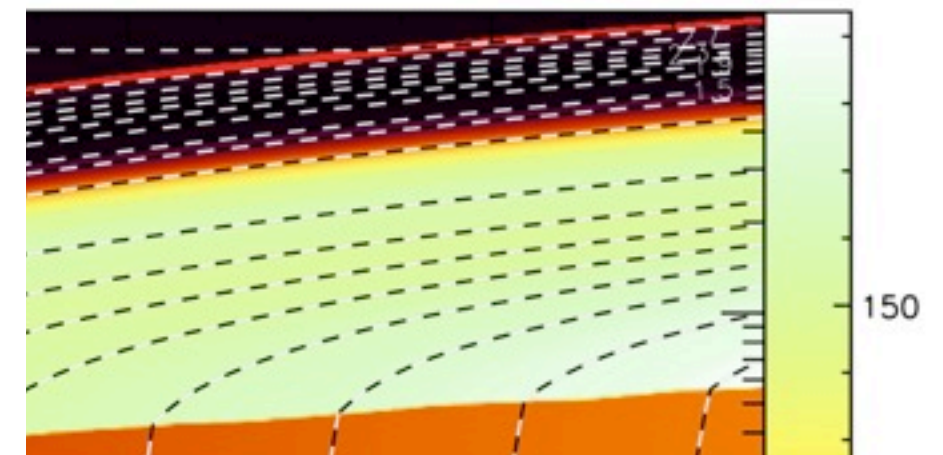
Arcones et al 2007



1D simulations for nucleosynthesis studies



Arcones et al 2007

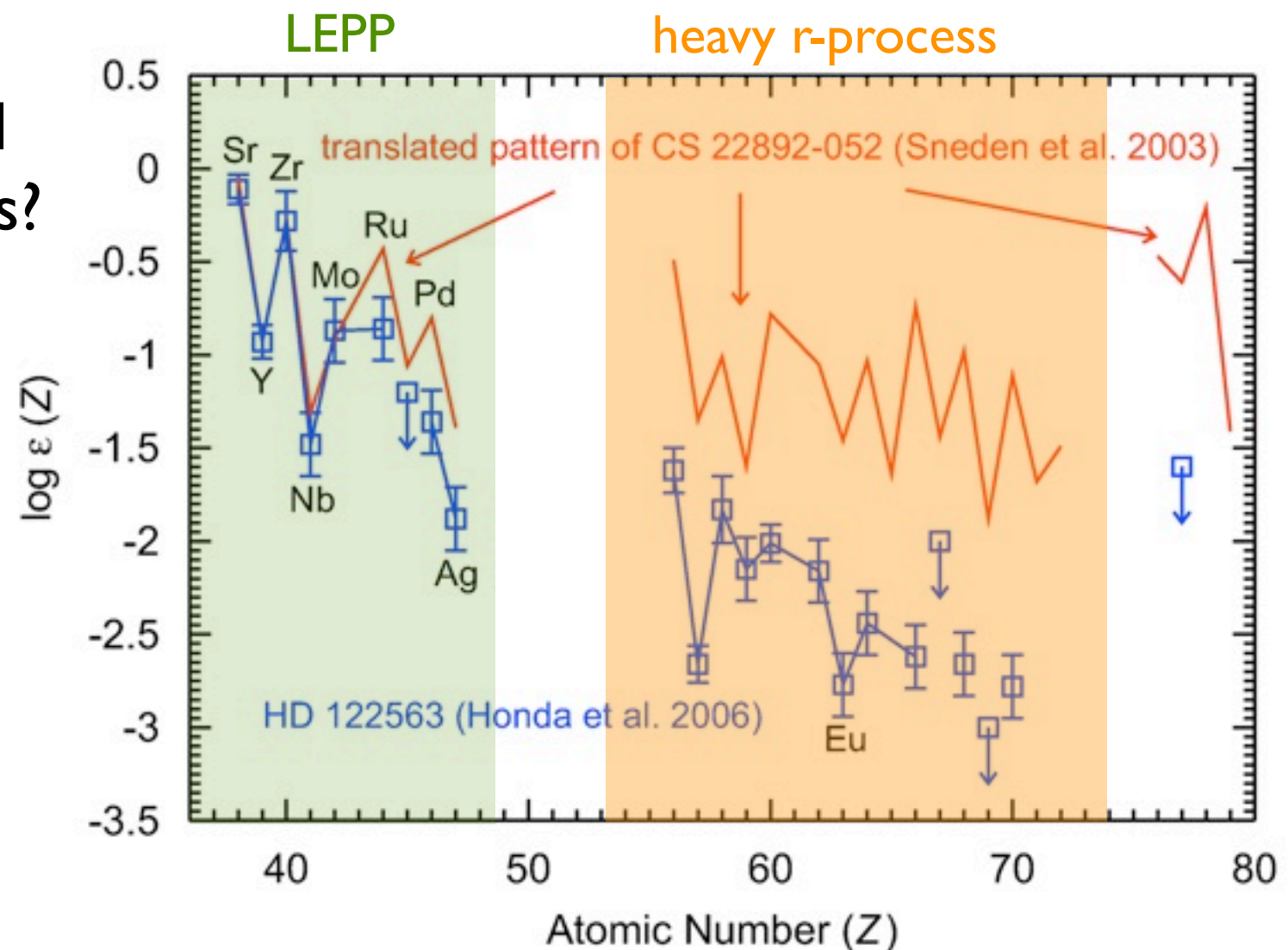


Sneden, Cowan, Gallino 2008

LEPP: Lighter Element Primary Process

Ultra metal-poor stars with **high** and **low** enrichment of heavy r-process nuclei suggest: two components or sites (Qian & Wasserburg, 2001, Travaglio et al. 2004)

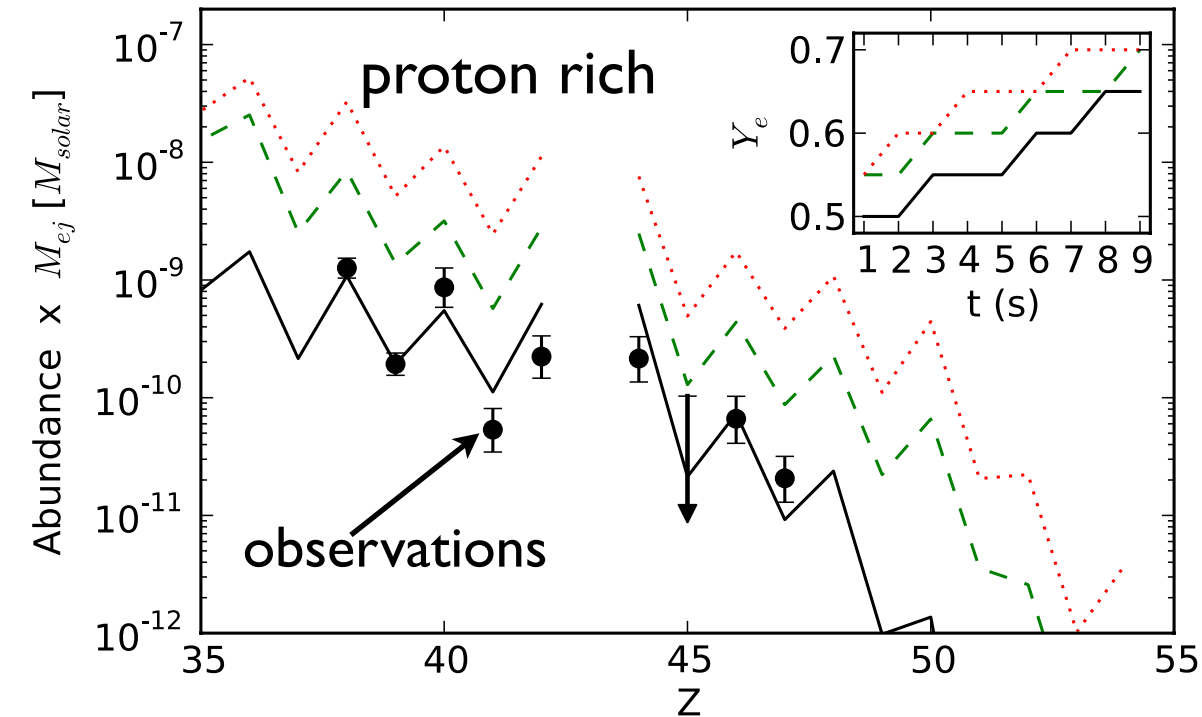
Can the LEPP pattern be produced in neutrino-driven wind simulations?



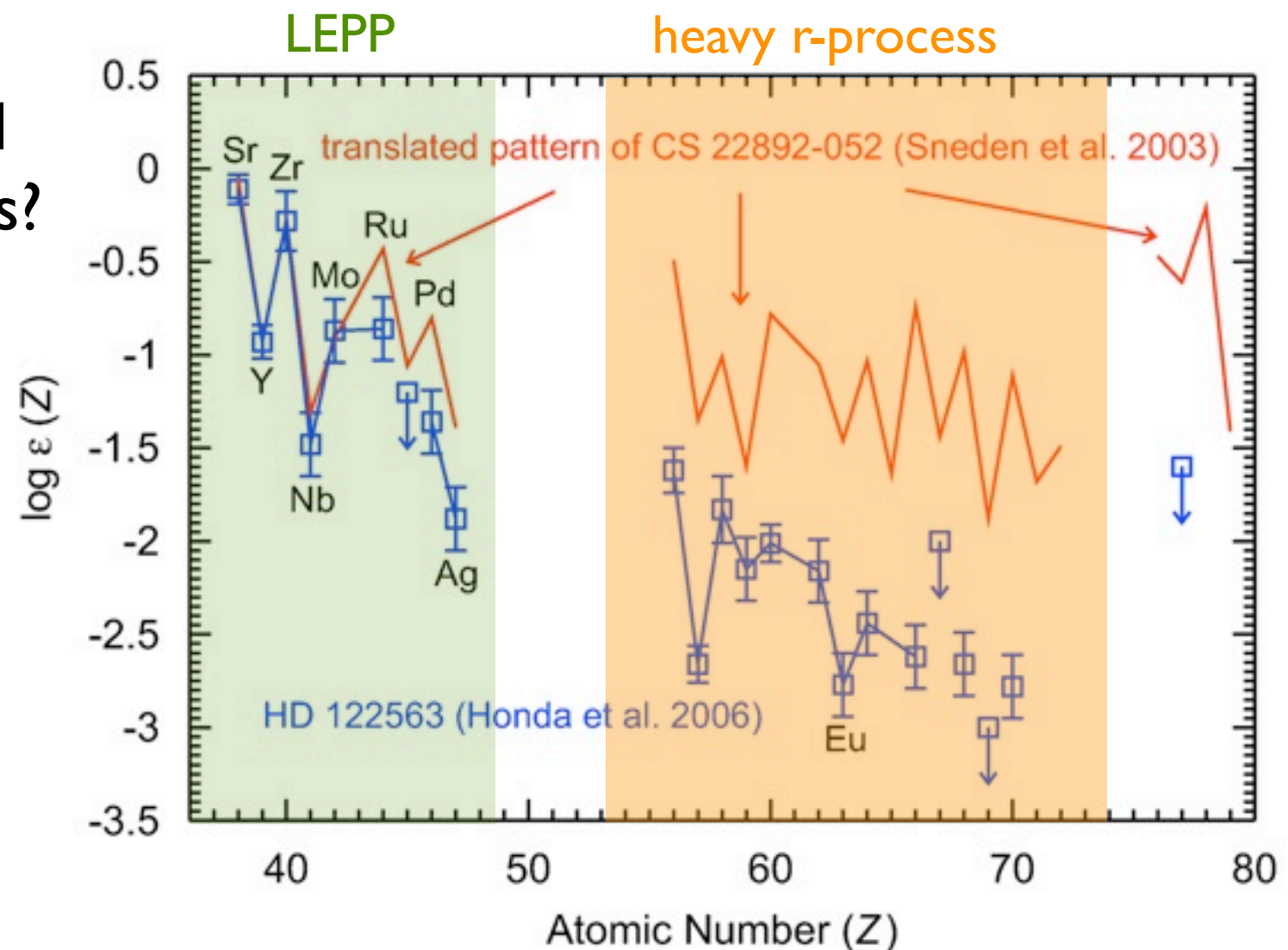
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Arcones & Montes, 2011

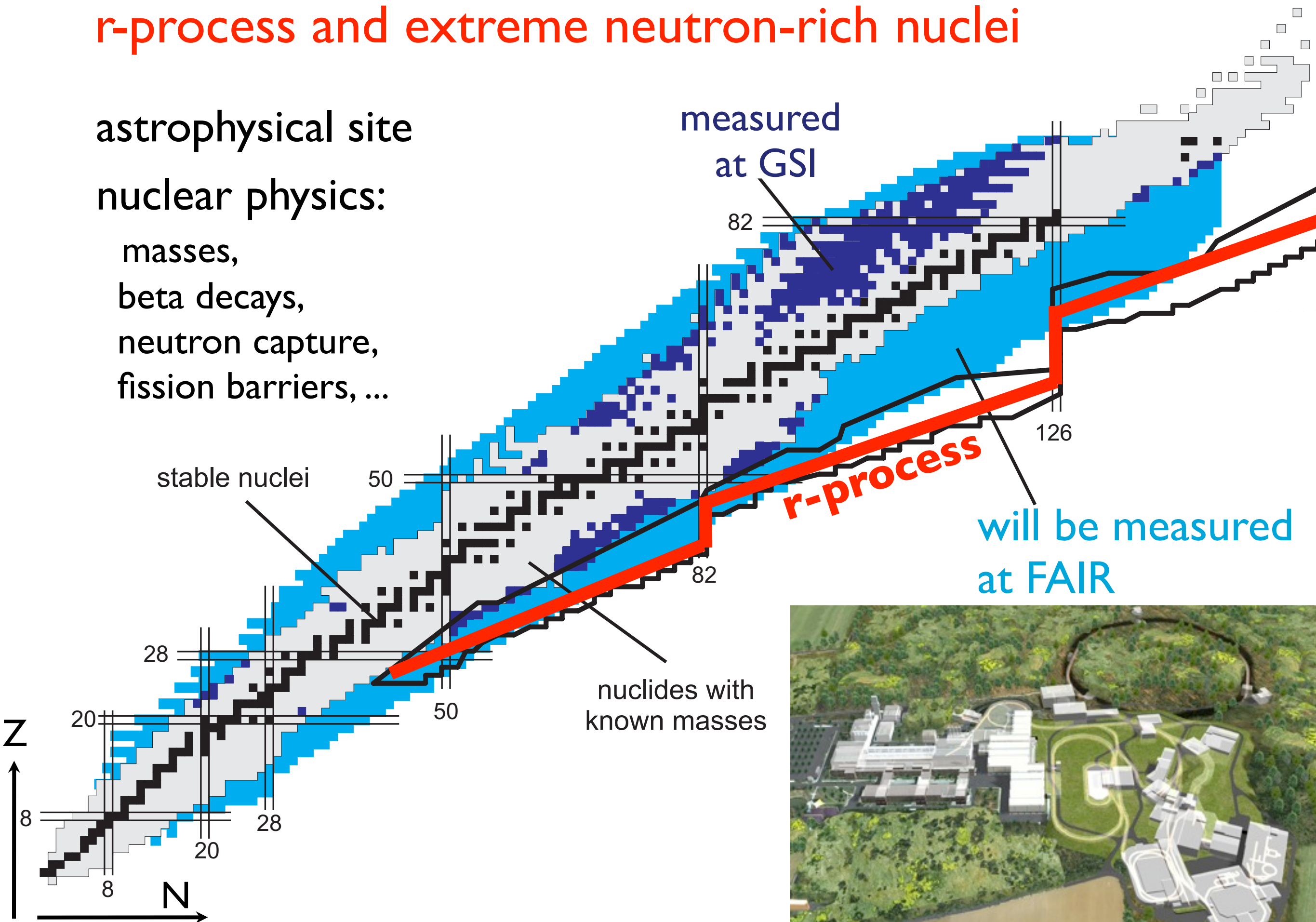


r-process and extreme neutron-rich nuclei

astrophysical site

nuclear physics:

masses,
beta decays,
neutron capture,
fission barriers, ...

