

Monte Carlo Rate Propagation for Experimentalists

R. Longland,
C. Iliadis, A.E. Champagne, J.R. Newton, C. Ugalde,
A. Coc, and R. Fitzgerald

Universitat Politècnica de Catalunya
Grup d'Astronomia i Astrofísica

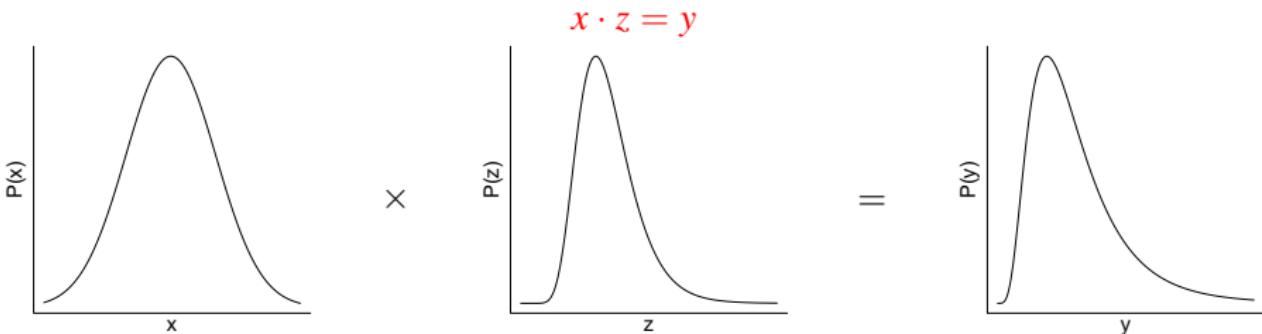
November 25, 2011



Monte Carlo Error Propagation

Reaction Rate per Particle Pair

$$\langle \sigma v \rangle = \sqrt{\frac{8}{\pi \mu}} \frac{1}{(kT)^{3/2}} \int_0^{\infty} E \sigma(E) e^{-E/kT} dE$$



- Draw a random number from each distribution
- Perform calculation
- Repeat (1) and (2) many times
- Analyse final y -distribution

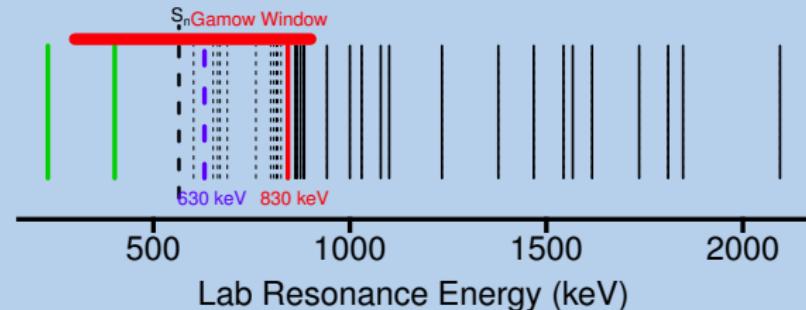
$$x = 0.5, z = 3.2$$

$$y = 1.6$$

$$y = 1.5, 2.3, 1.2, 1.4, 1.2, 1.8, 1.6, 1.5, 2.0$$

Median $y = 1.6$

A Fictional Example



• Directly Measured Resonances

- ▶ $E_r^{\text{lab}} = 830 \text{ keV } J^\pi = 2^+$
 $\Gamma_\alpha = 2.8 \times 10^{-5} \text{ eV}, \Gamma_\gamma = 3.0(15) \text{ eV},$
 $\Gamma_n = 2.5(17) \times 10^2 \text{ eV}$