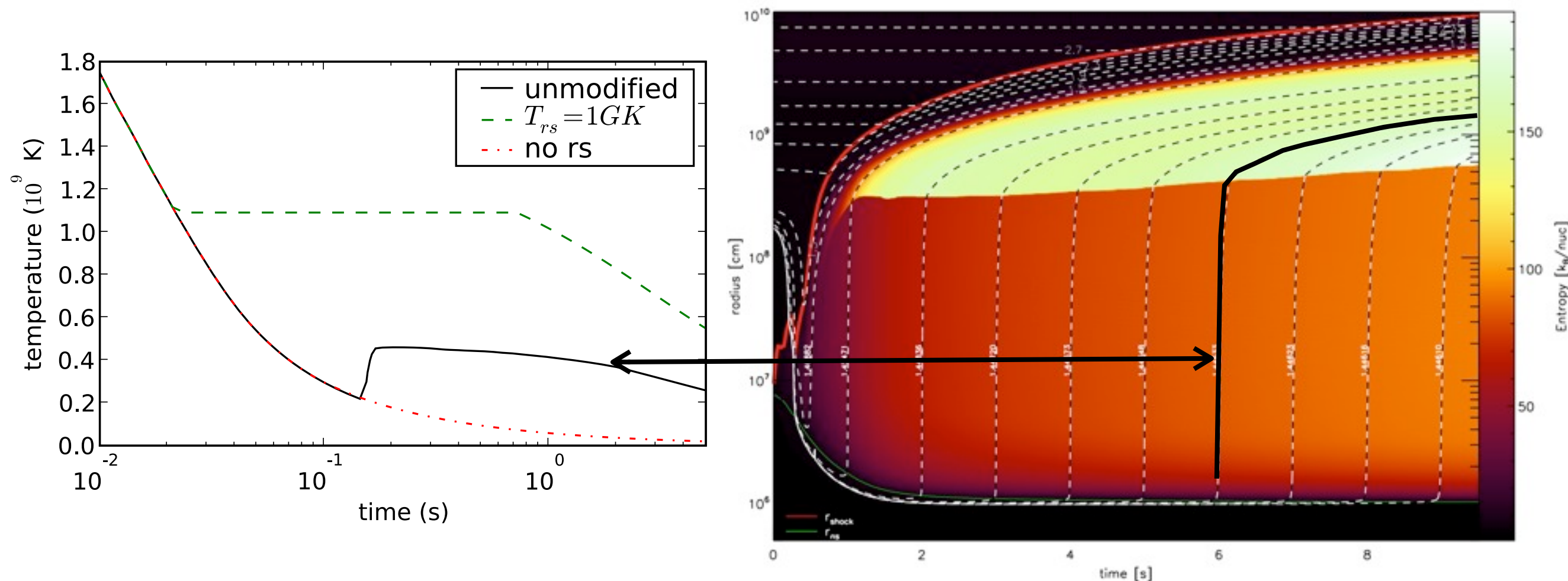


r-process: long-time evolution and reverse shock

We use one trajectory from our hydrodynamical simulations with entropy increased by factor two.

Vary the long-time evolution:

- reverse shock at 1 GK
- no reverse shock

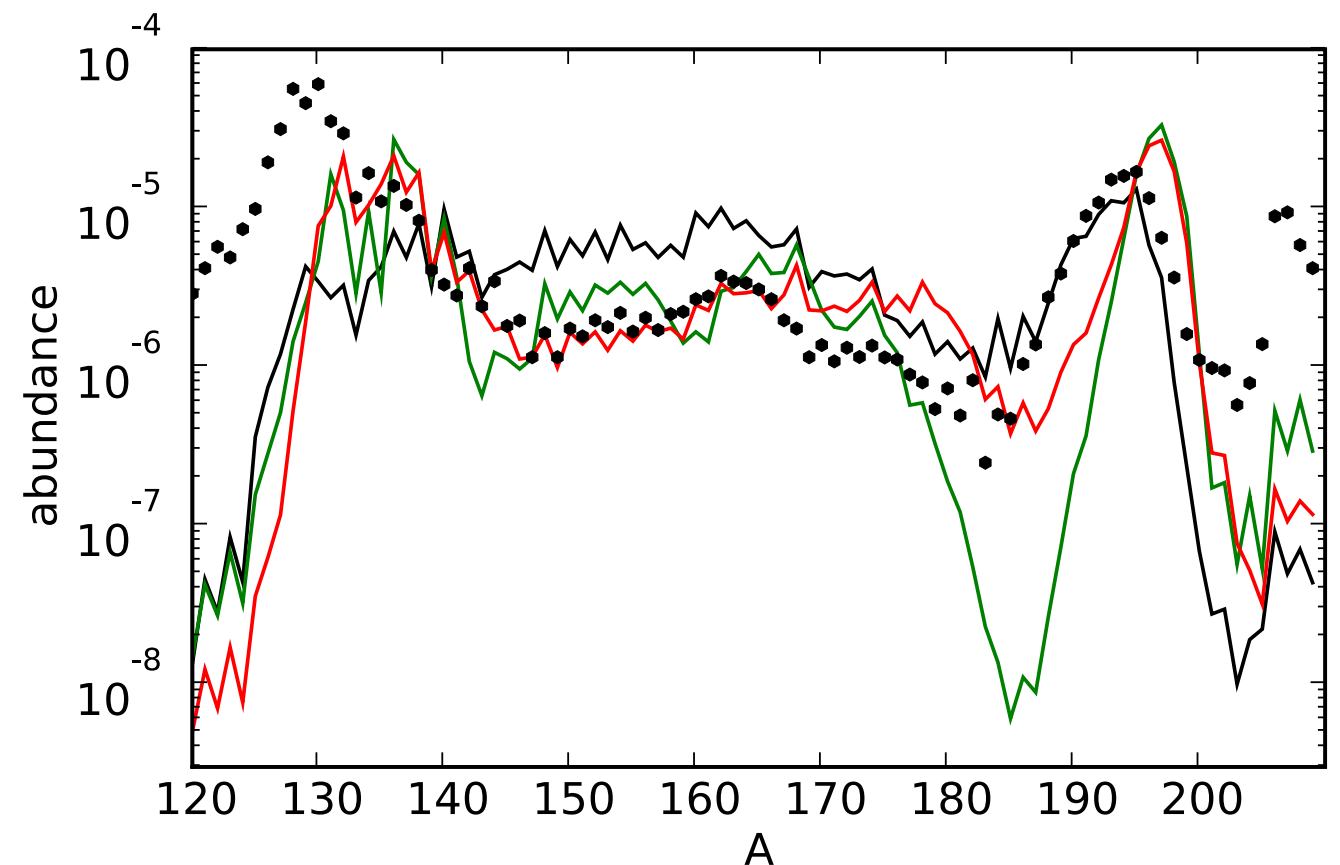
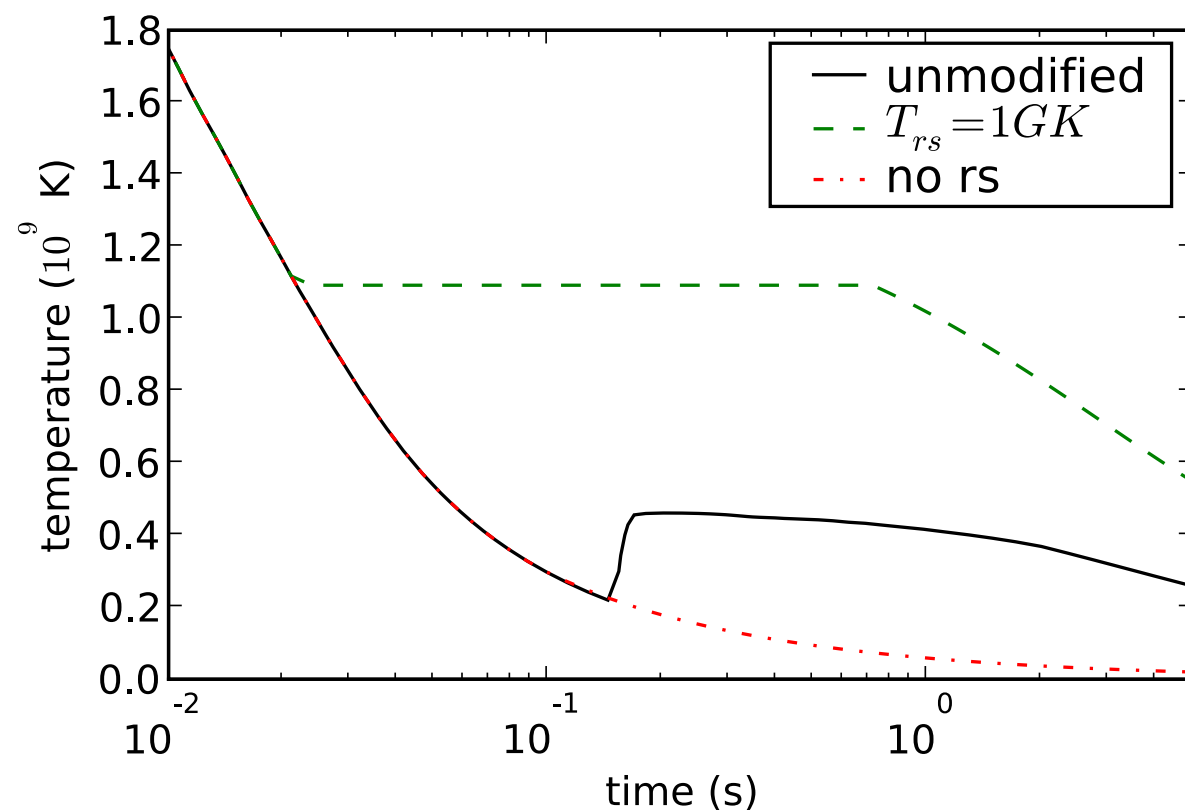


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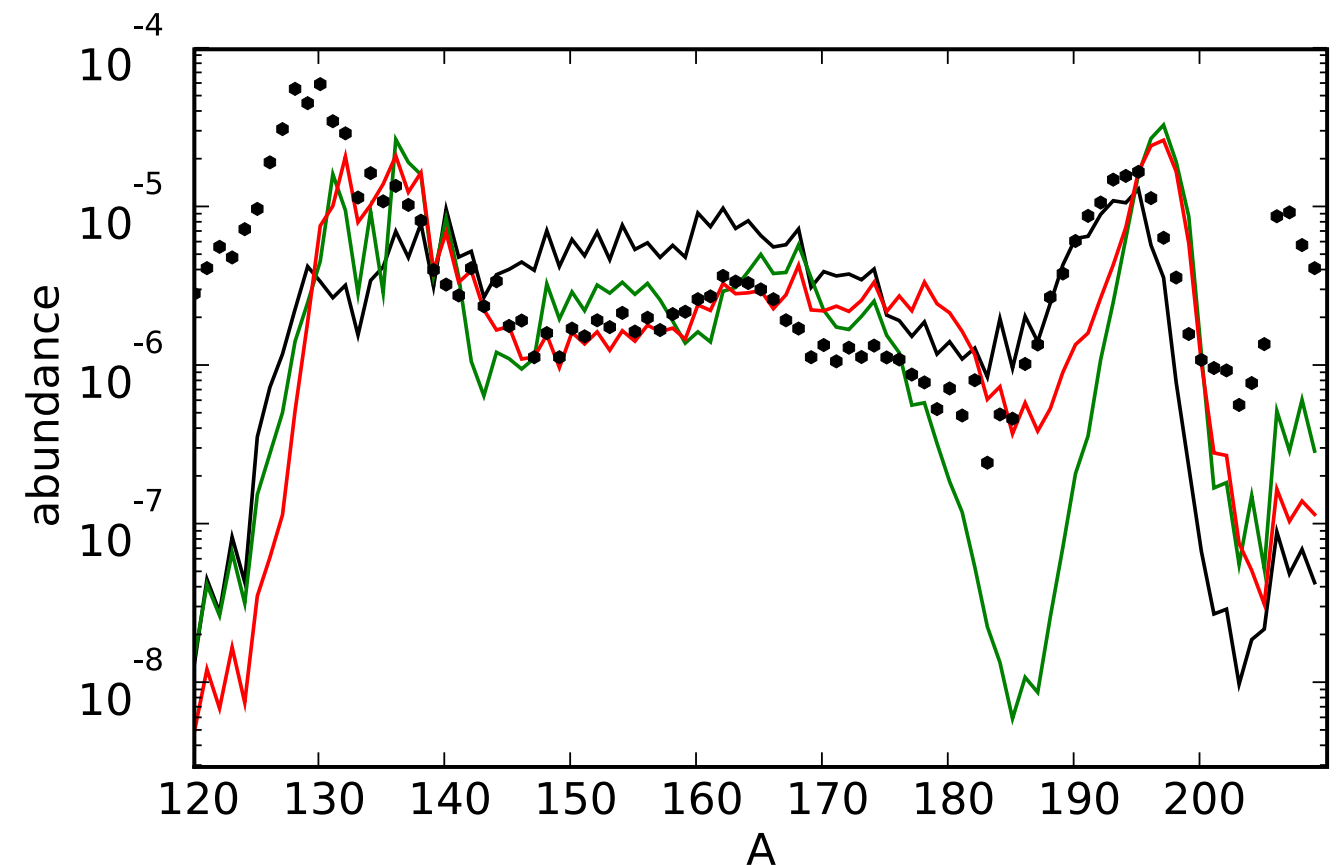
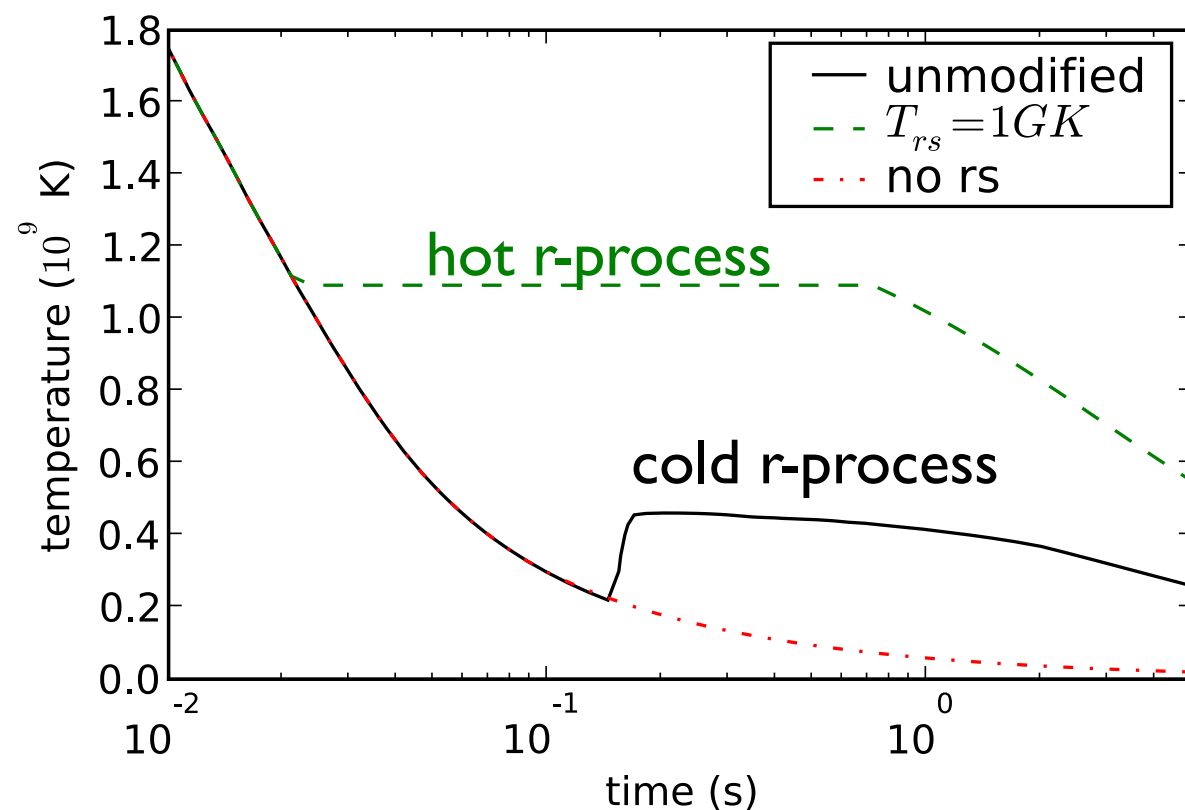


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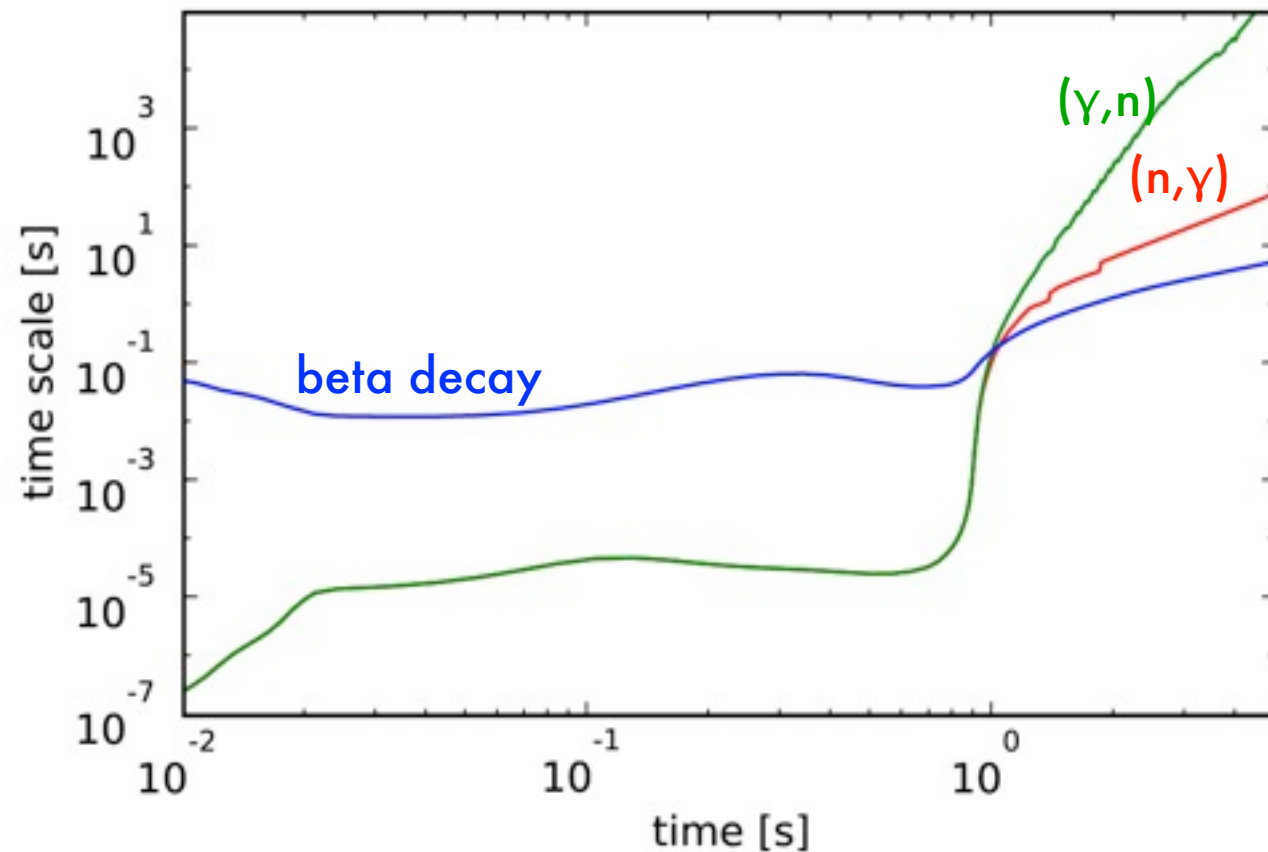
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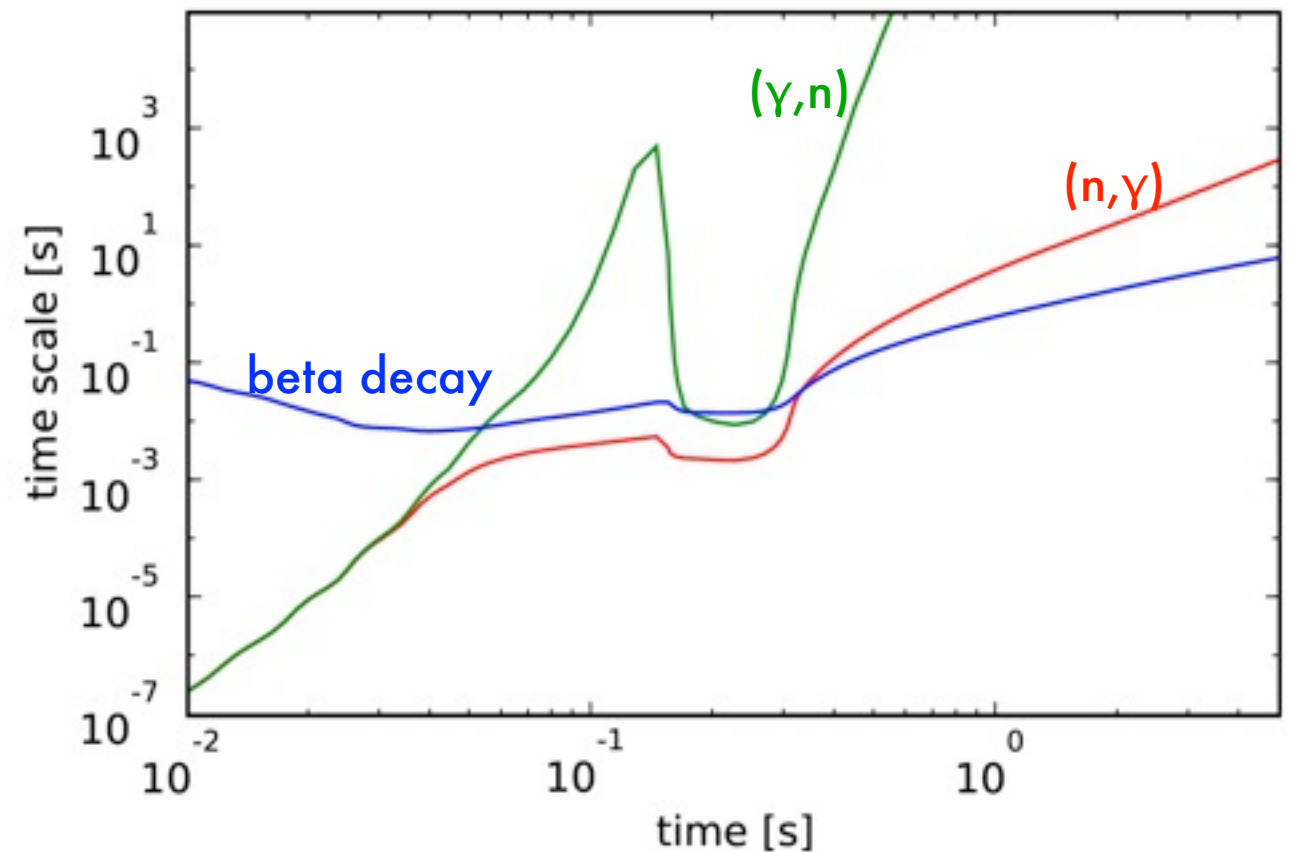
Long-time evolution: high vs. low temperature

Hot r-process



The evolution takes place under
 (n,γ) - (γ,n) equilibrium
(classical r-process, Seeger, Fowler and
Clayton 1965, Kratz et al. 1993).

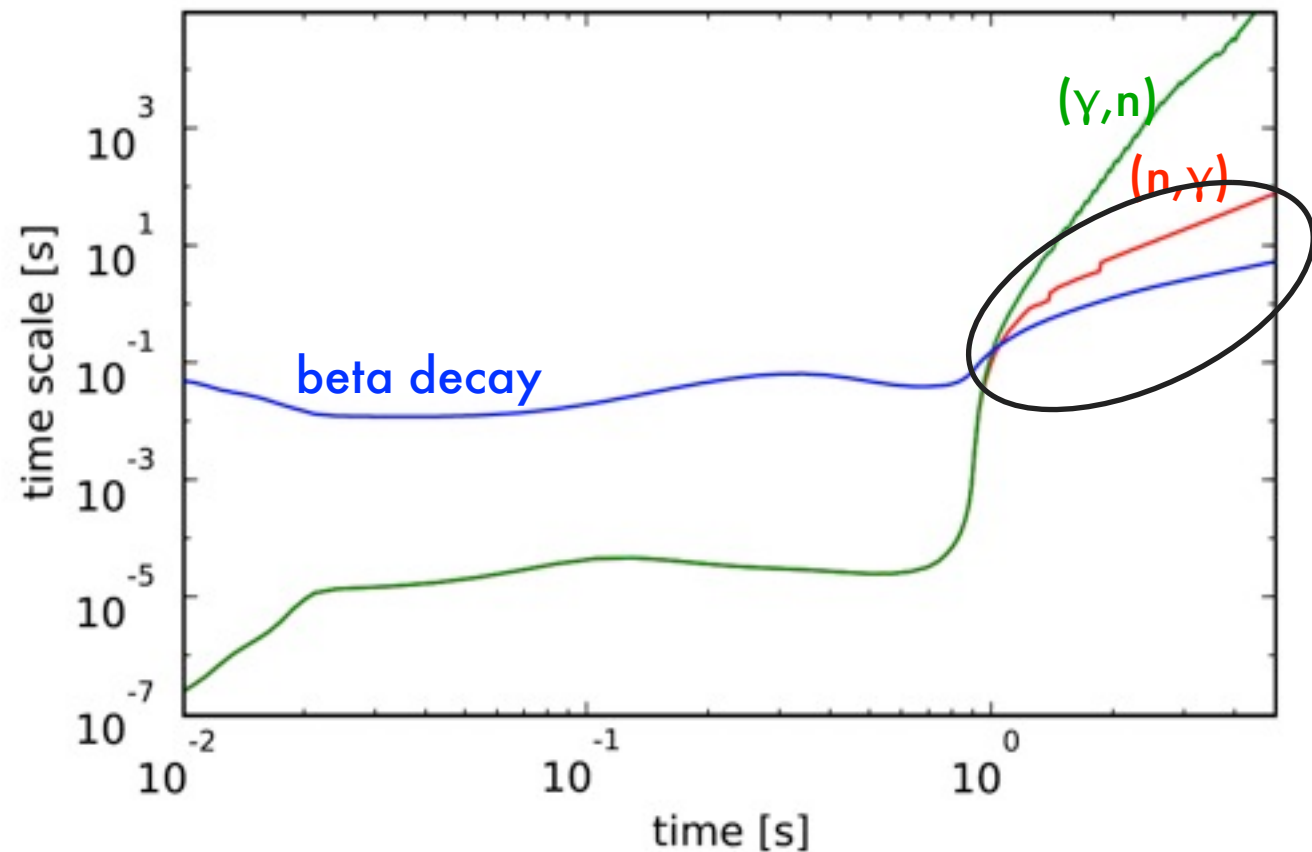
Cold r-process



Competition between beta decay and
neutron capture (Blake & Schramm 1976,
Wanajo 2007, Janka & Panov 2009)

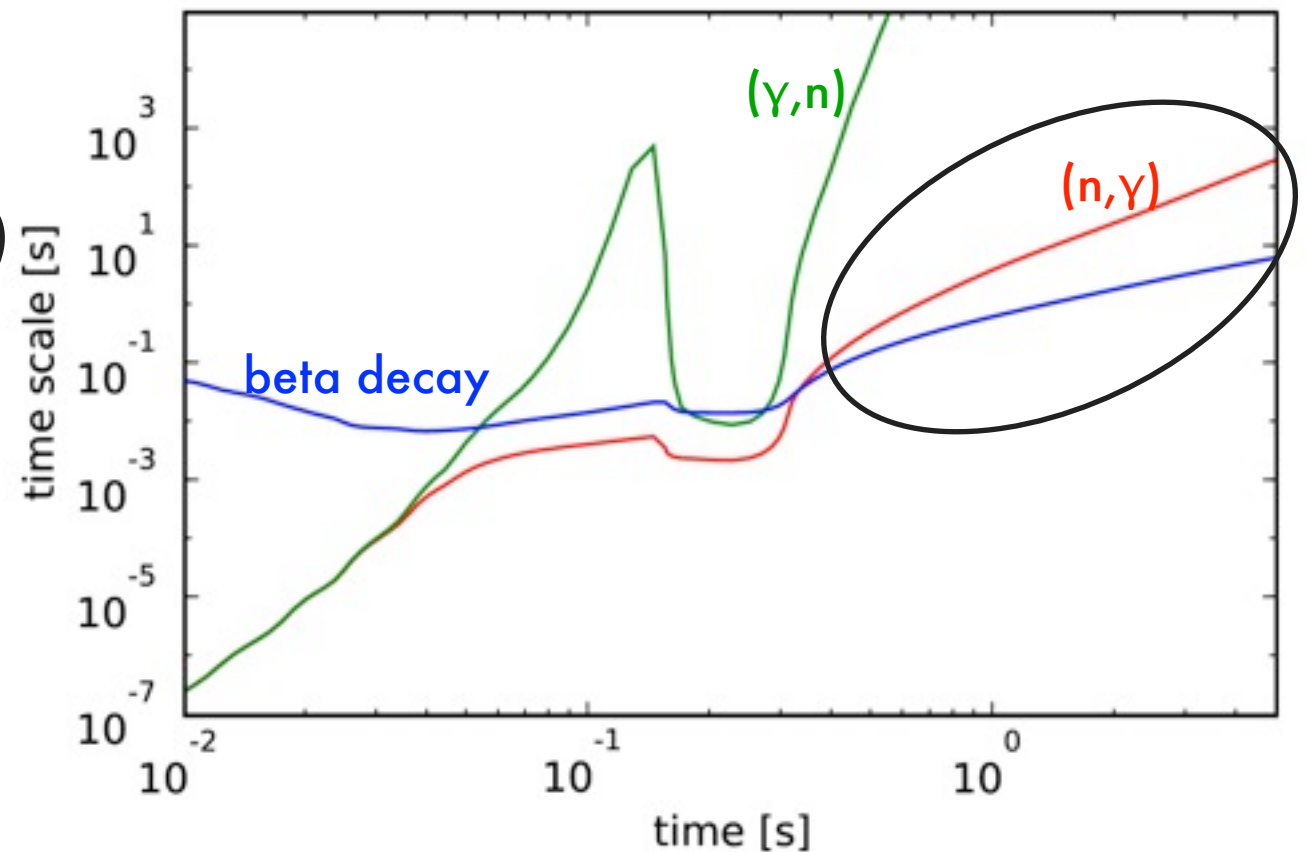
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Cold r-process



Competition between beta decay and
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Final abundances are strongly affected by neutron captures and beta decays that compete when matter moves back to stability.

Sensitivity to mass models

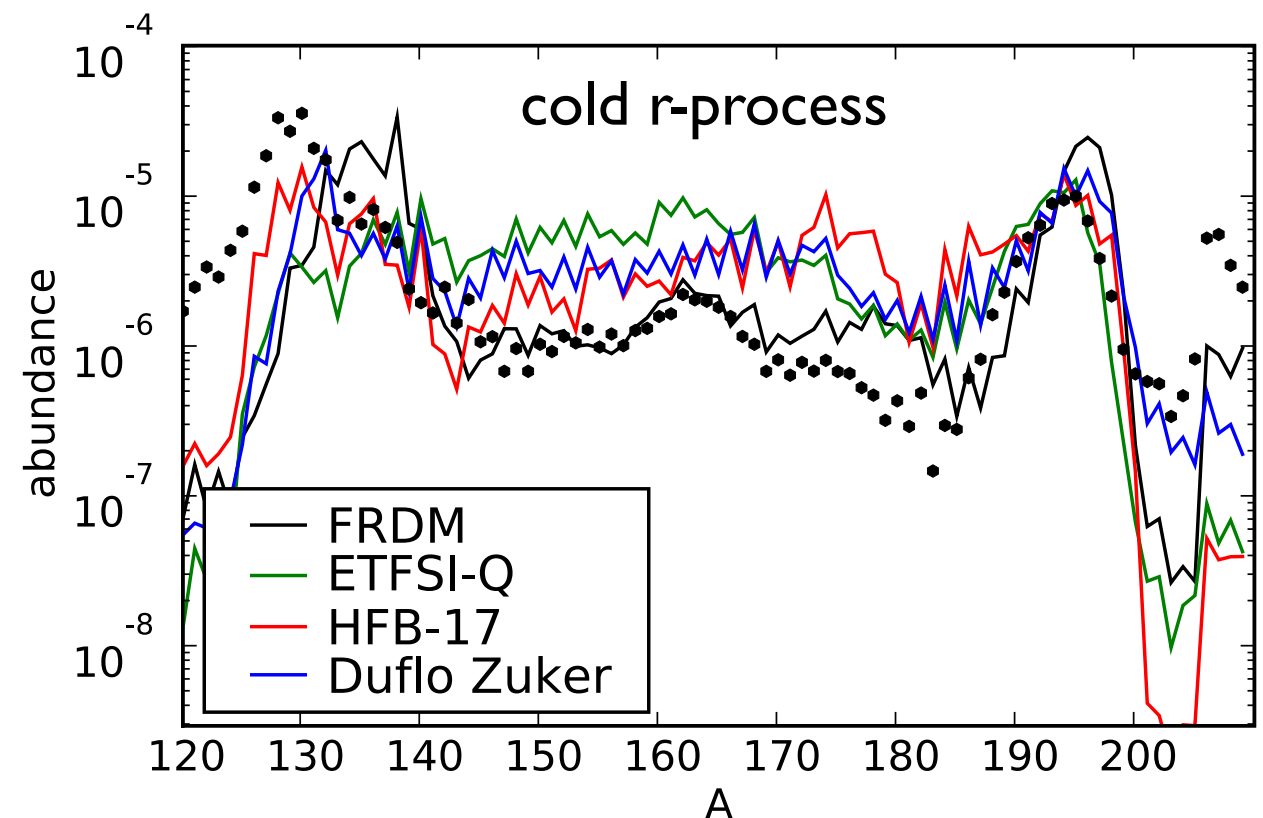
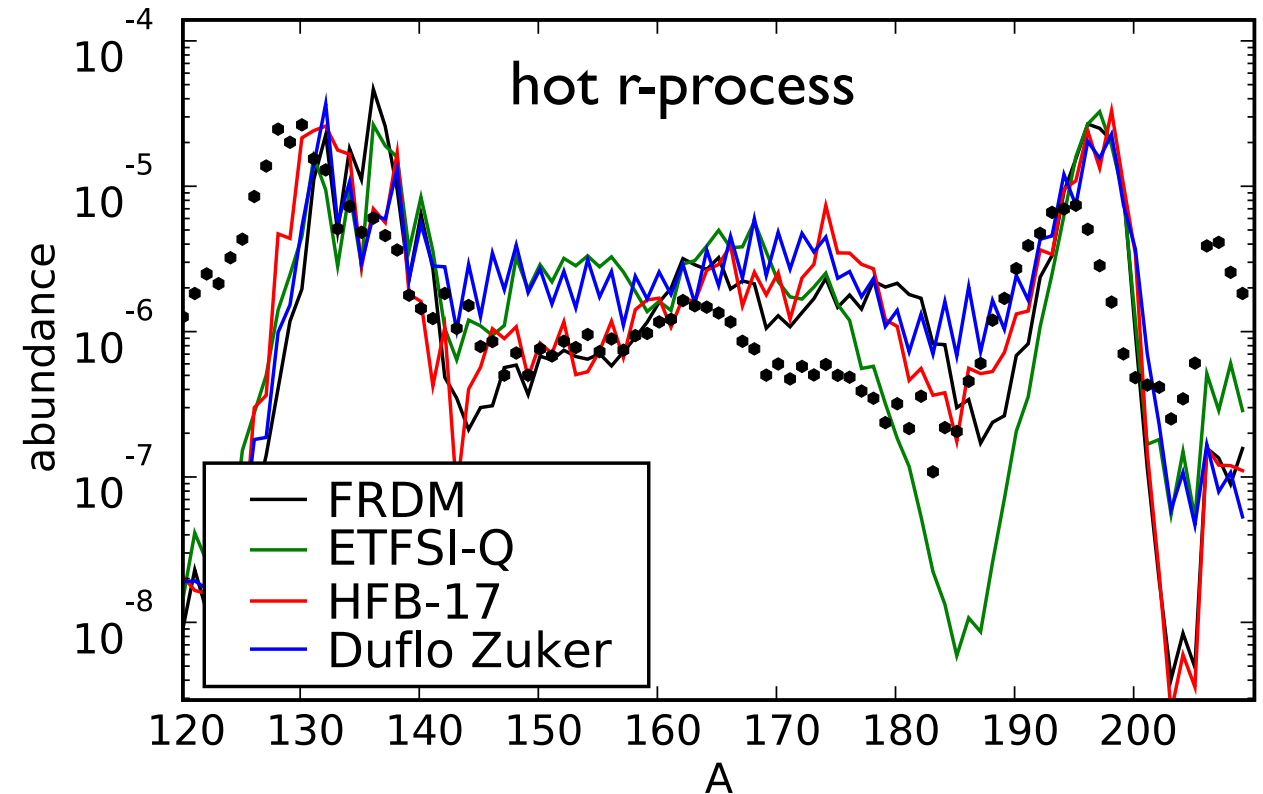
Compare four different mass models:

- FRDM (Möller et al. 1995)
- ETFSI-Q (Pearson et al. 1996)
- HFB-17 (Goriely et al. 2009)
- Duflo&Zuker mass formula

two cases: (n,γ) - (γ,n) equilibrium and non-equilibrium.

The nuclear physics input affects the final abundances differently depending on the long-time dynamical evolution.

Can we link the behavior of the masses (neutron separation energies) to the final r-process abundances?



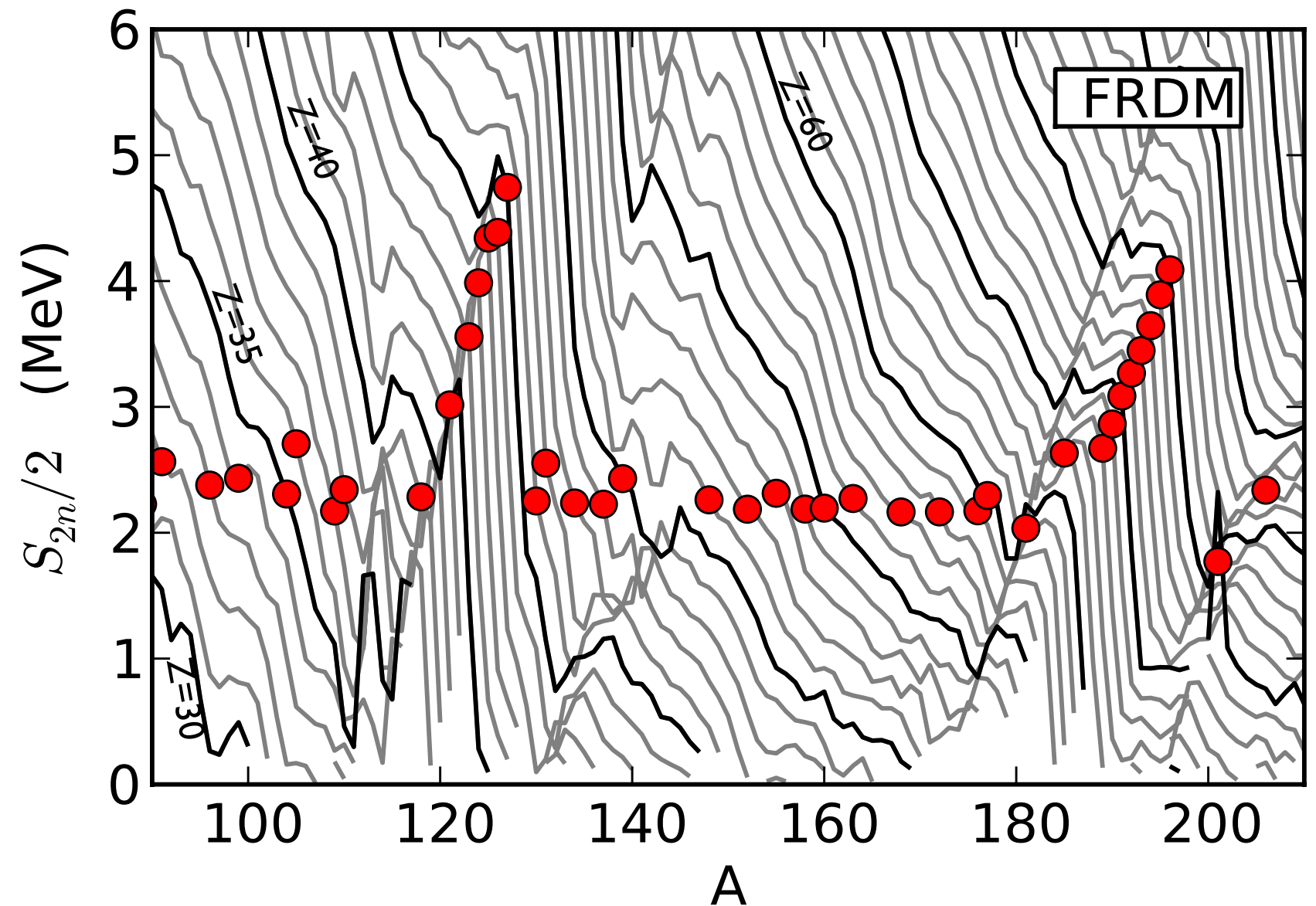
Two neutron separation energy

Abundances



S_{2n}

Nuclear
properties



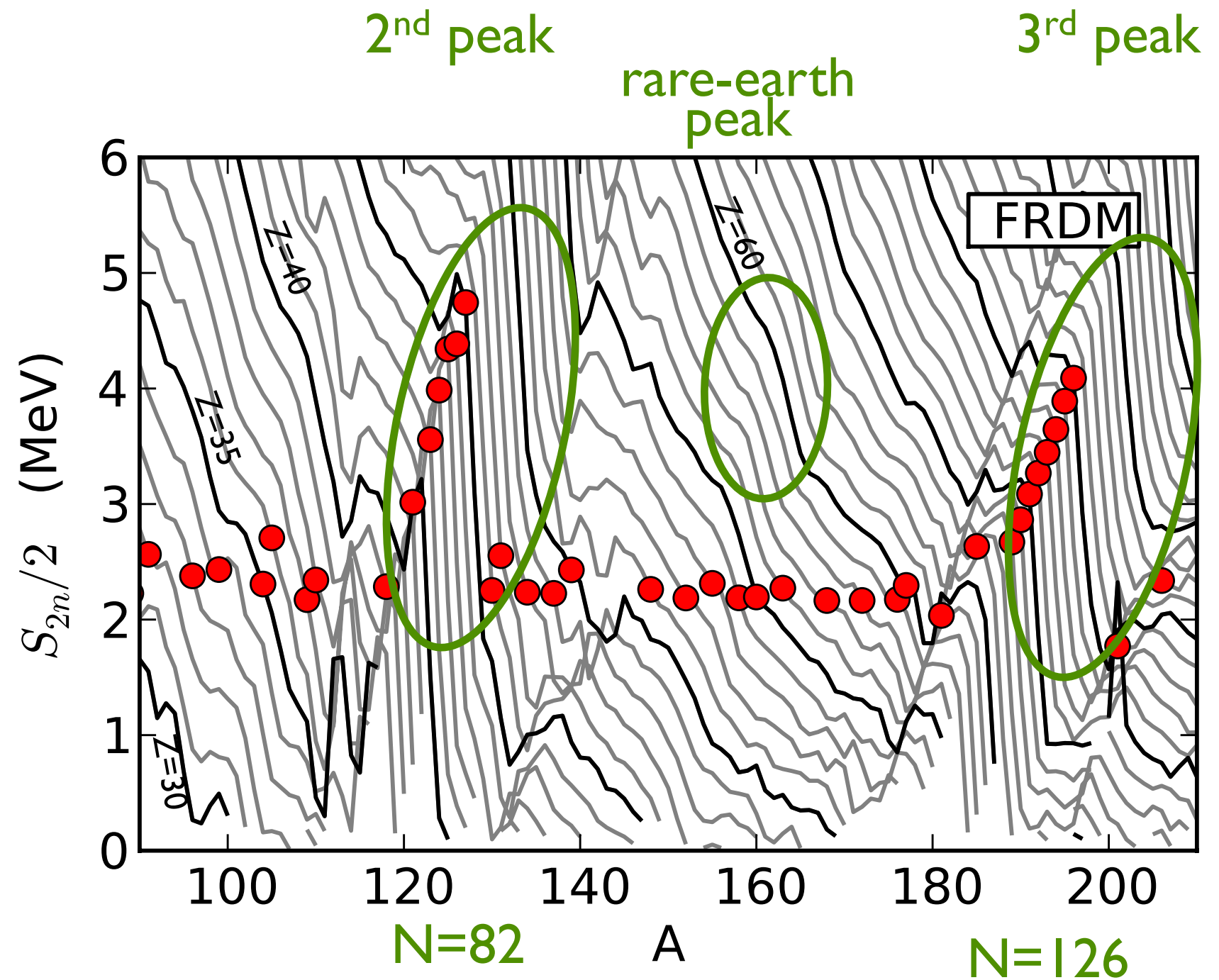
Two neutron separation energy

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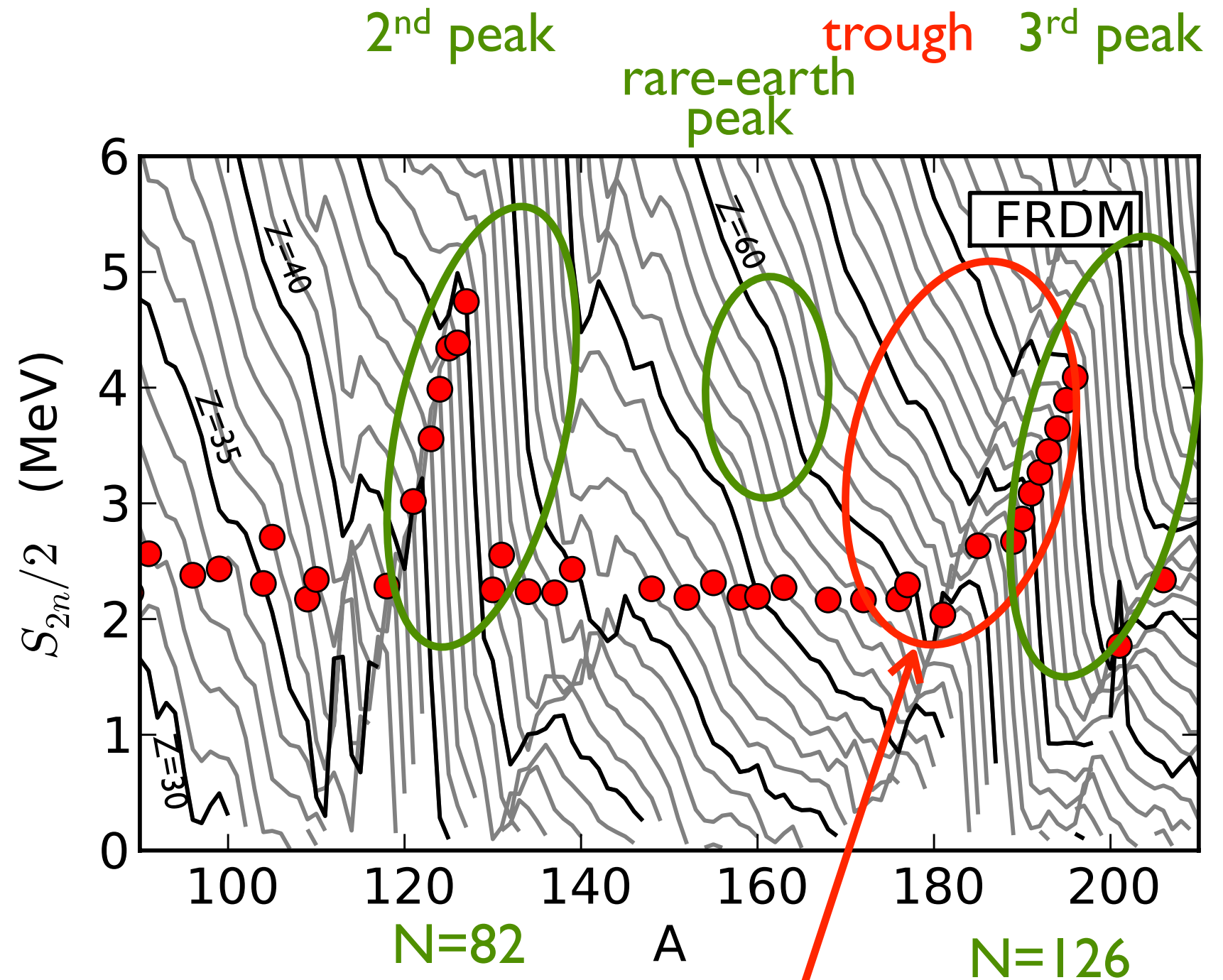
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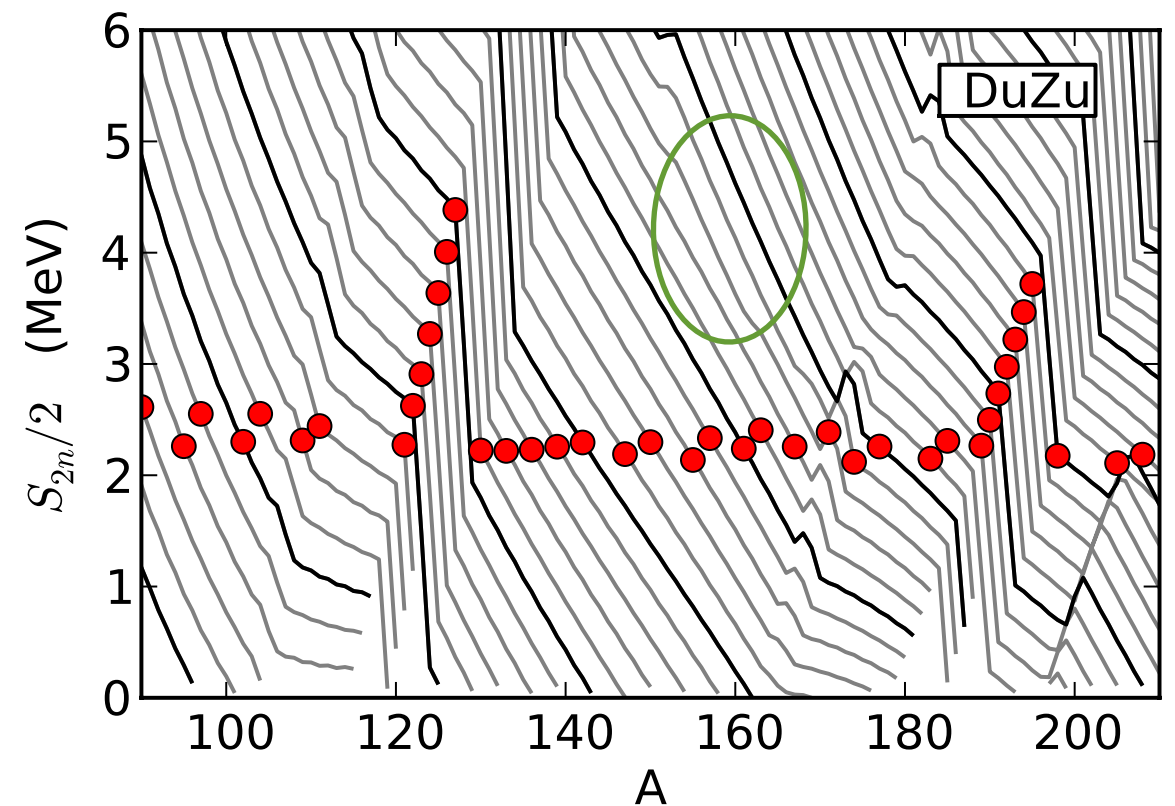
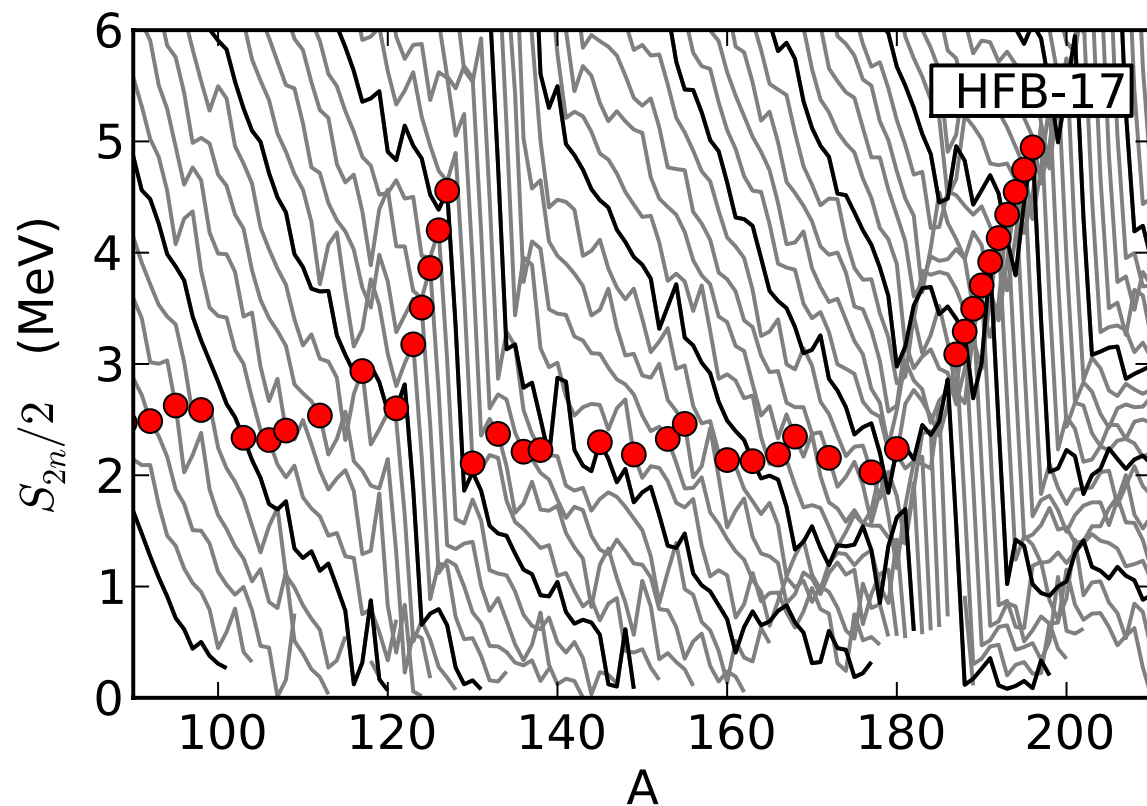