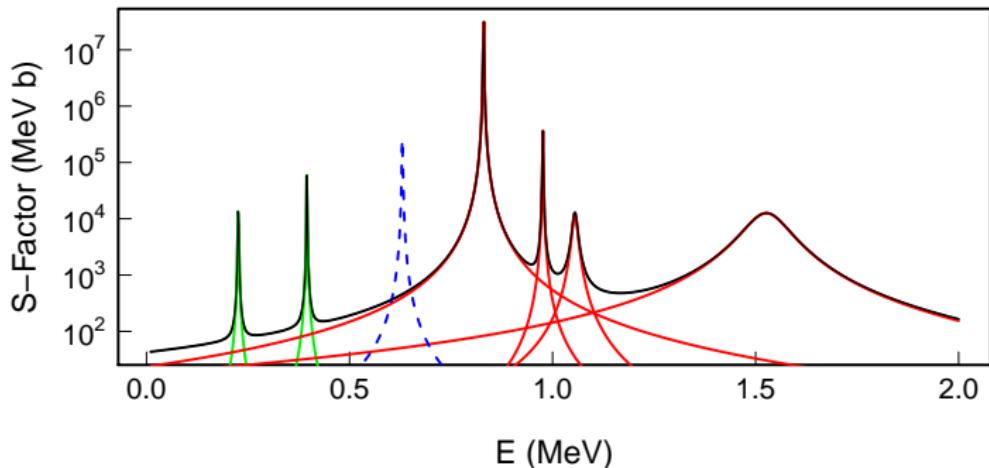
• Indirectly Measured Resonances

- ▶ $E_r^{\text{lab}} = 226 \text{ keV } J^\pi = 1^-$
 $S_\alpha = 1.9 \times 10^{-2}, \Gamma_\gamma = 0.72(18) \text{ eV}$
- ▶ $E_r^{\text{lab}} = 395 \text{ keV } J^\pi = 1^-$
 $S_\alpha = 2.8 \times 10^{-3}, \Gamma_\gamma = 1.9(3) \text{ eV}$

• Upper Limit Resonances

- ▶ $E_r^{\text{lab}} = 630 \text{ keV } J^\pi = 1^-$
 $\omega\gamma < 6 \times 10^{-3} \text{ neV}$
 $\Gamma_\gamma = 3.0(15) \text{ eV},$
 $\Gamma_n = 1.7(8) \times 10^3 \text{ eV},$
 $\Gamma_\alpha < 1.7 \times 10^{-9} \text{ eV}$

A Fictional Example



- Directly Measured Resonances

- ▶ $E_r^{\text{lab}} = 830 \text{ keV } J^\pi = 2^+$
 $\Gamma_\alpha = 2.8 \times 10^{-5} \text{ eV}, \Gamma_\gamma = 3.0(15) \text{ eV},$
 $\Gamma_n = 2.5(17) \times 10^2 \text{ eV}$

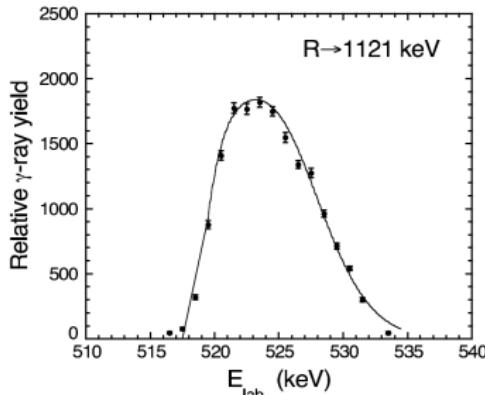
- Indirectly Measured Resonances

- ▶ $E_r^{\text{lab}} = 226 \text{ keV } J^\pi = 1^-$
 $S_\alpha = 1.9 \times 10^{-2}, \Gamma_\gamma = 0.72(18) \text{ eV}$
- ▶ $E_r^{\text{lab}} = 395 \text{ keV } J^\pi = 1^-$
 $S_\alpha = 2.8 \times 10^{-3}, \Gamma_\gamma = 1.9(3) \text{ eV}$

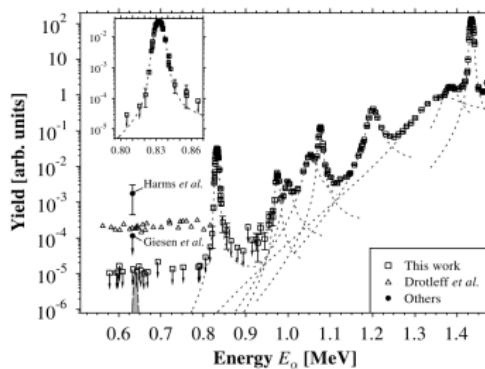
- Upper Limit Resonances

- ▶ $E_r^{\text{lab}} = 630 \text{ keV } J^\pi = 1^-$
 $\omega\gamma < 6 \times 10^{-3} \text{ neV}$
 $\Gamma_\gamma = 3.0(15) \text{ eV},$
 $\Gamma_n = 1.7(8) \times 10^3 \text{ eV},$
 $\Gamma_\alpha < 1.7 \times 10^{-9} \text{ eV}$