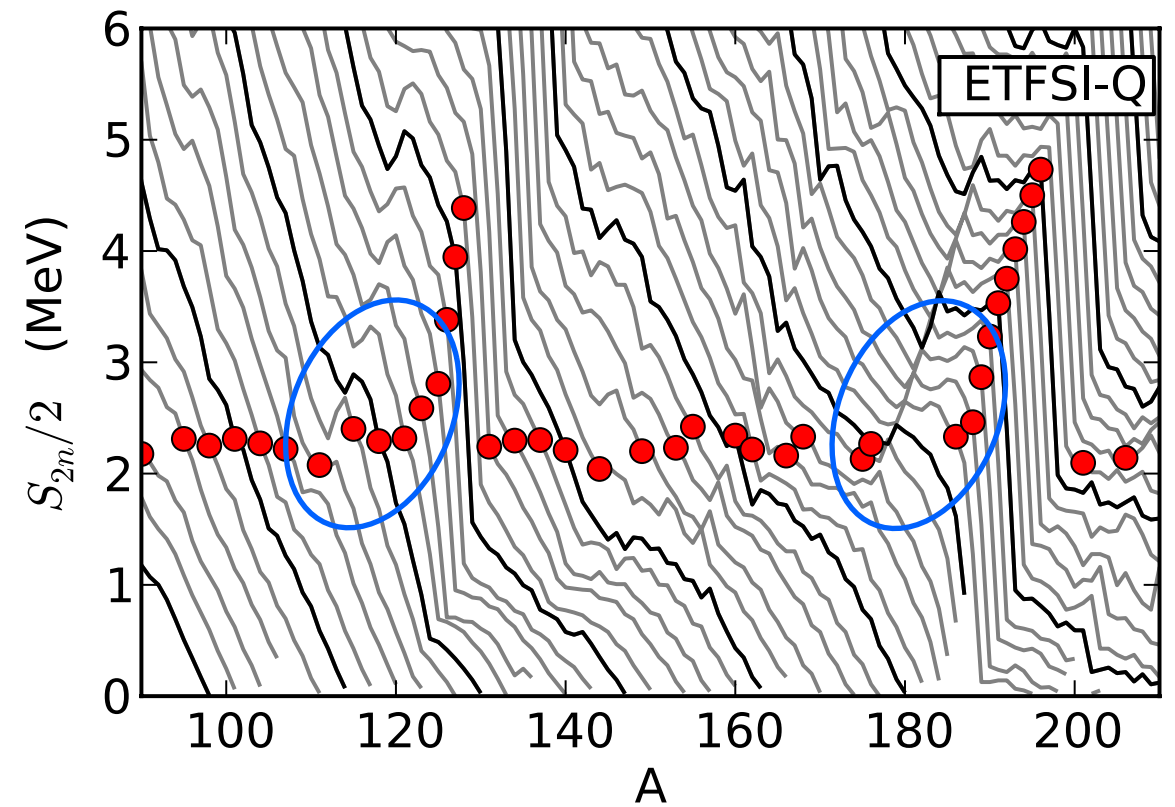
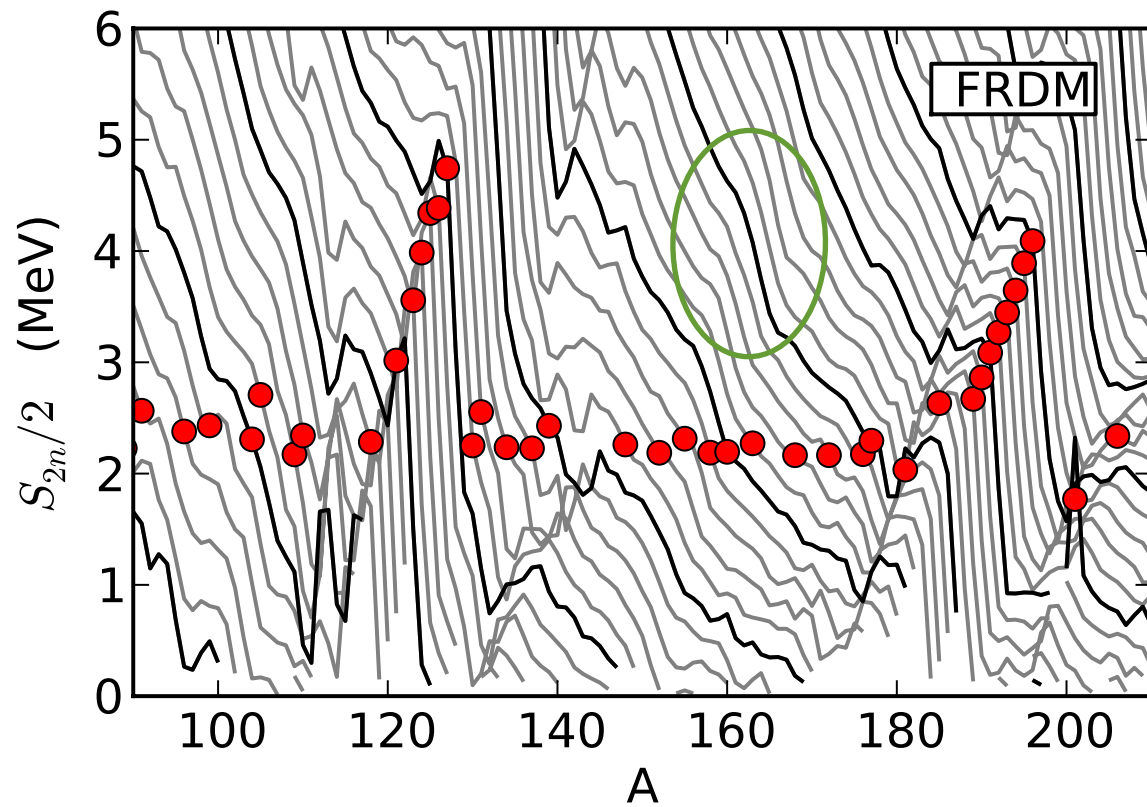
S_{2n}

Nuclear
properties



transition from
deformed to spherical

Aspects of different mass models



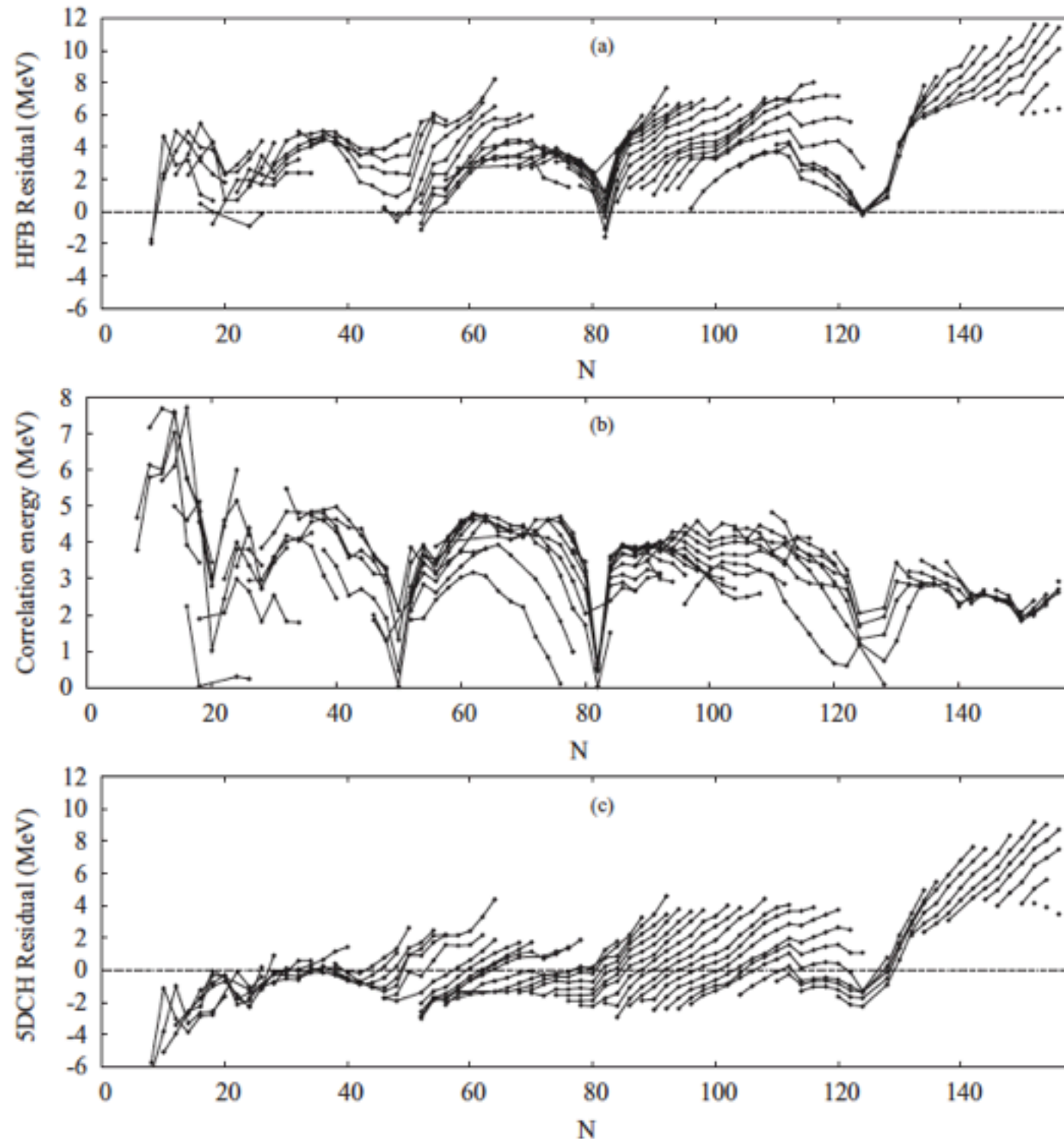
Structure of even-even nuclei using a mapped collective Hamiltonian and the D1S Gogny interaction

J.-P. Delaroche,^{1,*} M. Girod,¹ J. Libert,² H. Goutte,¹ S. Hilaire,¹ S. Péru,¹ N. Pillet,¹ and G. F. Bertsch^{3,*}

¹CEA, DAM, DIF, F-91297 Arpajon, France

²Institut de Physique Nucléaire IN2P3-CNRS/Université Paris-Sud, 91406 Orsay Cedex, France

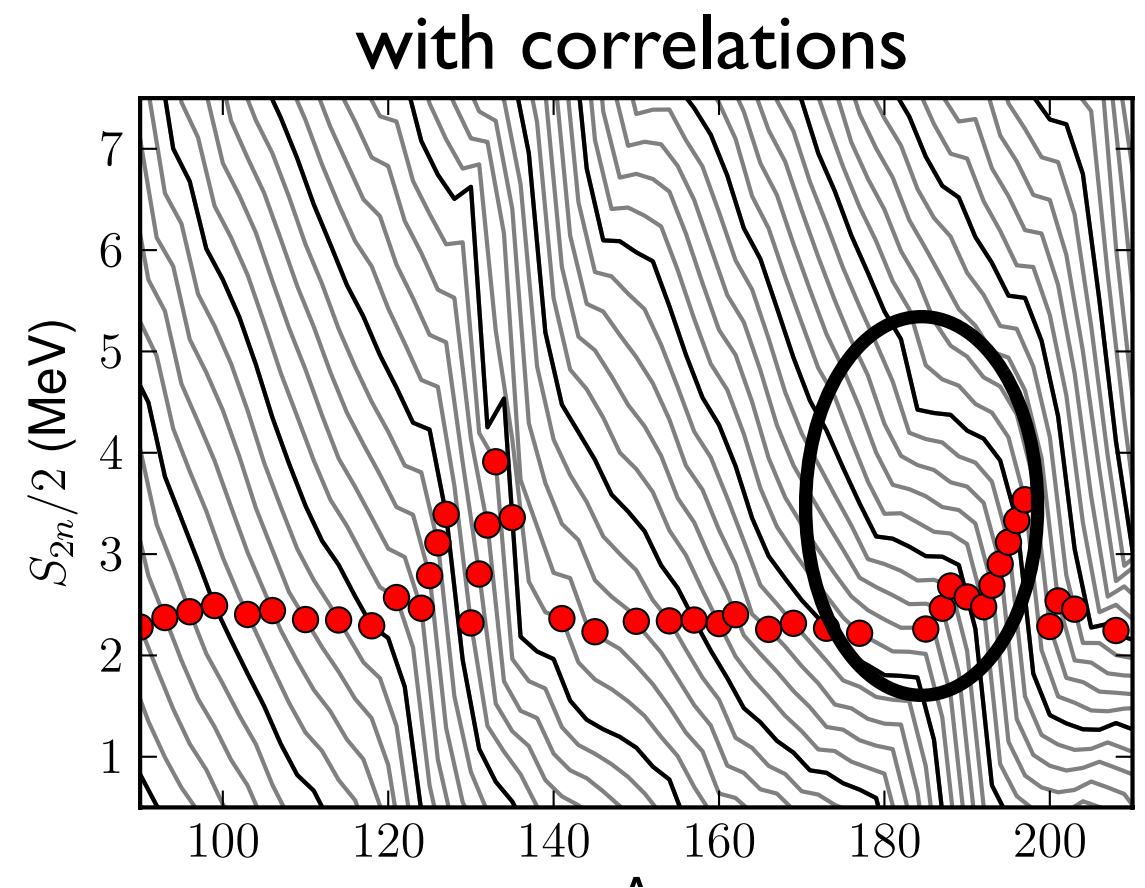
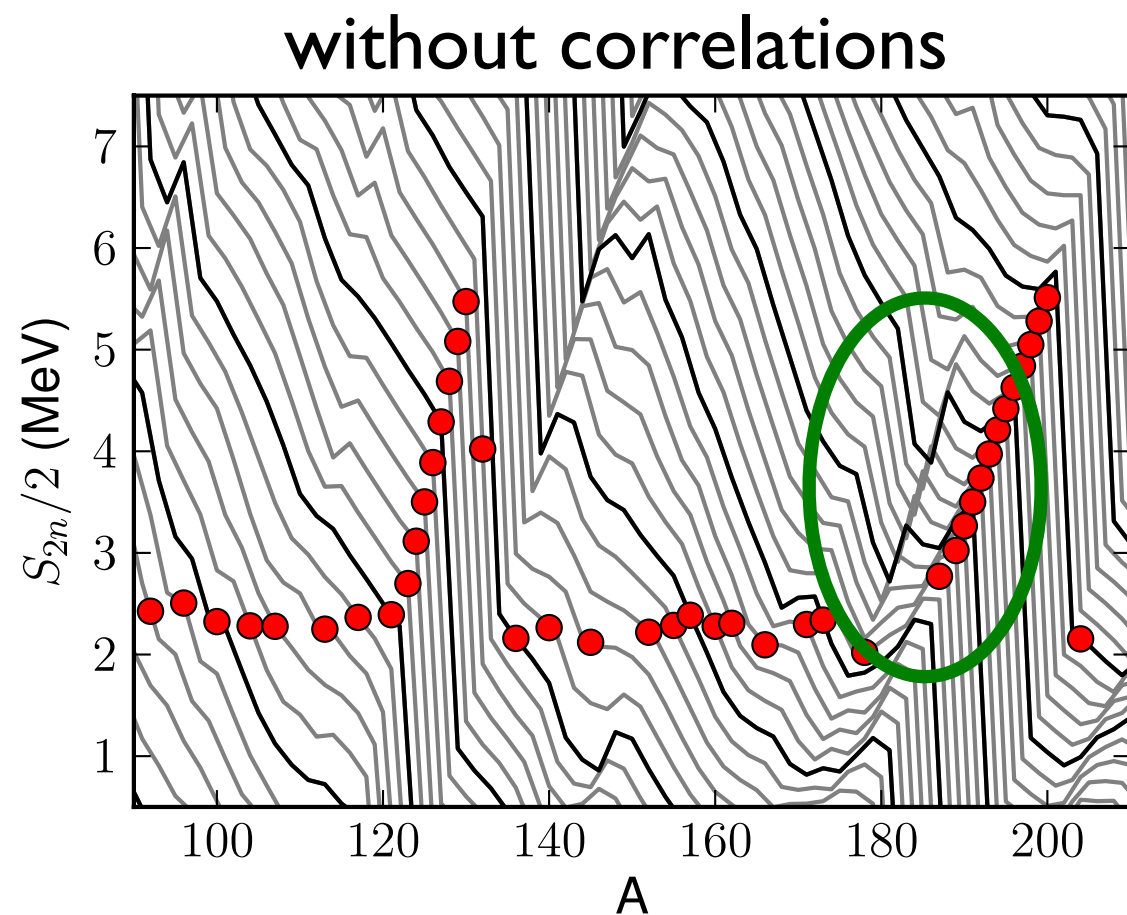
³Department of Physics and Institute of Nuclear Theory, Box 351560, University of Washington Seattle, Washington 98915, USA



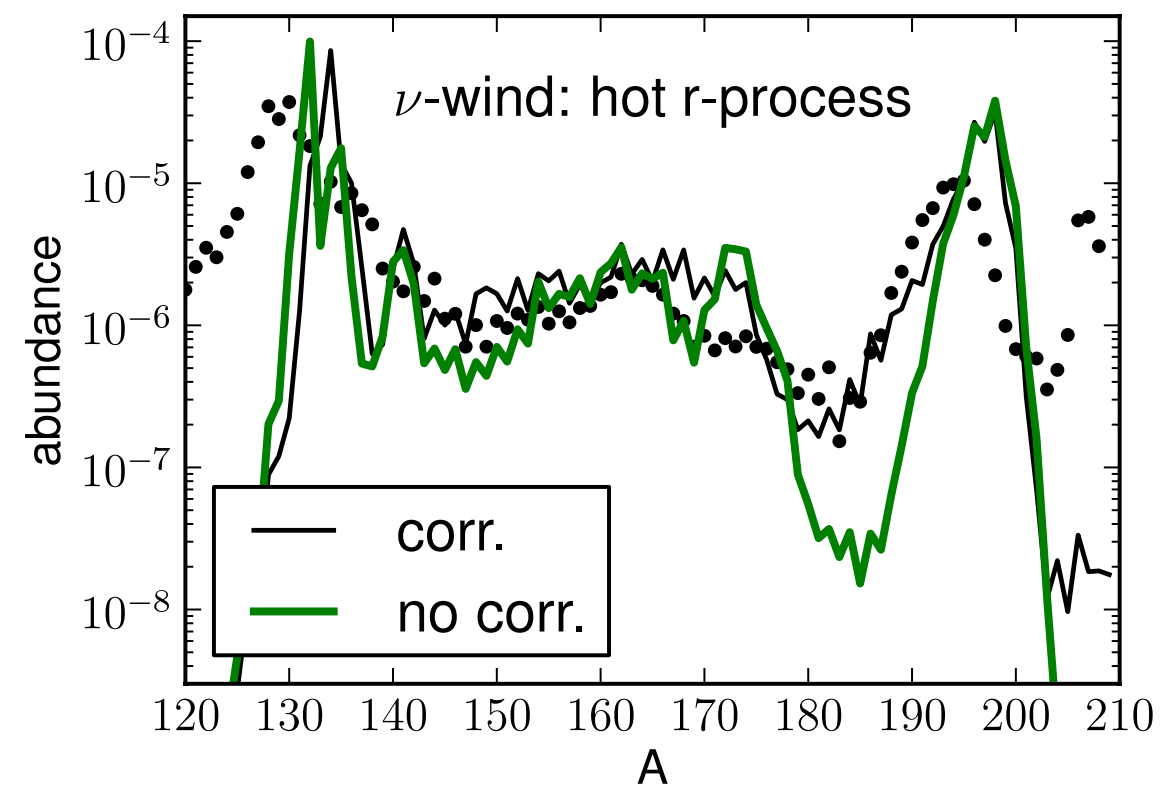
Impact of
nuclear correlations
on the r-process

Nuclear correlations and r-process

(Arcones & Bertsch,
arXiv:1111.4923)



nuclear correlations: strong impact
on trough before third peak!



Decay to stability

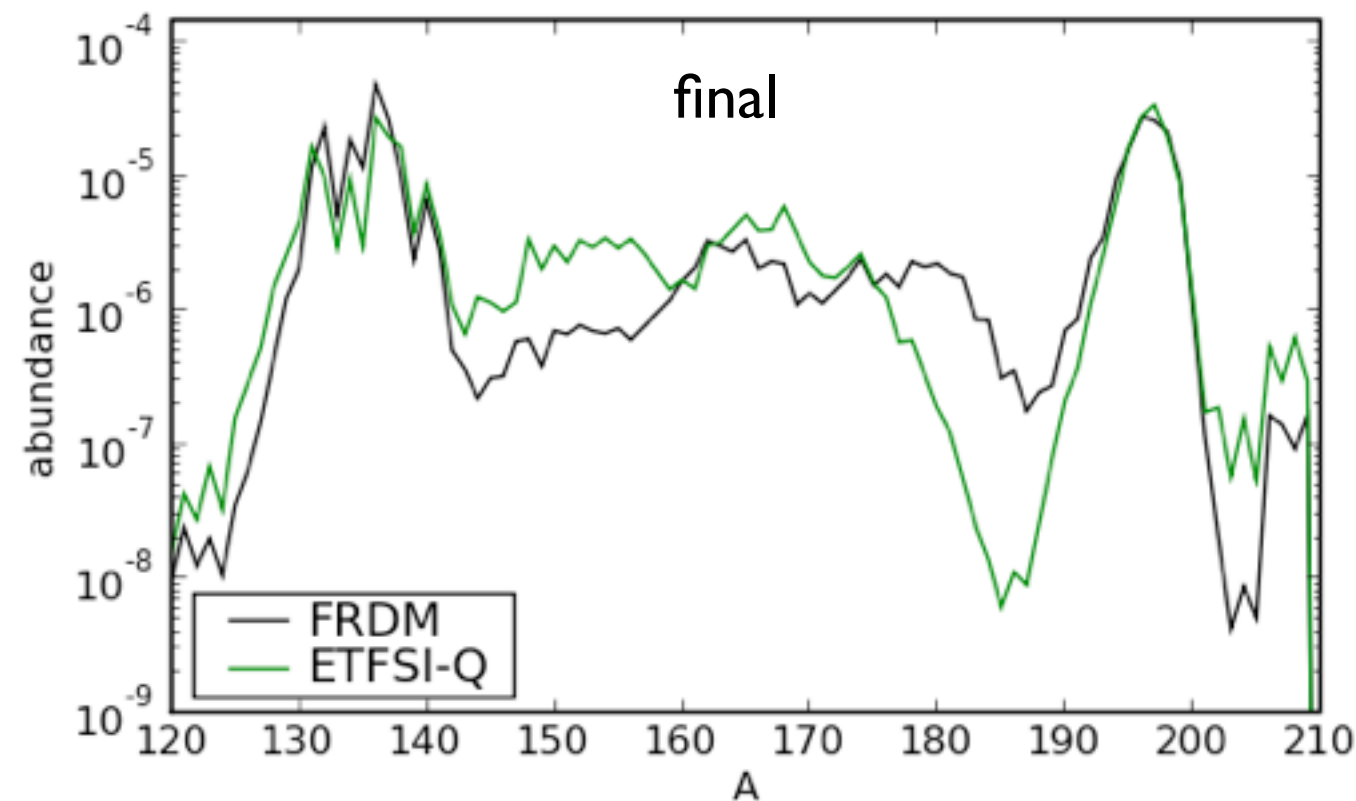
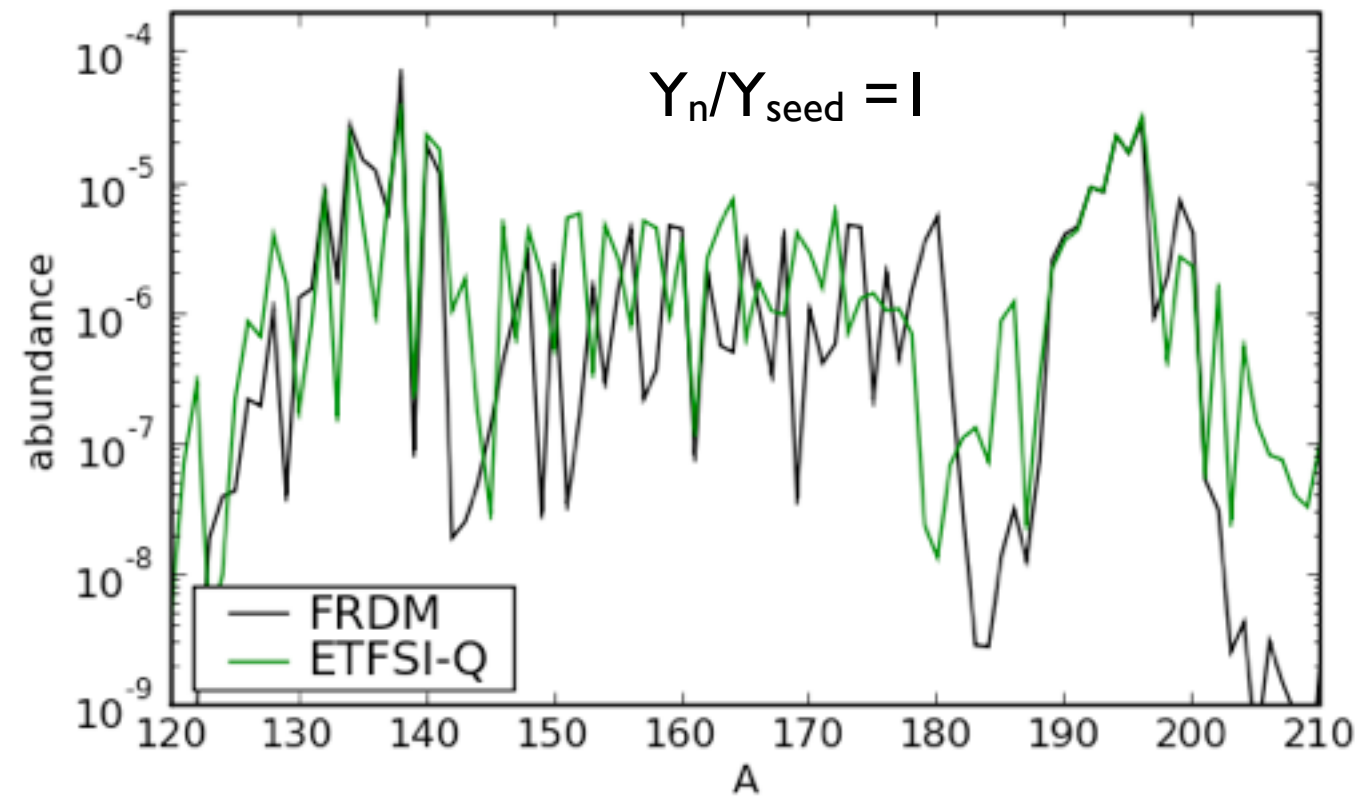
Abundances at freeze-out ($Y_n/Y_{\text{seed}}=1$):
odd-even effects

Final abundances are smoother like solar abundances.

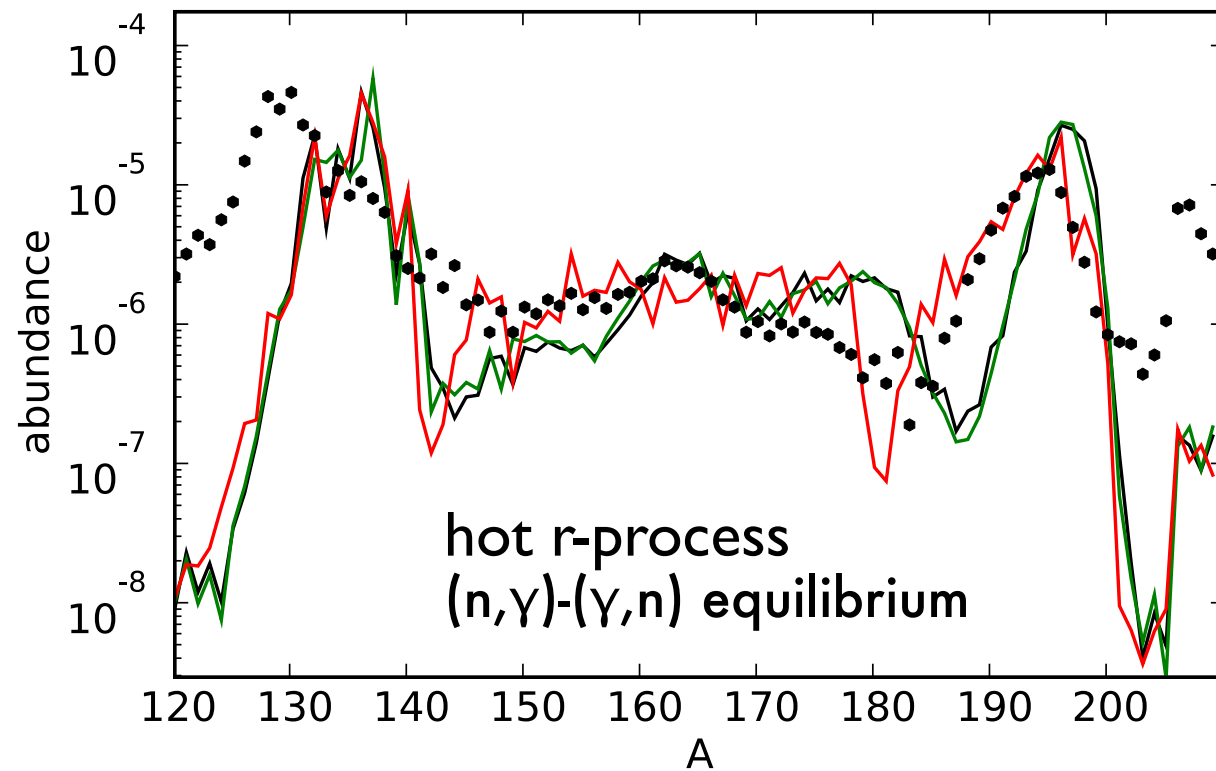
Why does the abundance pattern change?