Directly Measured Resonances



Fox, C. et al., Phys. Rev. C 71 (2005) 055801



Jaeger, M. et al., Phys. Rev. Lett. 87 (2001) 202501

• Narrow Resonances

- ▶ Thick targets
- ▶ Precise resonance energy
- ▶ Resonance strength $\omega\gamma$

$$\omega\gamma = \frac{2\varepsilon_{\text{eff}}}{\lambda_r^2} \frac{N_{\max}}{N_b B \eta W}$$

- ▶ Central Limit Theorem
Multiplication of quantities leads to Lognormal uncertainties.

• Wide Resonances

- ▶ Thin targets
- ▶ Total resonance width Γ
- ▶ Resonance strength $\omega\gamma$

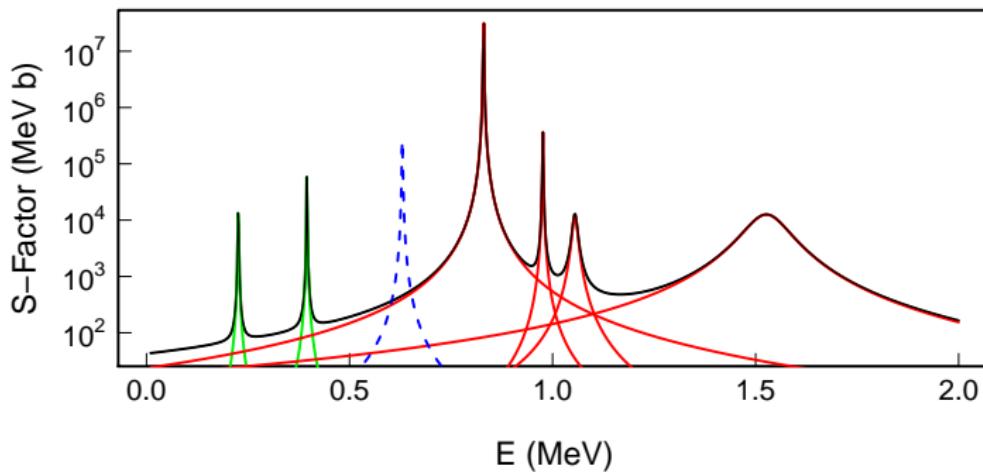
Wide Resonances Note

Cross Section

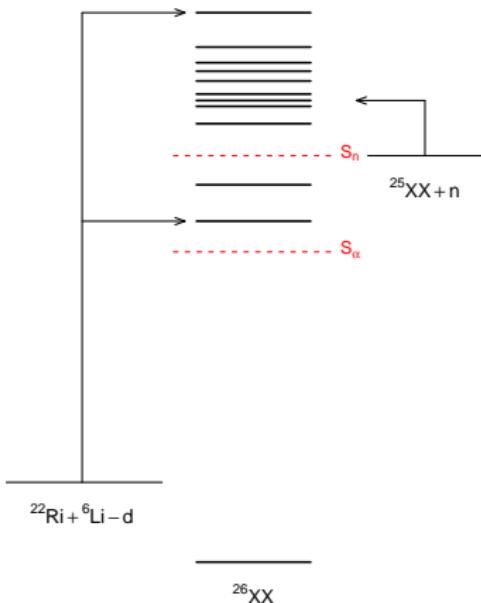
$$\sigma \sim \frac{\Gamma_a(E)\Gamma_b(E)}{(E_r - E)^2 + \Gamma(E)^2/4}$$

$$\Gamma(E) \sim P_c(E)\theta^2$$

$$\Gamma(E) = \Gamma(E_r) \frac{P_c(E)}{P_c(E_r)}$$



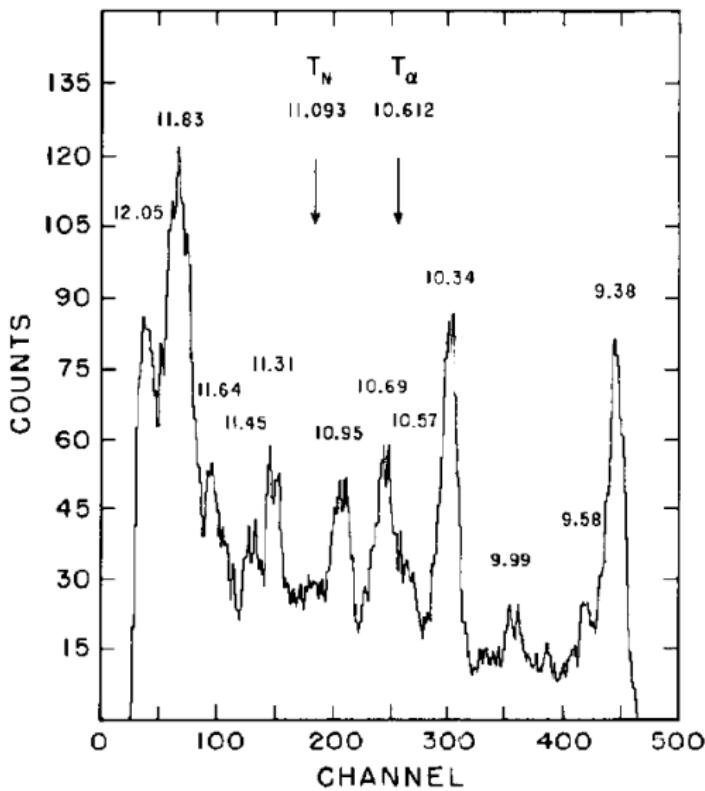
Indirectly Measured Resonances



- Transfer Reactions
 - ▶ Γ_α (S-factors, ANC's etc.)
 - ▶ Energies
 - ▶ Quantum Numbers
- Neutron Capture
 - ▶ Γ_n
 - ▶ Energies
 - ▶ Branching Ratios
- Photoexcitation
 - ▶ Γ_γ
 - ▶ Energies
 - ▶ Branching Ratios

- Indirect methods used mostly for low energy resonances
- Energy uncertainty can be large

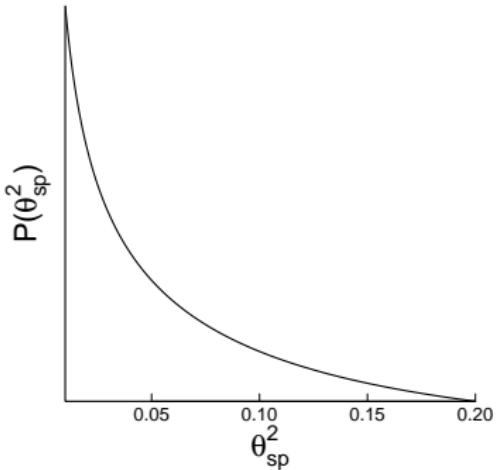
Indirectly Measured Resonances



Example: $^{22}\text{Ne}(\text{Li},\text{d})^{26}\text{Mg}$

- ($^6\text{Li},\text{d}$) reaction populates states either side of α -particle threshold
- Large energy uncertainties
- Upon subtracting S_α , one obtains resonance energy with finite probability of negative energy (sub-threshold resonance)
- Central Limit Theorem
Addition of quantities leads to Gaussian uncertainties.