Classical r-process (waiting point approximation): beta-delayed neutron emission (Kodama & Takahashi 1973, Kratz et al. 1993)

Dynamical r-process: **neutron capture** and **beta-delayed neutron emission** (Surman et al. 1997, Surman & Engel 2001, Surman et al. 2009, Buen et al. 2009)

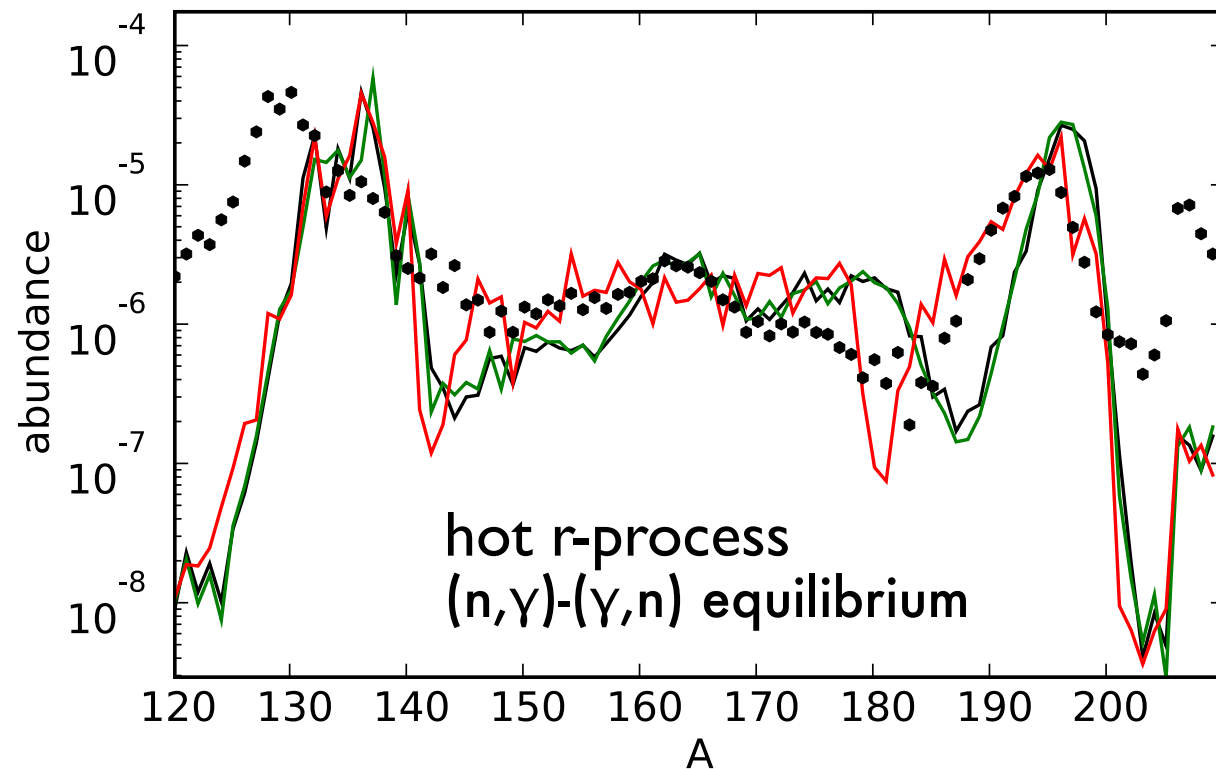


Neutron captures and beta-delayed neutron emission



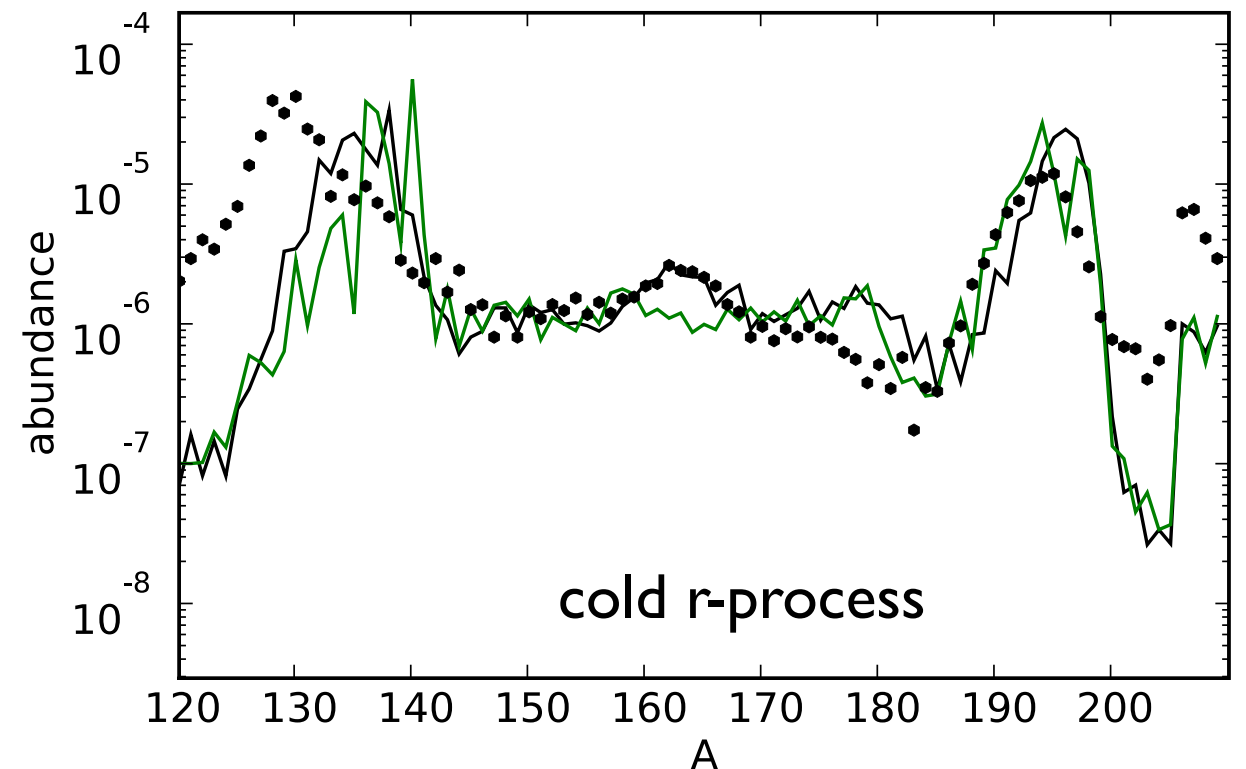
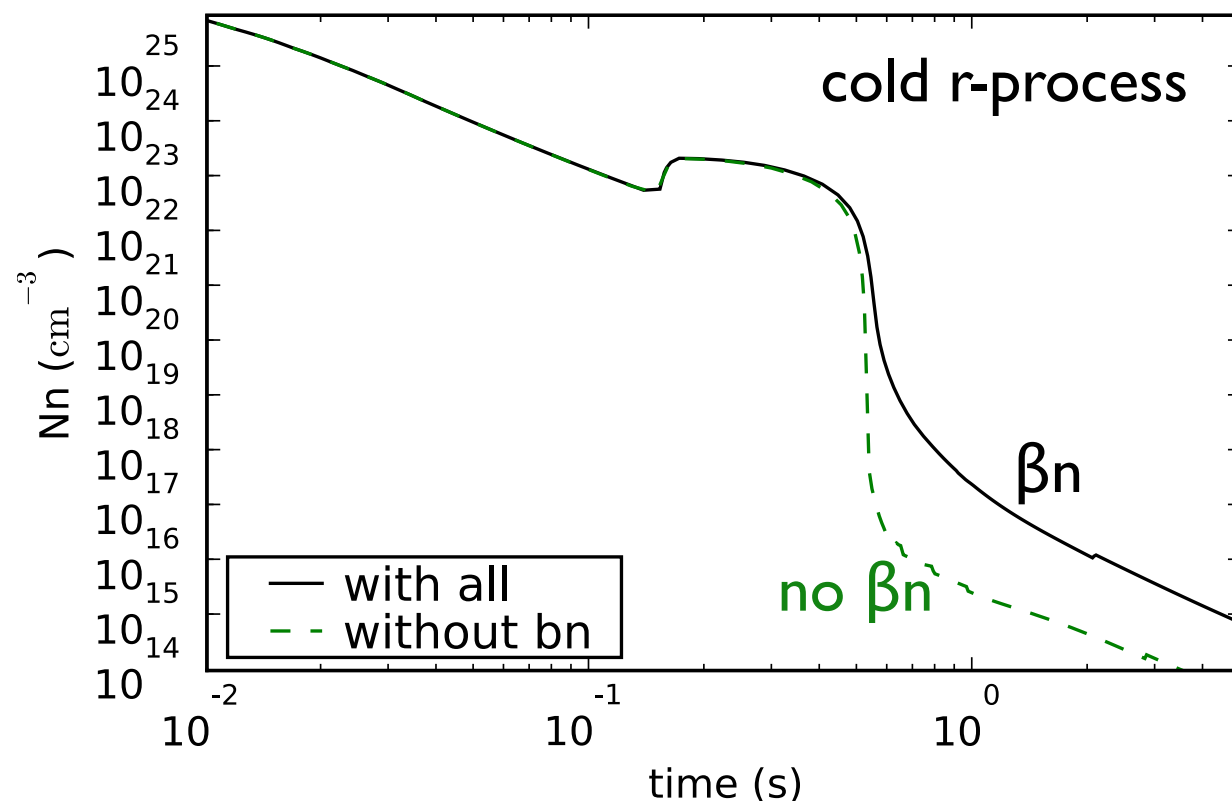
We compare final abundances with and without beta-delayed neutron emission and with and without neutron captures after freeze-out.

Neutron captures and beta-delayed neutron emission

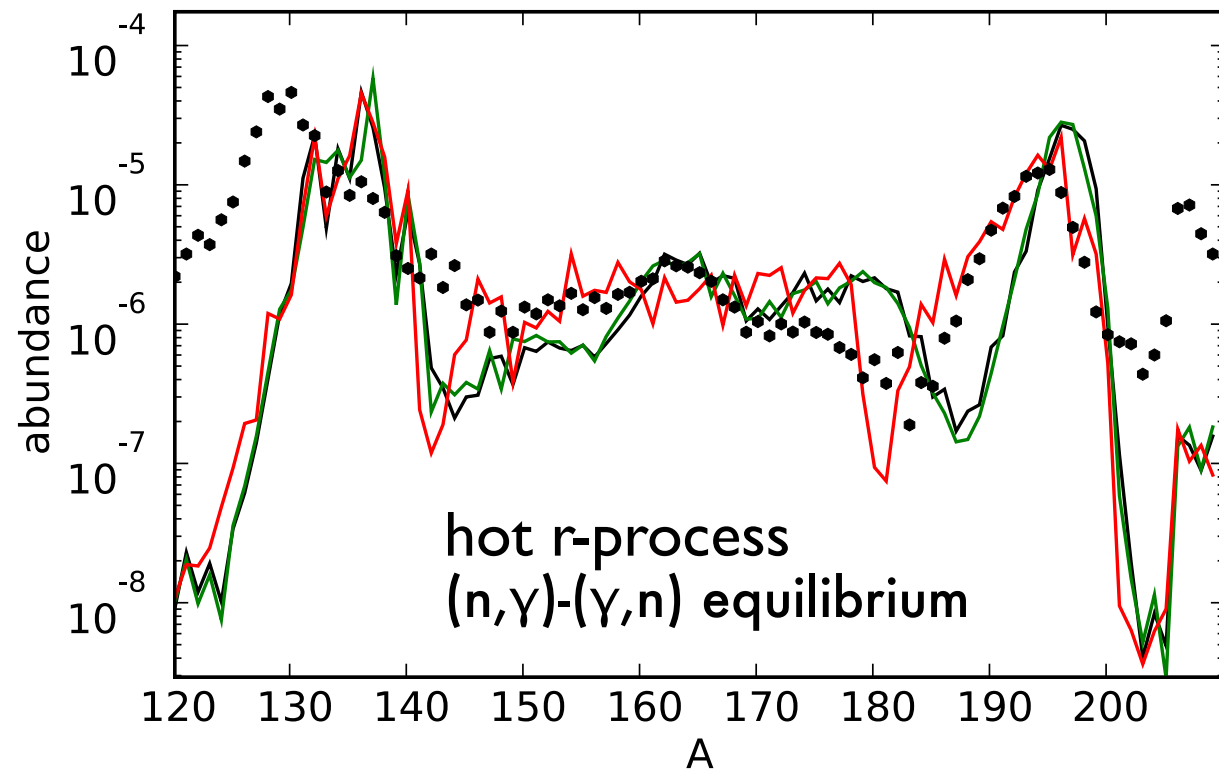


We compare final abundances with and without beta-delayed neutron emission and with and without neutron captures after freeze-out.

The main role of the beta-delayed neutron emission is to supply neutrons.

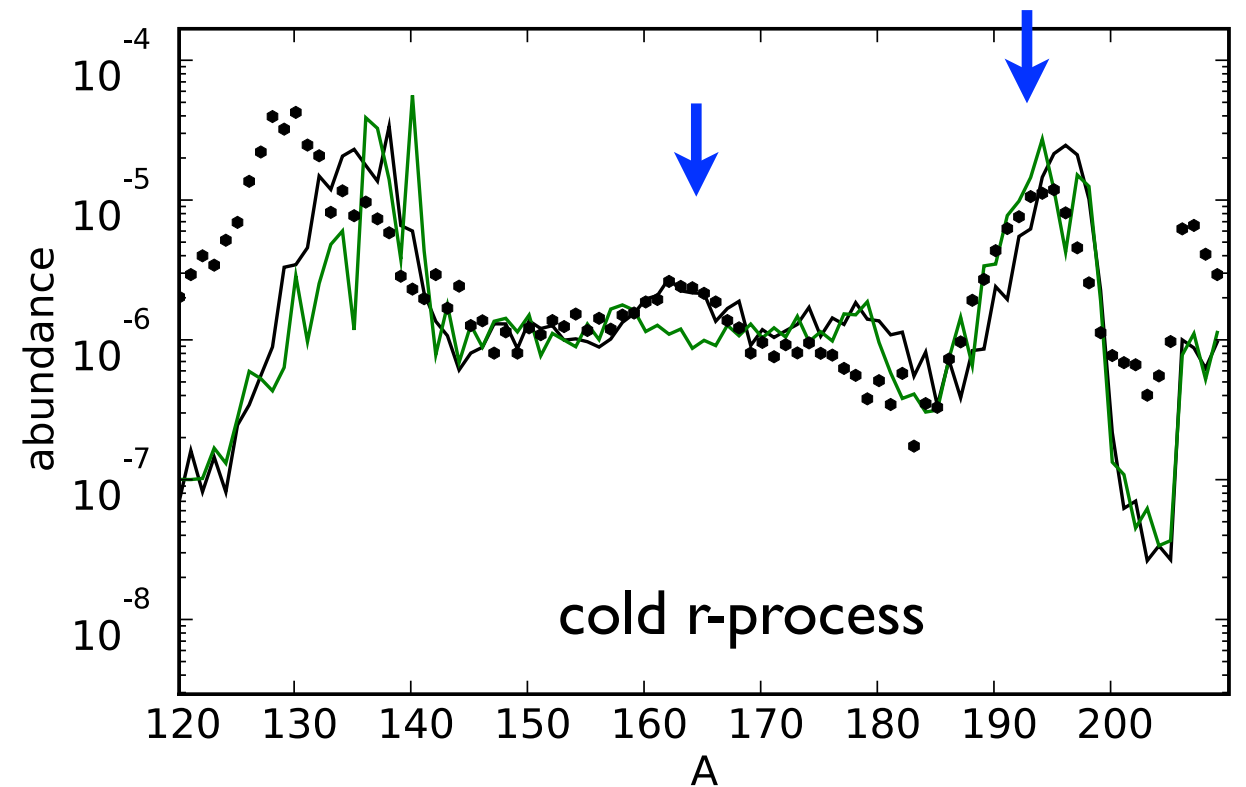
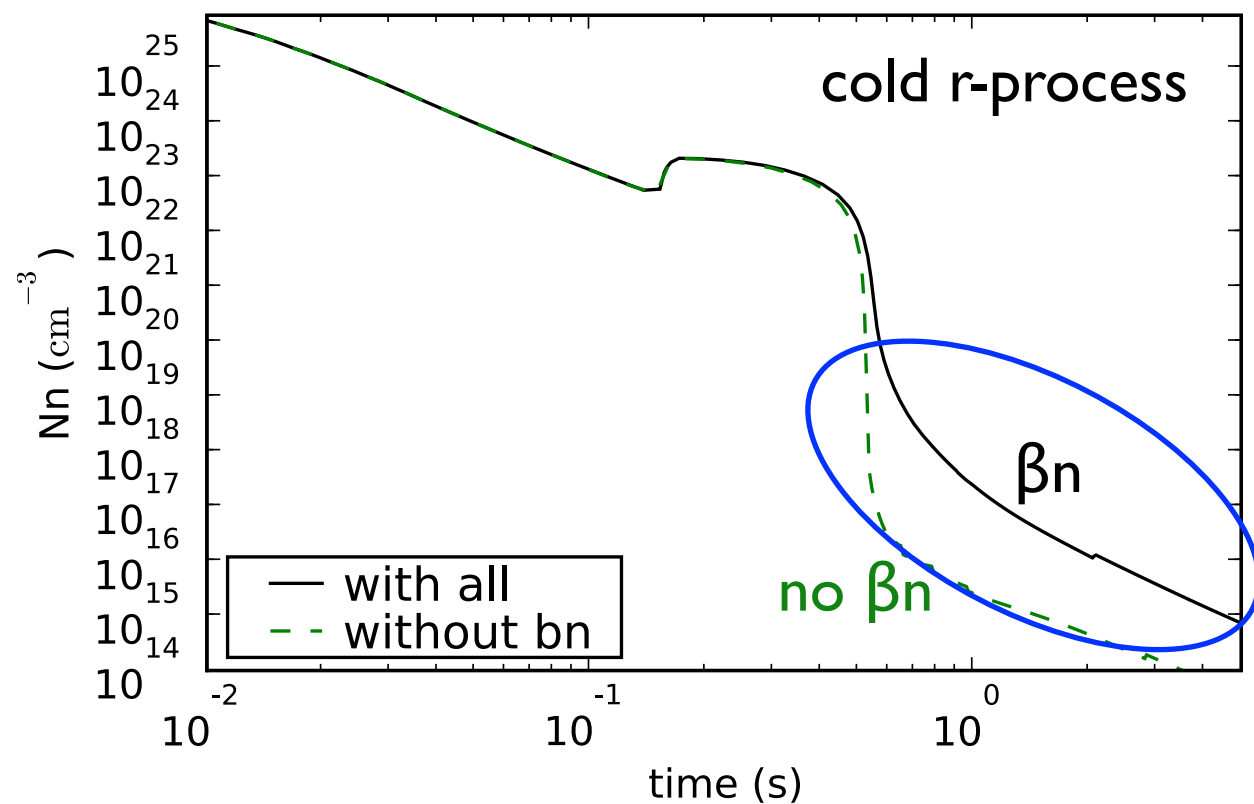