Upper Limit Resonances



- Random matrix theory predicts that dimensionless reduced widths (θ_{sp}^2) should be distributed *locally* according to a Porter-Thomas distribution

$$P(\theta^2) = \begin{cases} \frac{c}{\sqrt{\theta^2}} e^{-\theta^2/(2\langle\theta^2\rangle)} & \text{if } \theta^2 \leq \theta_{ul}^2 \\ 0 & \text{if } \theta^2 > \theta_{ul}^2 \end{cases}$$

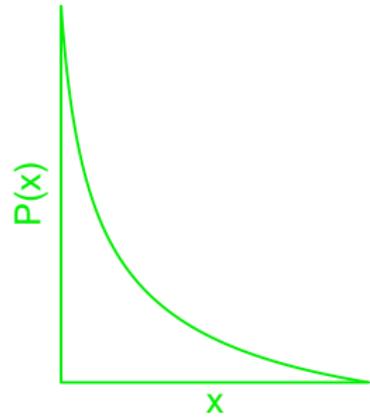
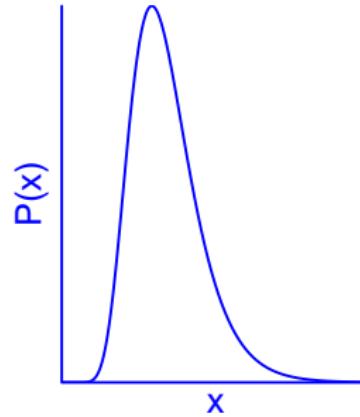
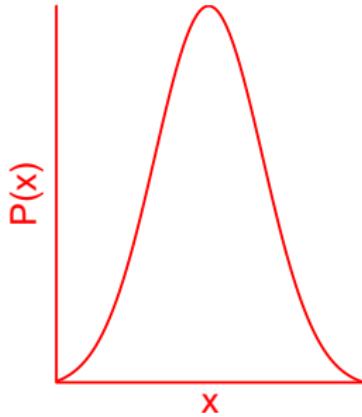
- To estimate $\langle\theta^2\rangle$ for $A = 20 - 40$
 - 360 α -particle widths
 $\langle\theta_\alpha^2\rangle = 0.010$
 - 1127 proton widths
 $\langle\theta_p^2\rangle = 0.0045$
- Nuclear theory and experiments are required

Reaction Rate

Reaction Rate per Particle Pair

$$\langle \sigma v \rangle = \sqrt{\frac{8}{\pi \mu}} \frac{1}{(kT)^{3/2}} \int_0^{\infty} E \sigma(E) e^{-E/kT} dE$$

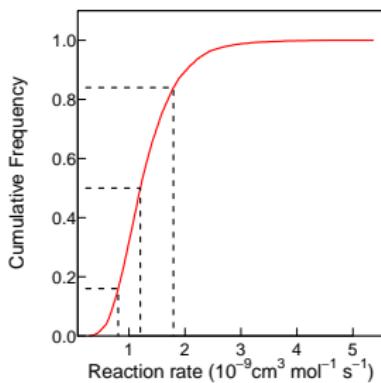
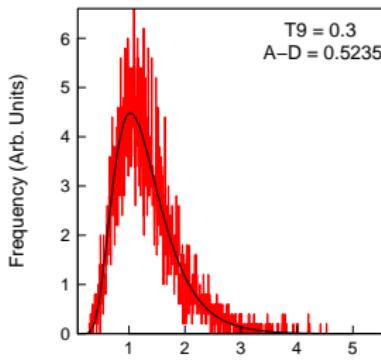
E_r	$\omega\gamma$	$S, S'S''$	Γ_a	Γ_{UL}	$\langle \sigma v \rangle$
Gaussian	Lognormal	Lognormal	Lognormal	Porter-Thomas	Lognormal



Putting it all Together

```
22Ri(a,g)26XX
*****
2          ! Zproj
12         ! Ztarget
0          ! Zexitparticle (=0 when only 2 channels open)
4.003      ! Aproj
21.991     ! Atarget
1.009      ! Aexitparticle (=0 when only 2 channels open)
0.0        ! Jproj
0.0        ! Jtarget
0.5        ! Jexitparticle (=0 when only 2 channels open)
10614.79   ! projectile separation energy (keV)
11093.08   ! exit particle separation energy (=0 when only 2 channels open)
1.25       ! Radius parameter R0 (fm)
2          ! Gamma-ray channel number (-2 if ejectile is a g-ray; -3 otherwise)
*****
1.0        ! Minimum energy for numerical integration (keV)
5000       ! Number of random samples (>5000 for better statistics)
0          ! =0 for rate output at all temperatures; =NT for rate output at selected temperatures
*****
Nonresonant Contribution
S(keVb) S'(b) S''(b/keV)  fracErr  Cutoff Energy (keV)
0.0 0.0 0.0 0.0 0.0
0.0 0.0 0.0 0.0 0.0
*****
Resonant Contribution
Note: G1 = entrance channel, G2 = exit channel, G3 = spectator channel !! Ecm, Exf in (keV); wg, Gx in (eV) !!
Note: if Er<0, theta^2=C2S*theta_sp^2 must be entered instead of entrance channel partial width
Ecm  DEcm  wg    Dwg   J  G1    DG1  L1  G2  DG2  L2  G3  DG3  L3  Exf  Int
191.08  0.15  0     0     1  1.25e-23  0.50e-23  1   0.7  0.2  1   0     0     0     0     0
334.31  0.1   0     0     1  1.20e-9   0.50e-9   1   1.9  0.3  1   0     0     0     0     0
703.78  2.11  0     0     2  2.8e-5   0.5e-5   2   3.0  1.5  1   2.5e2  1.7e2  1   0     1
826.04  0.19  0     0     4  3.78e-6  4.44e-7  4   3.0  1.5  1   1.47e3  8.0e1  2   0     1
893.31  0.90  0     0     1  1.17e-4  2.0e-5  1   3.0  1.5  1   1.27e4  2.5e3  1   0     1
1294.6   1.20  0     0     2  2.80e-6  0.1e-6  2   2.3  1.0  1   1.10e5  0.05e5  1   0     1
*****
Upper Limits of Resonances
Note: enter partial width upper limit by choosing non-zero value for PT, where PT=<theta^2> for particles and...
Note: ...PT=<B> for g-rays [enter: "upper_limit 0.0"]; for each resonance: # upper limits < # open channels!
Ecm  DEcm  Jr  G1    DG1  L1  PT  G2  DG2  L2  PT  G3  DG3  L3  PT  Exf  Int
533.1   0.8   1   1.7e-9  0   1   0.01  3.0  1.5   1   0   1.7e3  0.8e3  1   0   0   0
*****
```

Rate Distribution



- Distributions of rates are calculated
- Useful statistical parameters
 - Median
 - 1σ uncertainties (16th and 84th percentiles)
- Entire distribution of rates described by *lognormal distribution*

$$\mu = \ln(E[x]) - \frac{1}{2} \ln \left(1 + \frac{V[x]}{E[x]^2} \right), \quad \sigma = \sqrt{\ln \left(1 + \frac{V[x]}{E[x]^2} \right)}$$

- Lognormal distribution is an assumption
- Anderson-Darling statistic describes how well distribution is fit by the lognormal

$$t_{AD} = -n - \sum_{i=1}^n \frac{2i-1}{n} (\ln F(y_i) + \ln [1 - F(y_{n+1-i})])$$