Neutron captures and beta-delayed neutron emission



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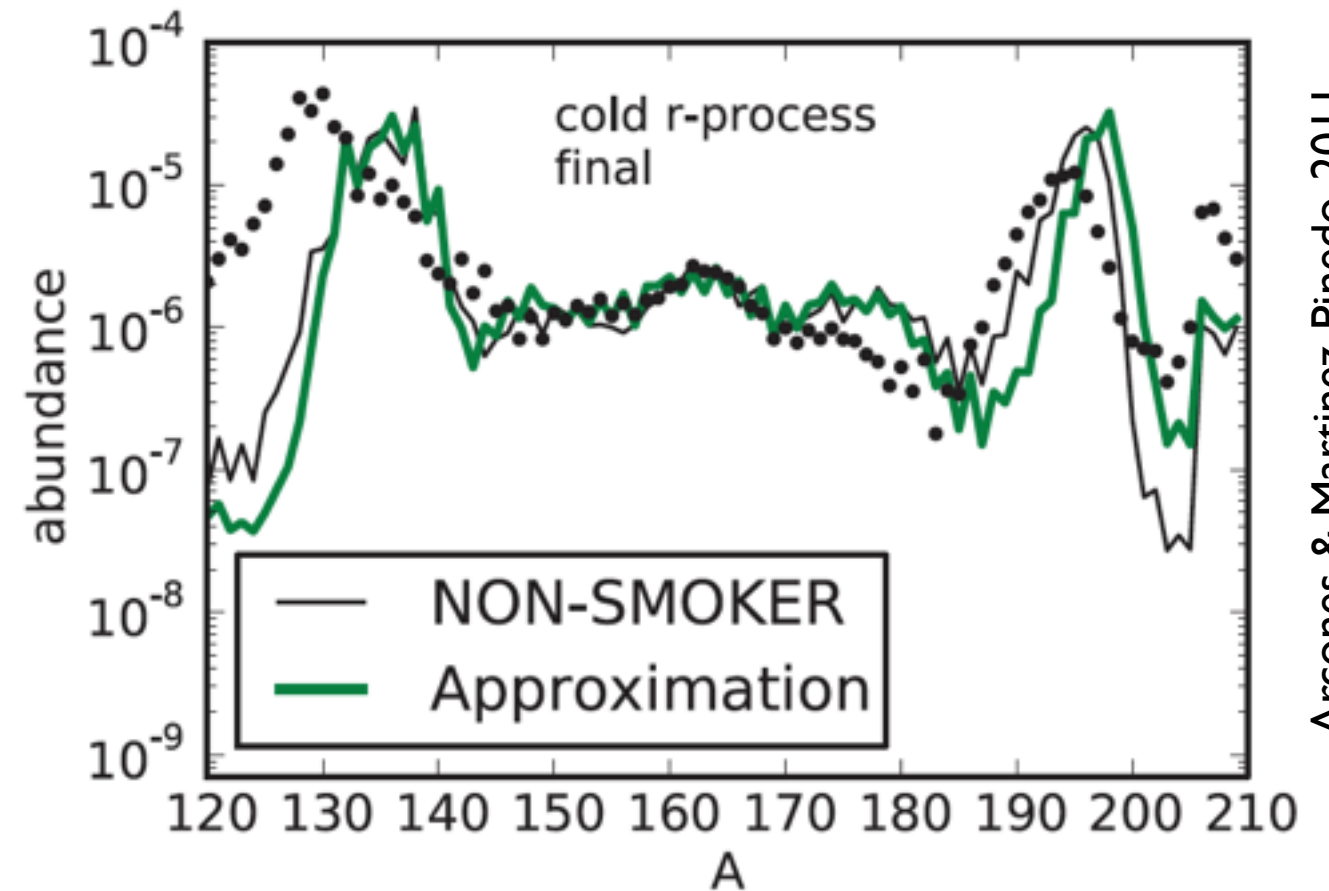
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Neutron captures

Compare neutron capture calculations:

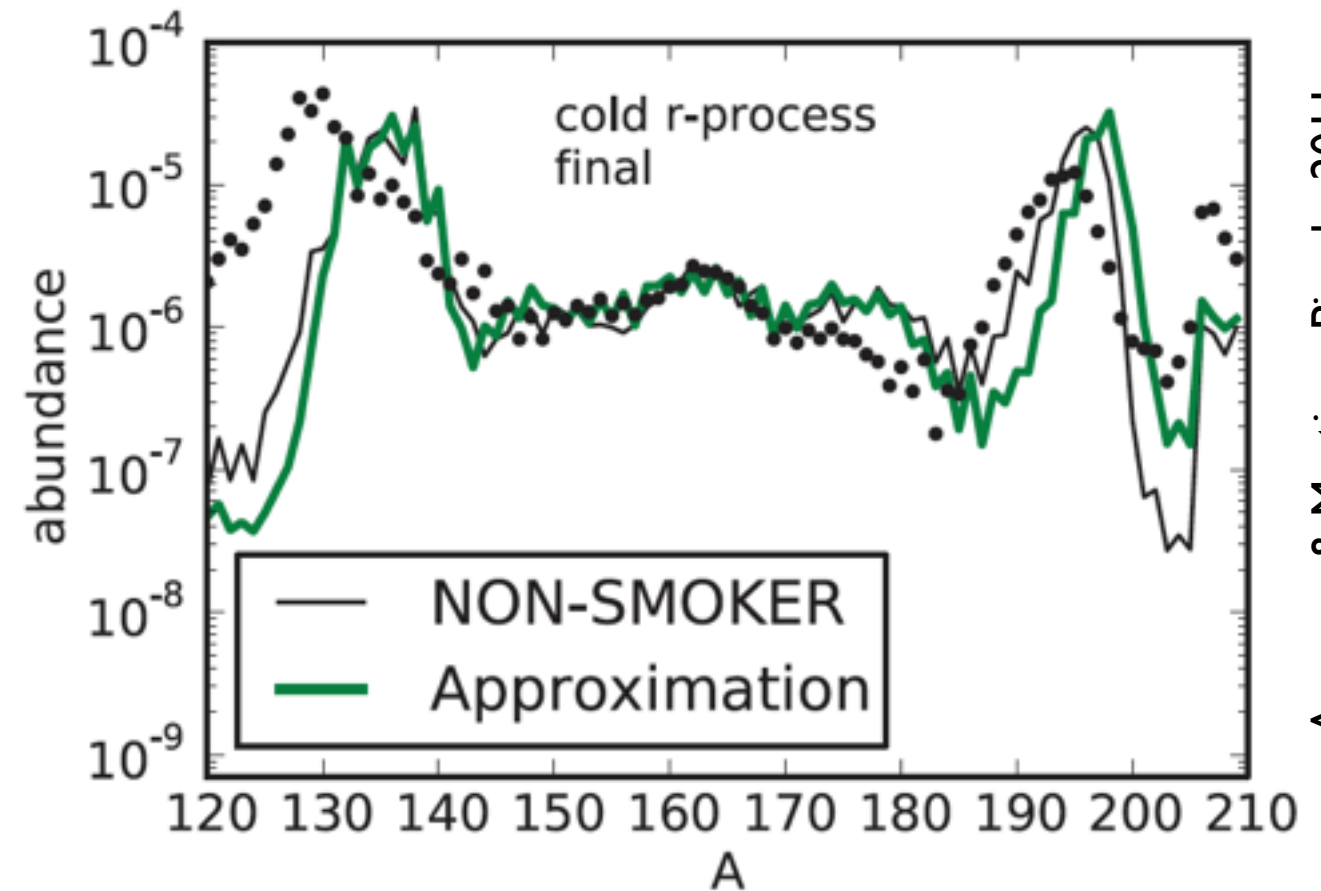
- NON-SMOKER
(Rauscher & Thielemann, 2000)
- Approximation
(Woosley, Fowler et al. 1975)



Neutron captures

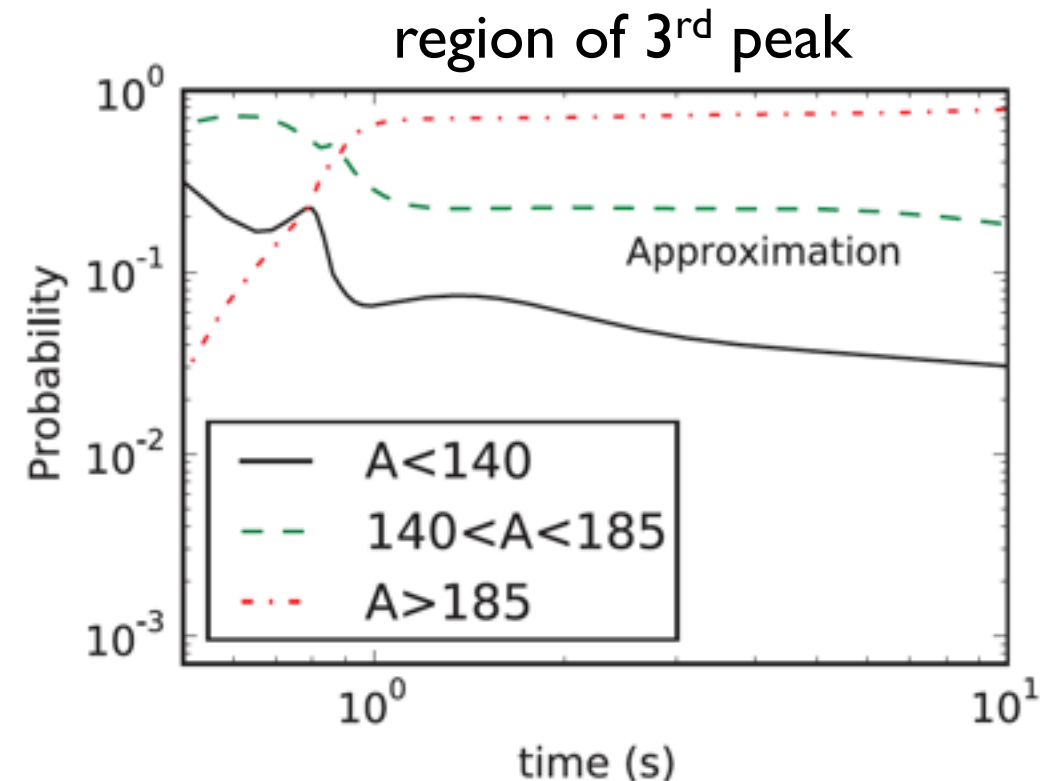
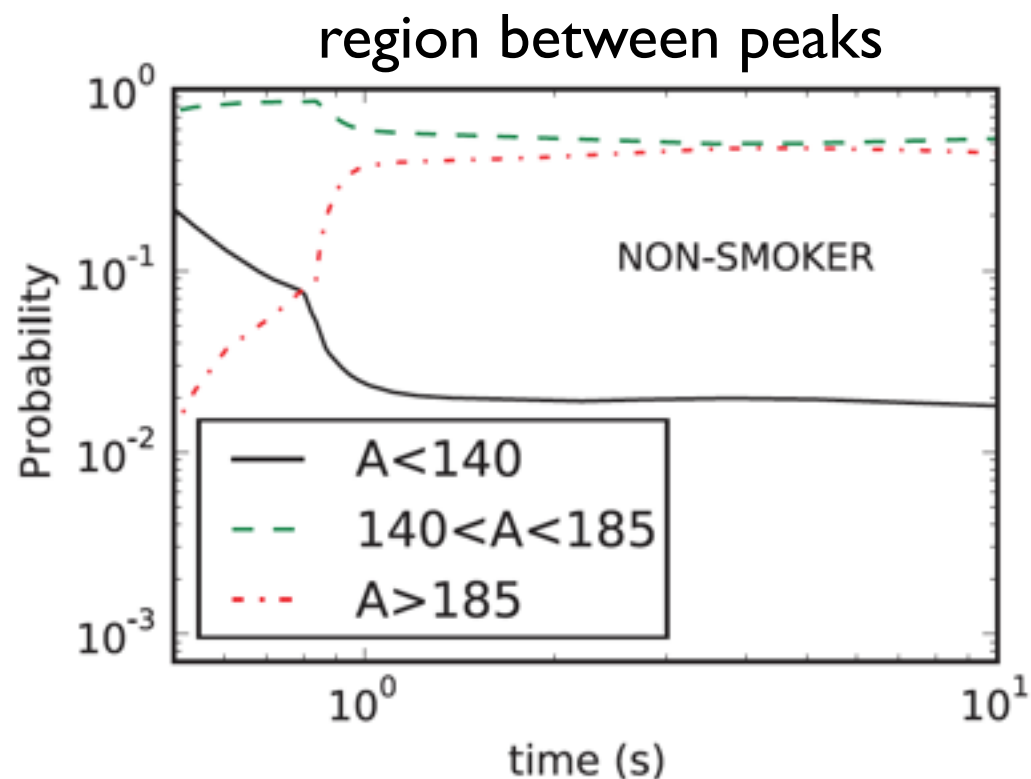
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Arcones & Martinez-Pinedo, 2011

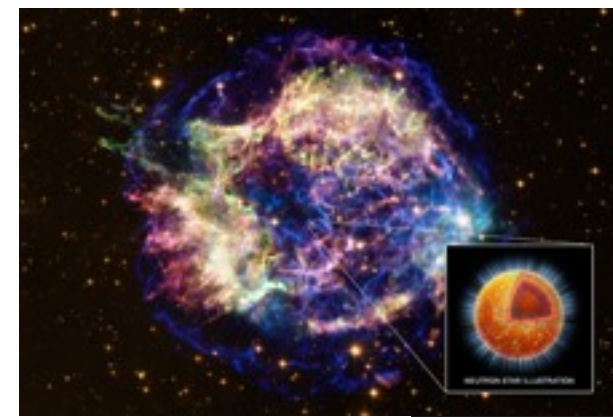
Neutron capture probability:



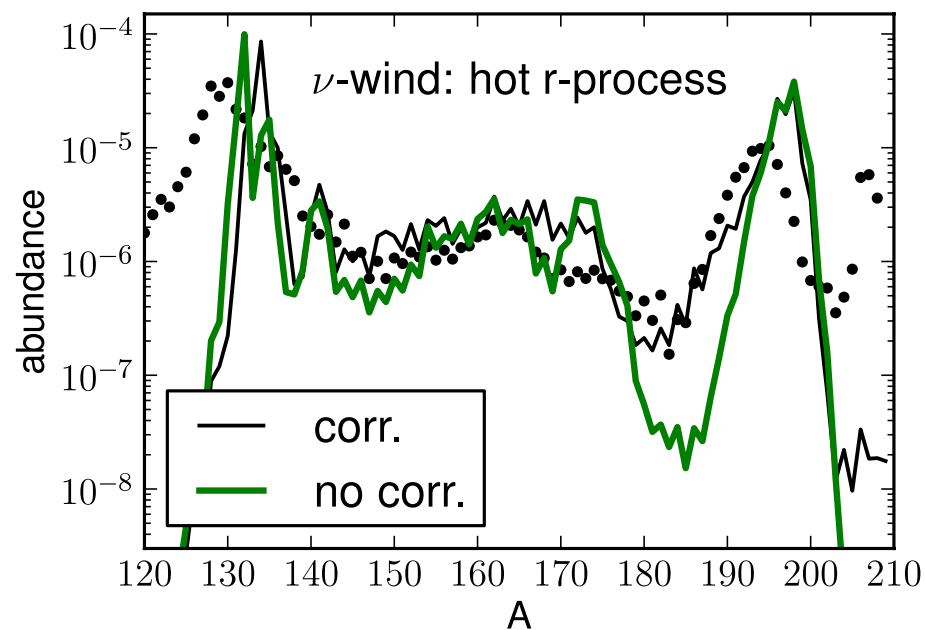
Conclusions

Where is the r-process?

Not found in recent supernova simulations



Long-time evolution and nuclear masses have big impact



Nuclear correlations:
masses in transition regions
from deformed to spherical
→ trough before 3rd peak

Decay to stability: beta-delayed neutron emission and
neutron captures still change the abundances