- $t_{AD} < 40$ corresponds to good *visual* agreement

Stephens, M.A., J. Am. Stat. Assoc. 69 (1973) 74

Rate Presentation

T (GK)	Low rate	Median rate	High rate	lognormal μ	lognormal σ	A-D
0.100	2.47×10^{-20}	3.67×10^{-20}	5.48×10^{-20}	$-4.475 \times 10^{+01}$	4.06×10^{-01}	5.29×10^{-01}
0.110	7.28×10^{-19}	1.08×10^{-18}	1.62×10^{-18}	$-4.137 \times 10^{+01}$	4.06×10^{-01}	5.29×10^{-01}
0.120	1.21×10^{-17}	1.79×10^{-17}	2.68×10^{-17}	$-3.856 \times 10^{+01}$	4.06×10^{-01}	5.30×10^{-01}
0.130	1.29×10^{-16}	1.91×10^{-16}	2.86×10^{-16}	$-3.619 \times 10^{+01}$	4.06×10^{-01}	5.30×10^{-01}
0.140	9.72×10^{-16}	1.44×10^{-15}	2.16×10^{-15}	$-3.417 \times 10^{+01}$	4.06×10^{-01}	5.30×10^{-01}
0.150	5.56×10^{-15}	8.25×10^{-15}	1.23×10^{-14}	$-3.243 \times 10^{+01}$	4.06×10^{-01}	5.30×10^{-01}
0.160	2.54×10^{-14}	3.77×10^{-14}	5.64×10^{-14}	$-3.091 \times 10^{+01}$	4.06×10^{-01}	5.29×10^{-01}
0.180	3.15×10^{-13}	4.68×10^{-13}	6.98×10^{-13}	$-2.839 \times 10^{+01}$	4.06×10^{-01}	5.29×10^{-01}
0.200	2.32×10^{-12}	3.45×10^{-12}	5.15×10^{-12}	$-2.639 \times 10^{+01}$	4.06×10^{-01}	5.29×10^{-01}
0.250	8.05×10^{-11}	1.19×10^{-10}	1.78×10^{-10}	$-2.285 \times 10^{+01}$	4.06×10^{-01}	5.28×10^{-01}
0.300	8.14×10^{-10}	1.21×10^{-09}	1.80×10^{-09}	$-2.053 \times 10^{+01}$	4.06×10^{-01}	5.23×10^{-01}
0.350	4.11×10^{-09}	6.10×10^{-09}	9.08×10^{-09}	$-1.891 \times 10^{+01}$	4.05×10^{-01}	4.96×10^{-01}
0.400	1.37×10^{-08}	2.02×10^{-08}	3.00×10^{-08}	$-1.772 \times 10^{+01}$	4.00×10^{-01}	4.35×10^{-01}
0.450	3.55×10^{-08}	5.14×10^{-08}	7.58×10^{-08}	$-1.678 \times 10^{+01}$	3.88×10^{-01}	4.13×10^{-01}
0.500	7.85×10^{-08}	1.12×10^{-07}	1.63×10^{-07}	$-1.600 \times 10^{+01}$	3.74×10^{-01}	4.65×10^{-01}
0.600	2.89×10^{-07}	4.16×10^{-07}	6.16×10^{-07}	$-1.468 \times 10^{+01}$	3.89×10^{-01}	$1.77 \times 10^{+00}$
0.700	8.29×10^{-07}	1.26×10^{-06}	2.06×10^{-06}	$-1.355 \times 10^{+01}$	4.74×10^{-01}	$1.10 \times 10^{+01}$
0.800	2.00×10^{-06}	3.29×10^{-06}	5.94×10^{-06}	$-1.258 \times 10^{+01}$	5.60×10^{-01}	$1.31 \times 10^{+01}$
0.900	4.26×10^{-06}	7.48×10^{-06}	1.44×10^{-05}	$-1.176 \times 10^{+01}$	6.18×10^{-01}	$1.00 \times 10^{+01}$
1.000	8.21×10^{-06}	1.51×10^{-05}	2.98×10^{-05}	$-1.106 \times 10^{+01}$	6.48×10^{-01}	$7.46 \times 10^{+00}$

Conclusions

- Statistically meaningful uncertainties of thermonuclear nuclear reaction rates formulated
- Treatment of upper limits based on nuclear physics theory
- Flexible method allows for reasonable uncertainty assignments
 - ▶ **Energies:** Gaussian
 - ▶ **Resonance Strengths:** Lognormal
 - ▶ **Partial Widths:** Lognormal
 - ▶ **Upper Limits:** Porter-Thomas Distribution
- Final uncertainties can be presented as lognormal parameters
- Allows for full Monte Carlo sensitivity studies to be performed

- See:
 - Longland, R. *et al.*, Nucl. Phys. A 841 (2010) 1–30
 - Iliadis, C. *et al.*, Nucl. Phys. A 841 (2010) 31–250
- Supported by:
 - ▶ US Department of Energy under Contract No. DE-FG02-97ER41041.

Monte Carlo Rate Propagation for Experimentalists

R. Longland,
C. Iliadis, A.E. Champagne, J.R. Newton, C. Ugalde,
A. Coc, and R. Fitzgerald

Universitat Politècnica de Catalunya
Grup d'Astronomia i Astrofísica

November 25, 2011

